

Lecture 16 : Fourier Series V

Note Title

★ Fourier Series on Intervals

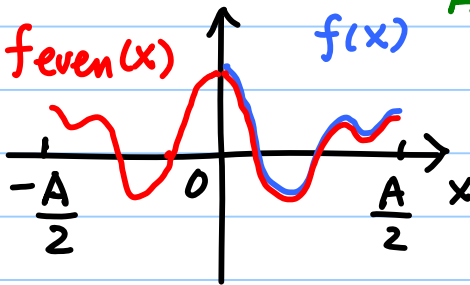
Suppose your fcn is defined on $[0, \frac{A}{2}]$ instead of $[-\frac{A}{2}, \frac{A}{2}]$.

$\Rightarrow \equiv$ two ways to make it an A -periodic fcn:

(1) Even extension

$$f_{\text{even}}(x) = \begin{cases} f(x) & \text{if } x \in [0, \frac{A}{2}] \\ f(-x) & \text{if } x \in [-\frac{A}{2}, 0] \end{cases}$$

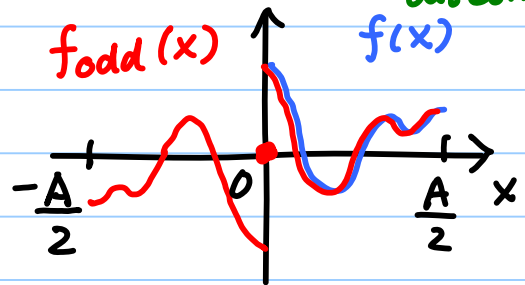
$f \in C[0, \frac{A}{2}] \Rightarrow f_{\text{even}} \in C(\mathbb{R})$
A-periodic



(2) Odd extension

$$f_{\text{odd}}(x) = \begin{cases} f(x) & \text{if } x \in (0, \frac{A}{2}] \\ 0 & \text{if } x = 0 \\ -f(-x) & \text{if } x \in [-\frac{A}{2}, 0) \end{cases}$$

$f \in C[0, \frac{A}{2}] \Rightarrow f_{\text{odd}}$:
discontinuous



Then their Fourier series expansions get simpler.

But before computing them, let's review the relationship between the Fourier coefficients

$\{C_k\}$ w.r.t. the ONB $\{\frac{1}{\sqrt{A}} e^{2\pi i k x/A}\}$ and
 $\{a_k, b_k\}$ w.r.t. the ONB $\{\frac{1}{\sqrt{A}}\} \cup \{\frac{\sqrt{2}}{\sqrt{A}} \cos(\frac{2\pi k x}{A})\} \cup \{\frac{\sqrt{2}}{\sqrt{A}} \sin(\frac{2\pi k x}{A})\}$

Let $g(x)$ be an A -periodic L^2 fcn.

$$g(x) \sim \frac{1}{\sqrt{A}} \sum_{k \in \mathbb{Z}} C_k e^{2\pi i k x/A}, \quad C_k = \frac{1}{\sqrt{A}} \int_{-A/2}^{A/2} g(x) e^{-2\pi i k x/A} dx$$

$$\text{Then } \frac{1}{\sqrt{A}} \sum_{k \in \mathbb{Z}} C_k e^{2\pi i k x/A}$$

$$= \frac{1}{\sqrt{A}} \left[C_0 + \sum_{k=1}^{\infty} (C_k + C_{-k}) \cos\left(\frac{2\pi k x}{A}\right) + i(C_k - C_{-k}) \sin\left(\frac{2\pi k x}{A}\right) \right]$$

$$= \frac{C_0}{\sqrt{A}} + \sum_1^{\infty} \left[\underbrace{\frac{C_k + C_{-k}}{\sqrt{2}}}_{a_k} \underbrace{\sqrt{\frac{2}{A}} \cos\left(\frac{2\pi k x}{A}\right)}_{\text{red wavy}} + \underbrace{\frac{C_k - C_{-k}}{\sqrt{2}} i}_{b_k} \underbrace{\sqrt{\frac{2}{A}} \sin\left(\frac{2\pi k x}{A}\right)}_{\text{red wavy}} \right]$$

If we want, we can write a_0 instead of C_0 .

Remark: In many books, the Fourier series is often defined in $[-\pi, \pi)$ and written as

$$\sum_{k=-\infty}^{\infty} C_k e^{ikx} = \frac{a_0}{2} + \sum_{k=1}^{\infty} (a_k \cos kx + b_k \sin kx),$$

where $a_0 = 2C_0$, $a_k = C_k + C_{-k}$, $b_k = i(C_k - C_{-k})$, $k \geq 1$

But $\{e^{ikx}\}$ is an orthogonal basis of $L^2[-\pi, \pi]$, but not normalized. To make it an ONB, one needs the factor $\frac{1}{\sqrt{2\pi}}$.

The same can be said for the orthogonal basis $\{1\} \cup \{\cos kx\} \cup \{\sin kx\}$.

The ONB is $\{\frac{1}{\sqrt{2\pi}}\} \cup \{\frac{1}{\sqrt{\pi}} \cos kx\} \cup \{\frac{1}{\sqrt{\pi}} \sin kx\}$

i.e., $A = 2\pi$.

Compare this notation with mine in this lecture, which is the orthonormalized version:

$$\sum_{k=-\infty}^{\infty} C_k \frac{1}{\sqrt{A}} e^{2\pi i k x / A} = \frac{1}{\sqrt{A}} a_0 + \sqrt{\frac{2}{A}} \sum_{k=1}^{\infty} \left(a_k \cos\left(\frac{2\pi k}{A} x\right) + b_k \sin\left(\frac{2\pi k}{A} x\right) \right)$$

where $a_0 = C_0$, $a_k = \frac{C_k + C_{-k}}{\sqrt{2}}$, $b_k = \frac{C_k - C_{-k}}{\sqrt{2}} i$, $k \geq 1$.

Now, let's go back to the Fourier series expansion of f_{even} & f_{odd} .

$$\begin{aligned} C_k[f_{\text{even}}] &= \frac{1}{\sqrt{A}} \int_{-A/2}^{A/2} f_{\text{even}}(x) e^{-2\pi i k x / A} dx \\ &= \frac{2}{\sqrt{A}} \int_0^{A/2} f(x) \cos\left(\frac{2\pi k x}{A}\right) dx \\ &= C_{-k}[f_{\text{even}}] \quad \text{thanks to the evenness of } \cos \theta. \end{aligned}$$

Recall the relationship:

$$a_0 = C_0, \quad a_k = \frac{C_k + C_{-k}}{\sqrt{2}}, \quad b_k = \frac{C_k - C_{-k}}{\sqrt{2}};$$

Hence, in this case of f_{even} , $a_0 = C_0[f_{\text{even}}]$,
 $a_k = \sqrt{2} C_k[f_{\text{even}}]$, $b_k \equiv 0$, $k \geq 1$.

In other words, f_{even} can be written as the Fourier **cosine** series:

$$f_{\text{even}}(x) \sim \frac{1}{\sqrt{A}} a_0 + \sum_{k=1}^{\infty} a_k \sqrt{\frac{2}{A}} \cos\left(\frac{2\pi k x}{A}\right)$$

$$\text{where } a_0 = \frac{2}{\sqrt{A}} \int_0^{A/2} f(x) dx = C_0[f_{\text{even}}],$$

$$\begin{aligned} a_k &= 2 \sqrt{\frac{2}{A}} \int_0^{A/2} f(x) \cos\left(\frac{2\pi k x}{A}\right) dx \\ &= \sqrt{2} C_k[f_{\text{even}}]. \end{aligned}$$

Similarly, for $f_{\text{odd}}(x)$,

$$C_k[f_{\text{odd}}] = \frac{1}{\sqrt{A}} \int_{-A/2}^{A/2} f_{\text{odd}}(x) e^{-2\pi i k x / A} dx$$

$$= \frac{-2i}{\sqrt{A}} \int_0^{A/2} f(x) \sin\left(\frac{2k\pi x}{A}\right) dx$$

$$= -C_{-k}[f_{\text{odd}}] \text{ due to the oddness of } \sin \theta.$$

$$\Rightarrow a_k = \frac{C_k + C_{-k}}{\sqrt{2}} \equiv 0, \quad k \in \mathbb{N}.$$

$$a_0 = C_0 = \frac{1}{\sqrt{A}} \int_{-A/2}^{A/2} f_{\text{odd}}(x) dx = 0.$$

No
cosines

$$b_k = \frac{C_k - C_{-k}}{\sqrt{2}} i = \sqrt{2} i C_k[f_{\text{odd}}]$$

$$\text{So, } f_{\text{odd}}(x) \sim \sum_{k=1}^{\infty} b_k \sqrt{\frac{2}{A}} \sin\left(\frac{2\pi k x}{A}\right)$$

$$b_k = 2 \sqrt{\frac{2}{A}} \int_0^{A/2} f(x) \sin\left(\frac{2\pi k x}{A}\right) dx.$$

So, if a fcn is given on $[0, \frac{A}{2}]$, say $f \in C[0, \frac{A}{2}]$

$\equiv$ three ways to extend it to a periodic fcn

- $O(1/k)$ (1) Brute force periodization with period $A/2$.
- $O(1/k^2)$ (2) Even extension followed by A -periodization.
- $O(1/k)$ (3) Odd " " "

(2) is the best among these 3 in terms of the decay of the Fourier coeff's.

However, $\equiv$ an even better way!

★ The Lanczos Method (1938)

Suppose $f \in C^{2m}[0, 1]$, but $f(0) \neq f(1)$ no match at $x=0, 1$.
also assume $f^{(2m+1)} \in BV[0, 1]$. $m=1, 2, \dots$

Lanczos's idea:

decompose $f(x) = u(x) + v(x)$

where $u(x)$ = a polynomial of degree $2m-1$.

s.t.
$$\begin{cases} u^{(2k)}(0) = f^{(2k)}(0) \\ u^{(2k)}(1) = f^{(2k)}(1) \end{cases}, k=0, 1, \dots, m-1.$$

Then, consider the **odd** extension of v .

$\Rightarrow v \in C^{2m-1}(\mathbb{R})$, $v^{(k)}(0) = v^{(k)}(1) = 0$
 $k=0, 1, 2, \dots, 2m-1.$

and the Fourier sine coefficients of $v(x)$ (with period 2) decay as

$$b_k = O(|k|^{-2m-1})$$

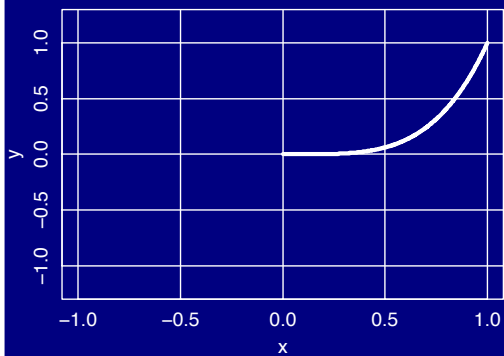
e.g., $m=1$ gives us $b_k = O(1/k^3)$.
and $u(x)$ is a straight line connecting $(0, f(0))$ & $(1, f(1))$.

Remarks:

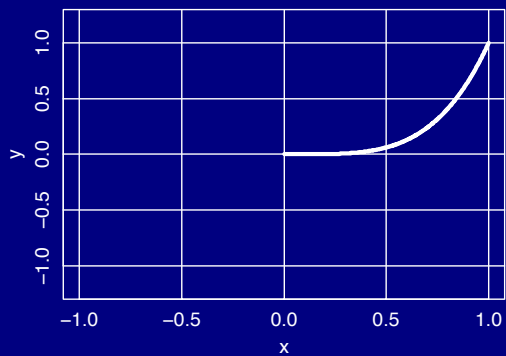
(1) $f = u + v =$ (an algebraic poly) + (a trig. poly).
can avoid both the **Runge** & **Gibbs** phenomena!

(2) NS-J.F. Remy (2006) generalized this to $\mathbb{R}^d$.

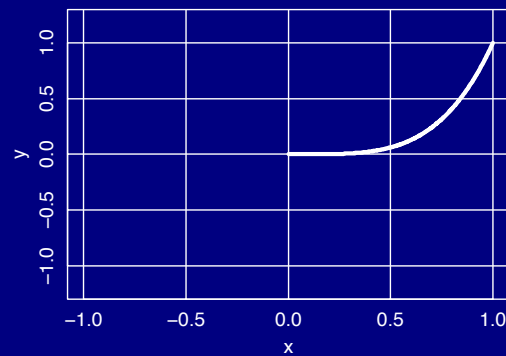
Original Signal Supported on [0,1]



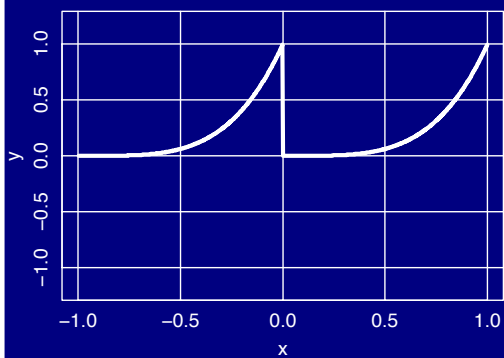
Original Signal Supported on [0,1]



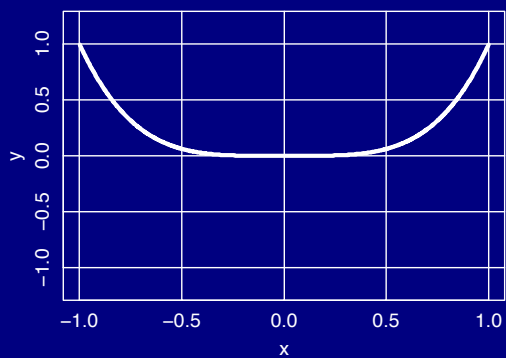
Original Signal Supported on [0,1]



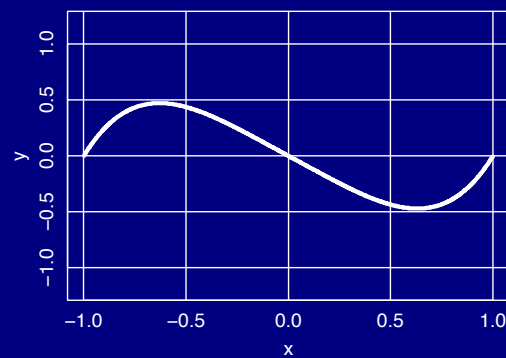
After Periodization



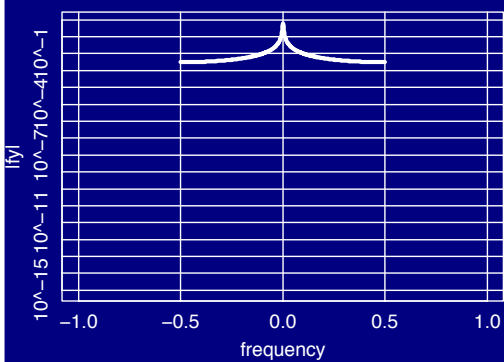
After Even Reflection



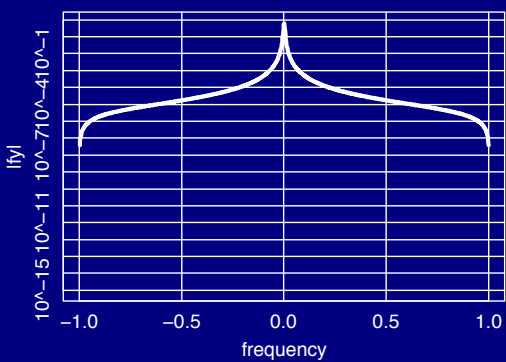
After Lin Removal+Odd Reflect



DFT Coefficients



DCT Coefficients



LLST Coefficients

