

SOLUTIONS: PROBLEM SET 7  
Math 207B, Winter 2012

1. (a) Explain why the total birth rate  $B(t)$  of a population with constant reproductive rate  $\lambda$  per individual and exponential survival rate  $e^{-\beta t}$  over time  $t$  satisfies the renewal equation

$$B(t) = N_0\lambda e^{-\beta t} + \lambda \int_0^t e^{-\beta(t-s)} B(s) ds$$

where  $N_0$  is the total initial population at  $t = 0$ . What are the dimensions of  $N_0$ ,  $\lambda$ ,  $\beta$ , and  $B(t)$ ?

(b) Solve this integral equation and discuss the behavior of  $B(t)$  as  $t \rightarrow \infty$ . Does your answer make sense?

HINT. One approach is to show that  $B(t)$  satisfies the IVP

$$\dot{B} = (\lambda - \beta)B, \quad B(0) = N_0\lambda.$$

**Solution**

- (a) After time  $t$ , there are  $N_0 e^{-\beta t}$  survivors from the initial population, and they contribute  $\lambda \cdot N_0 e^{-\beta t}$  to the birthrate. For  $0 < s < t$ , the population born between times  $s$  and  $s + ds$  is  $B(s) ds$ , and of these  $e^{-\beta(t-s)} B(s) ds$  survive to time  $t$ , so they contribute  $\lambda e^{-\beta(t-s)} B(s) ds$  to the birth rate. Integrating this over  $0 < s < t$  and adding the contribution from the initial population, we get the total birth rate  $B(t)$ .
- Let  $P$  denote the dimension of population and  $T$  the dimension of time. Then

$$[N_0] = P, \quad [\lambda] = \frac{1}{T}, \quad [\beta] = \frac{1}{T}, \quad [B] = \frac{P}{T}.$$

- (b) Differentiating the integral equation with respect to  $t$ , we get

$$\dot{B}(t) = -N_0\lambda\beta e^{-\beta t} - \lambda\beta \int_0^t e^{-\beta(t-s)} B(s) ds + \lambda B(t).$$

Using the integral equation to eliminate the integral from this equation, we find that

$$\dot{B} = (\lambda - \beta)B.$$

Setting  $t = 0$  in the integral equation, we get

$$B(0) = N_0\lambda.$$

The solution of this IVP is

$$B(t) = N_0\lambda e^{(\lambda-\beta)t}.$$

- If  $\lambda > \beta$ , meaning that the birth rate per individual exceeds the death rate, then the total birth rate (and the total population) grows exponentially in time. If  $\lambda < \beta$ , meaning that the death rate per individual exceeds the birth rate, then the total birth rate (and the total population) decays exponentially. If  $\lambda = \beta$ , then the total birth rate (and the total population) remains constant.

2. (a) Consider the following Fredholm equation of the second kind

$$u(x) - \lambda \int_0^1 xyu(y) dy = f(x), \quad 0 \leq x \leq 1$$

where  $\lambda \in \mathbb{C}$  is a constant and  $f$  is a given (continuous) function. If  $\lambda \neq 3$ , show that this equation has a unique (continuous) solution for  $u(x)$  and find the solution. If  $\lambda = 3$ , determine for what functions  $f$  a solution exists and find the solutions in that case.

(b) For what functions  $f$  is the following Fredholm equation of the first kind

$$\int_0^1 xyu(y) dy = f(x), \quad 0 \leq x \leq 1$$

solvable for  $u(x)$ ? Describe the solutions in that case.

HINT. Note that these Fredholm equations are degenerate.

### Solution

- (a) The integral equation implies that

$$u(x) = f(x) + cx \tag{1}$$

where the constant  $c$  is given by

$$c = \lambda \int_0^1 yu(y) dy.$$

Using the expression (1) for  $u$  in this equation for  $c$ , we get

$$\begin{aligned} c &= \lambda \int_0^1 yf(y) dy + \lambda \int_0^1 y \cdot cy dy \\ &= \lambda \int_0^1 yf(y) dy + \frac{1}{3}\lambda c, \end{aligned}$$

or

$$(3 - \lambda)c = 3\lambda \int_0^1 yf(y) dy.$$

- If  $\lambda \neq 3$ , then we get

$$c = \frac{3\lambda}{3 - \lambda} \int_0^1 yf(y) dy$$

and (1) with this value of  $c$  is the unique solution of the integral equation for any continuous function  $f$ .

- If  $\lambda = 3$ , then the integral equation is solvable if and only if

$$\int_0^1 yf(y) dy = 0.$$

In that case, the solution is (1) where  $c$  is an arbitrary constant.

- Note that the integral operator

$$(Ku)(x) = \int_0^1 xyu(y) dy$$

has eigenvalue  $\mu = 1/3$  and eigenfunction  $u(x) = x$  with  $Ku = \mu u$ , so  $cx$  is an arbitrary solution of the homogeneous equation for  $\lambda = 3$ .

- (c) Since

$$\int_0^1 xyu(y) dy = cx, \quad c = \int_0^1 yu(y) dy$$

the equation is only solvable if

$$f(x) = cx$$

for some constant  $c$ . In that case, one solution is the constant function

$$u(x) = 2c$$

The general solution is

$$u(x) = 2c + v(x)$$

where  $v(x)$  is any function orthogonal to  $x$ , meaning that

$$\int_0^1 xv(x) dx = 0.$$

- Note that the range of  $K$  is one-dimensional, and the nullspace of  $K$ , consisting of eigenfunctions with eigenvalue  $\mu = 0$ , is infinite-dimensional.

3. Show that the following IVP for a second-order scalar ODE for  $u(t)$

$$\begin{aligned}\ddot{u}(t) &= f(t, u(t)), \\ u(0) &= u_0, \quad \dot{u}(0) = v_0\end{aligned}$$

is equivalent to the Volterra integral equation

$$u(t) = \int_0^t (t-s)f(s, u(s)) ds + u_0 + v_0t.$$

### Solution

- Integrating the ODE once, we get

$$\dot{u}(t) = v_0 + \int_0^t f(s, u(s)) ds.$$

Integrating again, we get

$$u(t) = u_0 + v_0t + \int_0^t \left[ \int_0^r f(s, u(s)) ds \right] dr.$$

Integrating by parts in the  $r$ -integral, we obtain

$$\begin{aligned}\int_0^t \left[ \int_0^r f(s, u(s)) ds \right] dr &= \int_0^t 1 \cdot \left[ \int_0^r f(s, u(s)) ds \right] dr \\ &= \left[ r \int_0^r f(s, u(s)) ds \right]_0^t - \int_0^t r f(r, u(r)) dr \\ &= t \int_0^t f(s, u(s)) ds - \int_0^t s f(s, u(s)) ds \\ &= \int_0^t (t-s)f(s, u(s)) ds,\end{aligned}$$

which shows that a solution of the IVP satisfies the integral equation.

- Conversely, if  $u(t)$  satisfies the integral equation, then setting  $t = 0$  in the equation we get  $u(0) = u_0$ . Similarly, differentiating the integral equation once and setting  $t = 0$ , we get  $\dot{u}(0) = v_0$ . Differentiating once again, we get  $\ddot{u} = f(t, u)$ , so a solution of the integral equation satisfies the IVP.
- Note that the integral equation incorporates both the ODE and the initial conditions.

4. Consider the following BVP for  $u(x)$  in  $0 < x < 1$ :

$$\begin{aligned} -u'' &= k^2 [1 + \epsilon q(x)] u + f(x), \\ u(0) &= 0, \quad u(1) = 0. \end{aligned} \tag{2}$$

Here  $k > 0$  is a constant,  $\epsilon$  is a small parameter, and  $f(x)$ ,  $q(x)$  are given (continuous) functions. Assume that  $k \neq n\pi$  for any integer  $n \in \mathbb{N}$ , so that  $k^2$  is not an eigenvalue of  $-d^2/dx^2$  with Dirichlet BCs.

(a) Find the Green's function  $G(x, \xi)$  for (2) with  $\epsilon = 0$ , which satisfies

$$\begin{aligned} -\frac{d^2 G}{dx^2} &= k^2 G + \delta(x - \xi) \quad \text{in } 0 < x < 1, \\ G(0, \xi) &= 0, \quad G(1, \xi) = 0. \end{aligned} \tag{3}$$

(b) Use the Green's function from (a) to reformulate (2) as a Fredholm integral equation for  $u(x)$  of the form

$$u(x) = \epsilon \int_0^1 K(x, \xi) u(\xi) d\xi + g(x). \tag{4}$$

(c) Write out the first few terms in the Neumann series (or Born approximation) for  $u$  as integrals involving  $g(x)$ ,  $q(x)$ , and  $G(x, \xi)$ .

### Solution

- (a) Solutions of the homogeneous equation

$$-\frac{d^2 u}{dx^2} - k^2 u = 0$$

that vanish at  $x = 0$  and  $x = 1$  are

$$u_1(x) = \sin kx, \quad u_2(x) = \sin(k(1 - x))$$

respectively. The Wronskian of these solutions is

$$\begin{aligned} u_1 u_2' - u_2 u_1' &= -k \sin kx \cos(k(1 - x)) - k \cos kx \sin(k(1 - x)) \\ &= -k \sin k, \end{aligned}$$

which is nonzero if  $k \neq n\pi$ .

- The Green's function in (3), whose  $x$ -derivative jumps by  $-1$  across  $x = \xi$ , is therefore

$$G(x, \xi) = \begin{cases} \sin kx \sin (k(1 - \xi)) / k \sin k & \text{if } 0 \leq x \leq \xi, \\ \sin k\xi \sin (k(1 - x)) / k \sin k & \text{if } \xi \leq x \leq 1. \end{cases}$$

or

$$G(x, \xi) = \frac{\sin(kx_{<}) \sin(k(1 - x_{>}))}{k \sin k}.$$

(b) We write (2) as

$$\begin{aligned} -u'' - k^2u &= \epsilon q(x)u + f(x), \\ u(0) &= 0, \quad u(1) = 0. \end{aligned}$$

It follows from the Green's function representation that

$$\begin{aligned} u(x) &= \int_0^1 G(x, \xi) [\epsilon q(\xi)u(\xi) + f(\xi)] d\xi \\ &= \epsilon \int_0^1 G(x, \xi)q(\xi)u(\xi) d\xi + \int_0^1 G(x, \xi)f(\xi) d\xi, \end{aligned}$$

which gives (4) with

$$K(x, \xi) = G(x, \xi)q(\xi), \quad g(x) = \int_0^1 G(x, \xi)f(\xi) d\xi.$$

- (c) The Neumann series expansion for

$$u = g + \epsilon K u$$

is

$$u = g + \epsilon K g + \epsilon^2 K^2 g + \epsilon^3 K^3 g + \dots$$

Explicitly, we have

$$\begin{aligned} K g(x) &= \int_0^1 G(x, \xi)q(\xi)g(\xi) d\xi, \\ K^2 g(x) &= \int_0^1 \int_0^1 G(x, \xi_1)G(x, \xi_2)q(\xi_1)q(\xi_2)g(\xi_2) d\xi_1 d\xi_2, \\ K^3 g(x) &= \int_0^1 \int_0^1 \int_0^1 G(x, \xi_1)G(x, \xi_2)G(x, \xi_3)q(\xi_1)q(\xi_2)q(\xi_3)g(\xi_3) d\xi_1 d\xi_2 d\xi_3. \end{aligned}$$

- These terms correspond physically to single, double, and triple scattering corrections due to the perturbation  $\epsilon q(x)$  in the coefficient.