

## SECTION-10

### Entropy (According to Lax)

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# Math-280: A Mathematical Introduction to Shock Waves

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Entropy:

• compressible Euler Eqn's

$$\dot{S}_t + \operatorname{div} S u = 0 \quad (\text{MA})$$

$$(\rho u^i)_t + \operatorname{div} (\rho u^i u + p e^i) = 0 \quad (\text{Mo})$$

$$\dot{E}_t + \operatorname{div} (E + p) u = 0 \quad (\text{En})$$

$$\dot{S}'_t + \operatorname{div} S' u = 0 \quad (\dot{S}')$$

$$S = SS = \frac{\text{entropy}}{\text{vol}}, \quad E \equiv \frac{1}{2} \rho u^2 + \rho e \equiv \frac{\text{energy}}{\text{vol}}$$

$$S = \text{specific entropy}, \quad e = \text{specific energy}, \quad V = \frac{1}{\rho} \equiv \frac{\text{spec vol}}{\text{vol}}$$

• 2nd Law Thermo:  $de = Tds - pdv$

• Polytropic ( $\gamma$ -law) gas:  $e = C_v \rho^{\gamma-1} \exp \frac{s}{C_v}$

Solve for  $s$ :  $\frac{e}{C_v \rho^{\gamma-1}} = \exp \frac{s}{C_v}$

$$s = C_v \ln \left( \frac{e}{C_v \rho^{\gamma-1}} \right)$$

$$S = SS = S C_v \ln \left( \frac{e}{C_v \rho^{\gamma-1}} \right)$$

①

$$\text{Now } e = \frac{E}{\rho} - \frac{1}{2} u^2 \Rightarrow$$

$$S = S(V) = S C_v \ln \left\{ \frac{E}{C_v \rho^\gamma} - \frac{u^2}{C_v \rho^{\gamma-1}} \right\}$$

$$V = (\rho, m^2, E), \quad m^2 = \rho u^2$$

$$u^2 = \frac{m^2}{\rho^2}$$

$$\Rightarrow S = S(V) = S(\rho, m, E) = C_v S \ln \left\{ \frac{E}{C_v \rho^\gamma} - \frac{m^2}{C_v \rho^{\gamma+1}} \right\}$$

Defn:  $S$  is convex if its Hessian is

pos definite:  $H = \frac{\partial^2 S}{\partial v_i \partial v_j}$

satisfies

$$v^T H v \geq 0 \quad \forall v \in \mathbb{R}^3$$

Thm:  $-S$  is convex up (HW)

②

• "Entropy is constant" on smooth soln's ③

$$\frac{d}{dt} \int_{\Omega(t)} S d^3x = \int_S_t + \operatorname{div}(S u) d^3x = 0$$

For soln's  $u(x, t)$  with shock-waves,  
we want entropy to increase, so ask ④

$$\frac{d}{dt} \int_{\Omega(t)} S(u) d^3x \geq 0$$

$$\Leftrightarrow \int_{\Omega(t)} S_t + \operatorname{div}(S u) d^3x \geq 0 \quad \forall \Omega(t)$$

↑  
smooth  
solution

$$\Leftrightarrow S_t + \operatorname{div}(S u) \geq 0. \quad (*)$$

We look for weak formulation of (\*):

"mult by test fn and int. by parts"  $\phi \in C_0^\infty(\mathbb{R}^3 \times \mathbb{R}^4)$

Choose  $\phi \geq 0$ . Then (\*)  $\Leftrightarrow$

$$S_t \phi + \operatorname{div}(S u) \phi \geq 0 \quad \forall \phi \geq 0$$

$$\Leftrightarrow \iint_{\mathbb{R}^3 \times \mathbb{R}^4} S_t \phi + \operatorname{div}(S u) \phi \geq 0 \quad \forall \phi \geq 0$$

$(\phi \text{ compact supp})$

$$-\iint_{\mathbb{R}^3 \times \mathbb{R}^+} S \phi_t + S u \cdot \nabla \phi$$

$$+\iiint_{\mathbb{R}^3} (S \phi)_t + \operatorname{div}(S u \phi) d^3x \geq 0$$

0 by div. thm

$$\iint_{\mathbb{R}^3} S \phi_t + S u \cdot \nabla \phi d^3x \leq 0$$

$$\Leftrightarrow \boxed{\iint_{\mathbb{R}^3} S \phi_t + S u \cdot \nabla \phi d^3x \leq 0} \quad \left. \begin{array}{l} \\ (*) \end{array} \right\}$$

$$\forall \phi \in C_0^\infty(\mathbb{R}^3 \times \mathbb{R}^+) \quad \phi \geq 0.$$

Defn: Entropy increases on a weak solution  $u(x, t)$  if (\*) holds. We say that (\*) expresses that  $S_t + \operatorname{div}(S v) \geq 0$  in the weak sense, or in the sense of the theory of distributions.

(5)

- It turns out that  $-S$  is a convex <sup>up</sup> fn of  $(\rho, m, E) \in \mathcal{U}$ .

Defn: Let  $u_t + \operatorname{div} f = 0$  be an arbitrary system of conservation laws  $x \in \mathbb{R}^m$ . We say the system possesses a <sup>convex</sup> entropy if  $\exists$  a fn  $V(u)$ , convex up, such that, on smooth solutions  $u(x, t)$  of  $u_t + \operatorname{div} f = 0$ ,

$$V(u)_t + \operatorname{div} F(u) = 0.$$

Defn: we say entropy increases on a weak soln of  $u_t + \operatorname{div} f = 0$  if

$$(v) \quad U_t + \operatorname{div} F \leq 0 \quad \iint_{\mathbb{R}^3 \times \mathbb{R}^+} U \phi_t + F \cdot \nabla \phi \geq 0$$

in the weak sense.  $\Leftrightarrow \forall \phi \geq 0, \phi \in C_0^\infty$

(i.e.  $U = -S \Rightarrow S_t + \operatorname{div} S u \geq 0$ )

Note: Let  $U$  be a convex entropy for system of cons. laws  $U_t + \text{div}(u) = 0$ , and assume that on weak soln's

$$U_t + \text{div } F \leq 0. \quad (*)$$

Let  $u(x, t)$  be a <sup>weak</sup> soln with compact support in  $x \in \mathbb{R}^m$  at each  $0 \leq t \leq T$ .

claim:  $\int_{\mathbb{R}^m} U(x, T) dx - \int_{\mathbb{R}^m} U(x, 0) dx \leq 0$

(total entropy decreases)

Pf. Choose  $\phi = 1$  for  $0 < t < T$ ,  $\phi = 0$  for  $t < 0$  and  $t > T$ . Then (\*) implies

$$\iint U \phi_t + F \cdot \nabla \phi dx dt \geq 0$$

$\begin{matrix} xt & \uparrow & 0 \\ \delta f_n @ t=0 & & -\delta @ t=T \end{matrix}$

$$-\int_{\mathbb{R}^m} U(x, T) dx + \int_{\mathbb{R}^m} U(x, 0) dx \geq 0 \quad \checkmark$$

Expect: weak solns of a system of cons. laws with convex entropy are unique for given initial data if  $U$  holds. (8)

No proof for  $n \geq 2$ . ( $n = \#$  of equations)  
even in 1-space dimension. (Mostly resolved by Bressan)

$(U)$  can serve as a test that numerical schemes converge to the correct physical solution.

- Consider a system of cons. laws with ⑨  
a "convex entropy" in  $m$  dimensions  $(x^1, \dots, x^m)$

$$u_t + \operatorname{div} f = 0 \quad (\text{CL})$$

$$\nabla u_t + \operatorname{div} F(u) = 0 \quad (\text{Ent})$$

↑                      ↑  
 entropy              entropy  
 flux

$\nabla$  convex fn of  $U$ .

Then diff entropy:

$$0 = \nabla U \cdot u_t + \nabla F_1(u) u_{x_1} + \dots + \nabla F_m(u) u_{x_m} \quad (\text{A})$$

Mult (CL) by  $\nabla U$  & compare:

$$0 = \nabla U \cdot u_t + \nabla U \cdot df_1 u_{x_1} + \dots + \nabla U \cdot df_m u_{x_m} \quad (\text{B})$$

Equating (A), (B) we must have

$$\nabla U \cdot df_i = \nabla F_i \quad \forall i=1, \dots, m$$

Defn: A pair of functions  $(U(u), F(u))$  is ⑩  
an entropy-entropy flux pair for (CL) if

$$\boxed{\nabla F_i = \nabla U \cdot df_i}$$

In this case

$$U_t + \operatorname{div} F = 0$$

holds on smooth soln's of (CL).

In 1-D this is:

$$u_t + f(u)_x = 0$$

$$\nabla U_t + F(u)_x = 0$$

$$\nabla F = \nabla U \cdot df$$

(\*)

Theorem ① Assume that  $(U, F)$  is a convex<sup>⑪</sup> entropy-entropy flux pair for a system of IONS.

laws

(CL)

so that

$$u_t + \operatorname{div} f(u) = 0$$

$$U(u)_t + \operatorname{div} F(u) = 0$$

$$u = (u_1, \dots, u_n) \in \mathbb{R}^n$$

$$f = (f_1, \dots, f_m)$$

$$f_i(u) \in \mathbb{R}^n$$

$$x = (x_1, \dots, x_n) \in \mathbb{R}^m$$

holds on smooth solutions,

$$\nabla F_i = \nabla U \cdot d f_i$$

Then if  $u^\varepsilon(x, t)$  is a smooth soln of

$$u_t + \operatorname{div} f(u) = \varepsilon \Delta u$$

$$u_i^{\varepsilon}_t + \operatorname{div} f_i(u) = \varepsilon \Delta u^i \quad i = 1, \dots, n$$

coupled scalar eqns

such that  $u^\varepsilon(x, t) \xrightarrow{\varepsilon \rightarrow 0} u(x, t)$  boundedly & a.e. at each time  $t$ .  
 Then  $u(x, t)$  is a weak soln of (CL) that satisfies  $U_t + F_x \leq 0$  in weak sense. (CL)

Proof: We do case  $m=1$ ,  $x \in \mathbb{R}$ .

⑫

Homework: Prove this in case  $m > 1$ .

- First assume  $\phi(x, t)$  smooth test fn with compact support  $\Rightarrow$

$$\begin{aligned} & \iint_{xt} u_t^\varepsilon \phi + f(u^\varepsilon)_x \phi - \varepsilon u_{xx}^\varepsilon \phi \, dx dt \\ &= - \iint_{xt} u^\varepsilon \phi_t + f(u^\varepsilon) \phi_x - \int_{-\infty}^{\infty} u^\varepsilon(x, 0) \phi(x, 0) \, dx \\ & \quad - \varepsilon \iint_{xt} u^\varepsilon \phi_{xx} \, dx dt = 0 \end{aligned}$$

Since  $u^\varepsilon \rightarrow u$  boundedly at each time  $t$ , taking  $\varepsilon \rightarrow 0$  gives

$$\iint_{xt} u \phi_t + f(u) \phi_x = \int_{-\infty}^{\infty} u(x, 0) \phi(x, 0) \, dx$$

$\Rightarrow u(x, t)$  is a weak soln of CL.

For (1), multiply

$$u_t^\varepsilon + f(u^\varepsilon)_x = \varepsilon u_{xx}^\varepsilon.$$

by  $\nabla U(u^\varepsilon)$ , and use  $\nabla F(u^\varepsilon) = \nabla U^2 \cdot df(u^\varepsilon)$

$$\nabla U \cdot u_t^\varepsilon + \nabla U \cdot df u_x^\varepsilon = \varepsilon \nabla U \cdot u_{xx}^\varepsilon$$

$$U(u^\varepsilon)_t + F(u^\varepsilon)_x = \varepsilon \nabla U \cdot u_{xx}^\varepsilon.$$

- consider:  $U_x = \nabla U \cdot u_x$  - Hessian of  $U$   
 $U_{xx} = (\nabla U \cdot u_x)_x = \underbrace{H\bar{U}[u_x, u_x]}_{+ \nabla U u_{xx}}$

HJ per def  $\Rightarrow$

$$U_{xx} \geq \nabla U \cdot u_{xx}$$

$$\Rightarrow U(u^\varepsilon)_t + F(u^\varepsilon)_x \leq \varepsilon U_{xx}$$

(13)

(14)

∴ assuming  $u^\varepsilon$  smooth and  $u^\varepsilon \rightarrow u$  boundedly b pme,  $\phi \geq 0$   $\phi \in C_0^\infty(\mathbb{R} \times \mathbb{R}^+)$

$$\iint_{\mathbb{R} \times \mathbb{R}^+} \{U(u^\varepsilon)_t + F(u^\varepsilon)_x - \varepsilon U_{xx}\} \phi dx dt \leq 0$$

$\mathbb{R} \times \mathbb{R}^+$

$$\Rightarrow - \iint_{\mathbb{R} \times \mathbb{R}^+} \{U(u^\varepsilon) \phi_t + F(u^\varepsilon) \phi_x + \varepsilon U \phi_{xx}\} dx dt \leq 0$$

since  $u^\varepsilon \rightarrow u$  strongly b no deriv's on  $U, F$

$$\xrightarrow{\varepsilon \rightarrow 0} - \iint_{\mathbb{R} \times \mathbb{R}^+} \{U(u^0) \phi_t + F(u^0) \phi_x\} dx dt \leq 0$$

$\Rightarrow u^0$  satisfies entropy inequality

$$U_t + F_x \leq 0$$

in weak sense ✓

Cor: Let  $u_t + f(u)_x = 0$  be a system of PDEs.  
laws that has a convex entropy

$$U(u)_t + F(u)_x = 0$$

on smooth soln's. Assume  $u_s(x,t) = [u_L, u_R]$  is a shock that has structure. Then

$$U(u_s)_t + F(u_s)_x \leq 0 \quad (*)$$

in weak sense.

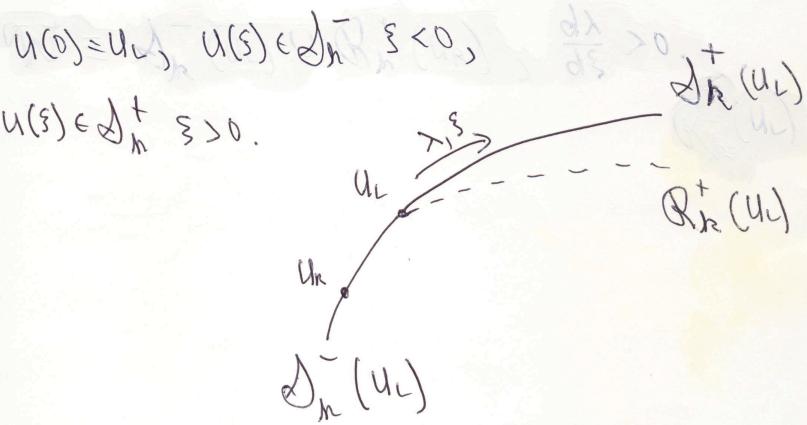
P.F.  $u_\epsilon(\frac{x-st}{\epsilon})$  solves  $u_t + f(u)_x = \epsilon u_{xx}$

&  $u_\epsilon \rightarrow u_s$  bddly & a.e at each time  $\Rightarrow$

$\star$  holds ✓

(15)

• Assume  $u_t + f(u)_x = 0$  is strictly hyperbolic  
so df has e-pairs  $(\lambda_i, R_i)$ ,  $\lambda_1 < \dots < \lambda_n$   
& assume  $(U, F)$  is a convex entropy-entropy  
flux pair. Assume that  $(\lambda, R) \equiv (\lambda_n, R_n)$   
is a genuinely nonlinear family,  $\forall \lambda \cdot R > 0$ ,  
& let  $u(\xi)$  be a parametrization of  $\mathcal{D}_R(u)$



Let  $u_R \in \mathcal{D}_R(u_L)$  so  $s[u] = [f]$ ,

$$u_R = u(\xi) \text{ for } \xi < 0.$$

(16)

Let  $u_s(x,t) = [u_L, u_R]$  denote the shock wave <sup>(17)</sup>  
 soln  $\frac{u_L}{u_R} \frac{dx}{dt} = s \quad s[u] = [f]$

Defn: we say the shock wave  $u_s$  satisfies  
 the entropy condition (EC) if

$$U(u_s)_t + F(u_s)_x < 0$$

in the weak sense.

Theorem <sup>(1)</sup>: Assume  $\nabla \lambda_k \cdot R_k > 0$ , and assume  
 $U$  is strictly convex in direction  $R_k$ , so that

$$\nabla^2 U(R_k, R_h) > 0.$$

Then for  $\xi$  suff small, a  $k$ -shock  $u_s$  satisfies (EC)  
 iff it satisfies the Lax char cond iff  $u_R = u(\xi) \in \bar{\mathcal{J}_h}$ ,  $\xi < 0$ .

(2) Assume  $\nabla \lambda_h \cdot R_h \equiv 0$ . Then we must have  $U$   
 const along  $\mathcal{J}_k = R_h$  &  $\nabla^2 U(R_k, R_h) \equiv 0$ .

Lemma ①: A  $k$ -shock satisfies (EC) iff <sup>(18)</sup>  
 $s[U] > [F]$   $\star$

$$s(U(u_R) - U(u_L)) > (F(u_R) - F(u_L))$$

(Homework)

Lemma ②: Assume  $\nabla \lambda_k \cdot R_k > 0$  &  $\nabla^2 U(R_k, R_h) > 0$ .  
 Then for  $\xi$  suff small, a  $k$ -shock  $u_s$  satisfies  
 $\star$  iff  $u_R = u(\xi) \in \bar{\mathcal{J}_h}$ ,  $\xi < 0$ .

Proof of Lemma 2:

$$\text{1st: } [v] = v(u(\xi)) - v(u)$$

The pt: until 3rd deriv,  
everything that doesn't have  
 $v(u)$  on it cancels out!

(19)

$$(s[v] - [f])^{\circ} \stackrel{(19)}{=} \dot{s}[v] + s\dot{v} - \dot{f}$$

But  $v(\xi)$ ,  $s(\xi)$  satisfy (R+H) identically:

$$0 = (s[u] - [f])^{\circ} = \dot{s}[u] + s\dot{u} - \dot{f}$$

$$\dot{f} = \nabla f \cdot \dot{u}, \quad \nabla v \cdot df = \nabla f$$

$$\Rightarrow \dot{f} = \nabla v \cdot \dot{f} = \nabla v \{ \dot{s}[u] + s\dot{u} \}$$

$$\Rightarrow (1) = \dot{s}[v] + \cancel{\dot{s}\dot{v}} - \nabla v \dot{s}[u] - \cancel{s\dot{v}}$$

$$= \dot{s}[v] - \dot{s}\nabla v[u]. \quad (=0 \text{ at } \xi=0)$$

$$(s[v] - [f])^{\circ} = \dot{s}[v] + \cancel{\dot{s}\dot{v}} - \dot{s}\nabla v[u] - \dot{s}\nabla^2 v \cdot \dot{u}[u] \\ - \cancel{\dot{s}\nabla v \cdot \dot{u}}$$

$$= \dot{s}[v] - \dot{s}\nabla v[u] - \dot{s}\nabla^2 v \dot{u}[u] \quad \overbrace{\dot{v}}$$

(=0 when  $\xi=0$ )

$$(s[v] - [f])^{\circ\circ} = \ddot{s}\dot{v} - \dot{s}\nabla v \cdot \dot{u} - \dot{s}\nabla^2 v[\dot{u}, \dot{u}] \\ + \{ \}_{I}^{II}[v] + \{ \}_{II}^{I}[u]$$

(20)

Set  $\xi=0$ :

$$(s[v] - [f])^{\circ\circ} \Big|_{\xi=0} = \ddot{s}\dot{v} - \cancel{\dot{s}\dot{v}} - \dot{s}\nabla^2 v[\dot{u}, \dot{u}]$$

< 0 (since  $\dot{u} = R_k$ )

$\Rightarrow$  for  $s$  suff small,  $s[v] - [f] \leq 0$  iff

&  $s[v] - [f] < 0$  iff  $s < 0$ .

$s \leq 0$

✓

Note: In case of a linearly degenerate field,  
 $s = \lambda \equiv \text{const fn of } \xi \Rightarrow (s[v] - [f])^{\circ} \equiv 0$

$\Rightarrow$  (EC) holds on sofn & its time reversal.

$$- \dot{s}[v] - \dot{s}\nabla v[u] \leq 0$$

Note: in 1-d, ( $m=1, x \in \mathbb{R}$ ), condt on  $U, F$   
to be an entropy-entropy flux is

$$\nabla F = \nabla U \cdot df$$

This is  $n$  equations in 2 unknowns  $U, F$   
with indept var's  $u$ .

$\Rightarrow$  Expect: can solve if  $n=2$ , but  $n>2$

$\Rightarrow$  only special systems have entropies

(21)

Thm: (Lax)  $n=2 \Rightarrow \exists$  infinite family of  
entropy-entropy flux pairs  
(Used by DiPerna in method of compensated  
compactness)

(22)

Thm: Lax, Harten-ref Young: an entropy-  
entropy flux pair exists iff the system of  
conservation laws is symmetrizable: if  
 $\exists$  mapping  $u \rightarrow v$  st

the system in  $v$ -variables is symmetric. 23

$$u_t + f_x = 0$$

$$\left( \frac{\partial u}{\partial v} v_t + df \frac{\partial u}{\partial v} v_x \right)_v = 0$$

symmetric      symmetric

### Conclusions:

- ① For shocks, (U) is equivalent to the condition that char. impinge on shock when  $|u_i - u_R| \ll 1$ .
- ② Every solution  $u(x,t)$  that is a limit, boundedly a.e. of  $u_t + f(u)_x = \epsilon u_{xx}$  as  $\epsilon \rightarrow 0$  satisfies (U)
- ③ Solutions generated by Godunov's method satisfy (U).

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