

Elastic Waves in Exterior Domains Part II: Global Existence with a Null Structure

Jason Metcalfe¹ and Becca Thomases²

¹Department of Mathematics, University of California, Berkeley, CA 94720-3840 and ²Courant Institute of Mathematical Sciences, New York University, New York, NY 10012

Correspondence to be sent to: Jason Metcalfe, Department of Mathematics, University of California, Berkeley, CA 94720-3840. e-mail: metcalfe@math.berkeley.edu

In this paper, we prove that solutions to a problem in nonlinear elasticity corresponding to small initial displacements exist globally in the exterior of a nontrapping obstacle. The medium is assumed to be homogeneous, isotropic, and hyperelastic, and the nonlinearity is assumed to satisfy a null condition. The techniques contained herein would allow for more complicated geometries provided that there is a sufficient decay of local energy for the linearized problem.

1 Introduction

In this paper, we shall prove global existence of elastic waves in exterior domains subject to a null condition. The elastic medium is assumed to be homogeneous, isotropic, and hyperelastic, and we only consider small initial displacements. We solve the said equations exterior to a bounded, nontrapping obstacle with smooth boundary, though our techniques would allow for any bounded, smooth obstacle for which there is a sufficiently fast decay of local energy.

For the boundaryless problem, global existence was first shown in [20] under a fairly restrictive null condition. This null condition was later relaxed in [1] and [21]. See,

Received May 3, 2006; Accepted April 16, 2007
Communicated by Sergiu Klainerman

See http://www.oxfordjournals.org/our_journals/imrn/ for proper citation instructions.

© The Author 2007. Published by Oxford University Press. All rights reserved. For permissions, please e-mail: journals.permissions@oxfordjournals.org.

also, [22]. As in studies of the wave equation, the null condition is essential in order to obtain global solutions. Without some assumption on the structure of the nonlinearity, blow up in finite time was shown in [5]. In this case, almost global existence of solutions was shown in [6] and a simplified proof was later given in [11]. In [12], the corresponding almost global existence result was shown outside of a star-shaped obstacle.¹

Several related studies have been conducted concerning multiple speed systems of wave equations in three-dimensional exterior domains. Almost global existence was established in [8, 9] and [17] exterior to star-shaped domains. Using the techniques of [15], these results can be extended to any domain for which there is a sufficiently fast decay of local energy. For nonlinearities satisfying the null condition, global existence has been established in [15, 18] and [13, 14]. Our proof will use many of the techniques established in this context.

We shall utilize the method of commuting vector fields, though the equation of interest is not Lorentz invariant and thus we will be required to use a restricted set of vector fields. Moreover, we will rely upon the arguments of [9] and [15] to establish the necessary higher-order energy estimates. Here, we use elliptic regularity to establish energy estimates involving the generators of translations. Without the star-shapedness assumption, proving energy estimates involving the scaling vector field is more delicate, and we will use a boundary term estimate which is reminiscent of those in [15]. For energy estimates involving the full set of vector fields, a class of weighted mixed-norm estimates, called KSS estimates, will be called upon to handle the relevant boundary terms. These estimates are analogous to those of [8] for the wave equation and have previously been established in [12] for the current setting.

The nontrapping hypothesis is assumed to guarantee the existence of an exponential decay of local energy. See [27]. Similar well-known estimates (see [19]) were widely used in the studies of nonlinear wave equations in exterior domains. Our techniques would, in fact, permit any exterior domain² for which there is a sufficiently fast decay of local energy. In particular, the method of proof would allow for the loss of regularity in the local energy decay which would be necessary if there were trapped rays. In the setting of the wave equation, see [3, 4] for such local energy decays and [15], [13, 14] for proofs of long-time existence in such domains. Moreover, by using the technique of proof contained herein, the almost global existence result of [12] may be extended to any domain for which there is such a decay of local energy.

¹By finite propagation speed and [5], blow up in finite time is still possible.

²By this, we mean the exterior of any bounded obstacle with smooth boundary.

Let us more precisely describe the initial boundary value problem. We fix a bounded, nontrapping³ obstacle $\mathcal{K} \subset \mathbb{R}^3$ with smooth boundary. We may, without loss of generality, assume that $0 \in \mathcal{K} \subset \{|x| < 1\}$, and we will do so throughout.

The linearized equation of elasticity is

$$Lu = \partial_t^2 u - c_2^2 \Delta u - (c_1^2 - c_2^2) \nabla(\nabla \cdot u) = 0, \quad (1.1)$$

where $u = (u^1, u^2, u^3)$ is the displacement vector. The constants c_1, c_2 may be assumed to satisfy

$$c_1^2 - \frac{4}{3}c_2^2 > 0, \quad c_2^2 > 0.$$

The two quantities in the above equations represent the bulk and shear moduli respectively.

The nonlinear problem that results from hyperelastic mediums is

$$(Lu)^I = \partial_l (B_{lmn}^{IJK} \partial_m u^J \partial_n u^K) + \dots^4 \quad (1.2)$$

where the real constants B_{lmn}^{IJK} satisfy the symmetry condition

$$B_{lmn}^{IJK} = B_{mln}^{JKI} = B_{lnm}^{IKJ}. \quad (1.3)$$

The terms not explicitly stated in (1.2) are all of cubic or higher order in $\nabla u, \nabla^2 u$. The null condition that we require assumes that

$$B_{lmn}^{IJK} \xi_I \xi_J \xi_K \xi_l \xi_m \xi_n = 0, \quad \text{whenever } \xi \in S^2. \quad (1.4)$$

Rather than interpreting the equation and null condition further, we refer the reader to [1], [21], and the references therein. The interested reader may also wish to refer to the expository article [22].

As has become common, for convenience, we will truncate the equations at the quadratic level. It is straightforward to modify the compatibility conditions below to permit the general case, and by examining the proof of the main theorem, such a

³Recall that more general domains are possible, but the authors are not aware of the necessary decay of local energy results for such domains.

⁴Here and throughout, we use the summation convention. Latin indices indicate that the implicit summation runs from 1 to 3, while Greek indices are used when the summation shall run from 0 to 3.

truncation is easily seen not to affect the long-time existence as the higher order terms will permit additional applications of the decay estimates (5.14) and as such have better long time behavior. We then study the following initial boundary value problem:

$$\begin{cases} (Lu)^I = \partial_l(B_{lmn}^{IJK} \partial_m u^J \partial_n u^K), & (t, \mathbf{x}) \in \mathbb{R}_+ \times \mathbb{R}^3 \setminus \mathcal{K}, \\ u|_{\partial\mathcal{K}} = 0, \\ u(0, \cdot) = f, \quad \partial_t u(0, \cdot) = g \end{cases} \quad (1.5)$$

for “small” Cauchy data f, g . For convenience in the sequel, we shall set

$$Q^I(\nabla u, \nabla^2 u) = \partial_l(B_{lmn}^{IJK} \partial_m u^J \partial_n u^K).$$

In order to solve (1.5), we must assume that the initial data satisfy certain well-known compatibility conditions. Letting $J_k u = \{\partial_x^\alpha u : 0 \leq |\alpha| \leq k\}$ and noticing that $\partial_t^k u(0, \cdot) = \psi_k(J_k f, J_{k-1} g)$, $0 \leq k \leq m$ for formal H^m solutions, we require that the compatibility functions ψ_k — which depend on Q , $J_k f$, and $J_{k-1} g$ — vanish on $\partial\mathcal{K}$ for $0 \leq k \leq m-1$. Smooth data $(f, g) \in C^\infty$ are said to satisfy the compatibility conditions to infinite order if this holds for all m .

Under these assumptions, we may prove that solutions to (1.5) corresponding to small initial displacements exist globally. This is our main result.

Theorem 1.1. Let \mathcal{K} be a bounded, nontrapping obstacle with smooth boundary as above. Assume that B_{lmn}^{IJK} are real constants satisfying (1.3) and (1.4). Suppose further that the data $(f, g) \in C^\infty(\mathbb{R}^3 \setminus \mathcal{K})$ satisfy the compatibility conditions to infinite order. Then, there are positive constants ε_0 and N so that for all $\varepsilon \leq \varepsilon_0$, if

$$\sum_{|\alpha| \leq N} \|\langle x \rangle^{|\alpha|} \partial_x^\alpha f\|_2 + \sum_{|\alpha| \leq N-1} \|\langle x \rangle^{|\alpha|+1} \partial_x^\alpha g\|_2 \leq \varepsilon, \quad (1.6)$$

then (1.5) has a unique global solution $u \in C^\infty([0, \infty) \times \mathbb{R}^3 \setminus \mathcal{K})$. \square

This paper is organized as follows. In the remainder of this section, we introduce some notations that will be used throughout. In particular, we introduce the vector fields that shall be utilized. The second section is devoted to estimates related to the energy inequality. These include the exponential decay of local energy, the energy inequality and its higher order variants, and weighted mixed-norm KSS estimates. In the third section,

⁵Here, and throughout, $\langle x \rangle = \langle r \rangle = \sqrt{1 + |x|^2}$ denotes the Japanese bracket.

we gather the main pointwise decay estimates that are used to give global existence. As a corollary to these, we obtain the main boundary term estimate that is utilized instead of the star-shapedness assumption which was used, e.g., in [12]. The fourth section is devoted to certain Sobolev-type estimates and estimates related to the null condition. Finally, in the last section, we prove Theorem 1.1.

1.1 Notation

When convenient, we shall set $x_0 = t$, $\partial_0 = \partial_t$. The space-time gradient will be denoted by $u' = \partial u = (\partial_t u, \nabla_x u)$, but we shall reserve the notation ∇ for $\nabla = \nabla_x = (\partial_1, \partial_2, \partial_3)$. The generators of the spatial rotations are denoted

$$\Omega = (\Omega_1, \Omega_2, \Omega_3) = \mathbf{x} \times \nabla,$$

where \times is the usual vector cross product.

When studying elasticity, it is natural to use the generators of the simultaneous rotations

$$\tilde{\Omega}_l = \Omega_l I + U_l,$$

where

$$U_1 = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{bmatrix}, \quad U_2 = \begin{bmatrix} 0 & 0 & -1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{bmatrix}, \quad U_3 = \begin{bmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}.$$

The scaling operator will be denoted

$$S = t\partial_t + r\partial_r - 1.$$

This differs slightly from the usual notion of the scaling vector field, but here it is more convenient to use this modification as it properly preserves the null structure. See [21].

We will set

$$Z = \{\partial, \tilde{\Omega}\}, \quad \Gamma = \{\partial, \tilde{\Omega}, S\}.$$

We differentiate these two sets as it will be necessary to produce estimates which require relatively few occurrences of the scaling vector field S . This is due to the fact that the

coefficients of S can be arbitrarily large in a neighborhood of $\partial\mathcal{K}$, while those of Z are bounded. Similar care with S had to be taken in [9], [15], [13, 14], and [12].

A key property of the vector fields is the commutation properties with L . In particular, we have that

$$[L, \partial] = [L, \tilde{\Omega}] = 0, \quad [L, S] = 2L. \quad (1.7)$$

Moreover, we note that the vector fields preserve the null structure of Q . See [21].⁶

We will use the projections onto the radial and transverse directions, defined as follows:

$$P_1 u = \frac{x}{r} \left\langle \frac{x}{r}, u \right\rangle, \quad P_2 u = [I - P_1]u = -\frac{x}{r} \times \left(\frac{x}{r} \times u \right).$$

We shall use $A \lesssim B$ to denote that there is a positive, unspecified constant C so that $A \leq CB$, and we will use $S_t = [0, t] \times \mathbb{R}^3 \setminus \mathcal{K}$ to denote a time-strip of height t in the exterior domain.

2 Energy and KSS estimates

2.1 Local energy decay

An important tool in previous studies of nonlinear problems in exterior domains is the decay of local energy. It is here that we require the geometric condition on the obstacle. In the current setting, we have the following result of [27]. This is an analog of the classical result of [19] for the wave equation.

Theorem 2.1. Let $\mathcal{K} \subset \{|x| < 1\} \subset \mathbb{R}^3$ be a nontrapping obstacle with smooth boundary. Let u solve

$$\begin{cases} Lu = 0, & (t, x) \in \mathbb{R}_+ \times \mathbb{R}^3 \setminus \mathcal{K}, \\ u|_{\partial\mathcal{K}} = 0, \\ \text{supp } u(0, \cdot), \partial_t u(0, \cdot) \subset \{|x| < 10\}. \end{cases} \quad (2.1)$$

Then

$$\left(\int_{\{x \in \mathbb{R}^3 \setminus \mathcal{K} : |x| < 10\}} |u'(t, x)|^2 dx \right)^{1/2} \lesssim e^{-ct} \|u'(0, \cdot)\|_2, \quad (2.2)$$

□

for some $c > 0$.

⁶See Proposition 3.1.

In what follows, we shall require a higher order version of (2.2). To establish this, we utilize the following version of elliptic regularity, which appeared in [12].

Lemma 2.2. Let $\mathcal{K} \subset \{|x| < 1\} \subset \mathbb{R}^3$ be an obstacle with smooth boundary. Suppose that $u \in C^\infty(\mathbb{R}_+ \times \mathbb{R}^3 \setminus \mathcal{K})$, $u|_{\partial\mathcal{K}} = 0$, and u vanishes for large $|x|$ for each t . Then,

$$\sum_{|\alpha| \leq M} \|S^\nu \partial^\alpha u'(t, \cdot)\|_2 \lesssim \sum_{\substack{j+\mu \leq M+\nu \\ \mu \leq \nu}} \|S^\mu \partial_t^j u'(t, \cdot)\|_2 + \sum_{\substack{|\beta|+\mu \leq M+\nu-1 \\ \mu \leq \nu}} \|S^\mu \partial^\beta Lu(t, \cdot)\|_2 \quad (2.3)$$

and

$$\begin{aligned} \sum_{|\alpha| \leq M} \|S^\nu \partial^\alpha u'(t, \cdot)\|_{L^2(|x| < 4)} &\lesssim \sum_{\substack{j+\mu \leq M+\nu \\ \mu \leq \nu}} \|S^\mu \partial_t^j u'(t, \cdot)\|_{L^2(|x| < 6)} \\ &+ \sum_{\substack{|\beta|+\mu \leq M+\nu-1 \\ \mu \leq \nu}} \|S^\mu \partial^\beta Lu(t, \cdot)\|_{L^2(|x| < 6)} \quad (2.4) \end{aligned}$$

for any M and ν . □

Using (2.4) and the fact that ∂_t preserves the Dirichlet boundary conditions, we can immediately establish the following result which is also from [12].

Lemma 2.3. Let $\mathcal{K} \subset \{|x| < 1\} \subset \mathbb{R}^3$ be a smooth, nontrapping obstacle. Suppose that $u|_{\partial\mathcal{K}} = 0$ and $Lu(t, x) = 0$ for $|x| > 4$ and $t > 0$. Suppose also that $u(t, x) = 0$ for $t \leq 0$. Then if M and ν are fixed and if $c > 0$ is as in (2.2),

$$\begin{aligned} \sum_{\substack{|\alpha|+\mu \leq M+\nu \\ \mu \leq \nu}} \|S^\mu \partial^\alpha u'(t, \cdot)\|_{L^2(|x| < 4)} &\lesssim \sum_{\substack{|\alpha|+\mu \leq M+\nu-1 \\ \mu \leq \nu}} \|S^\mu \partial^\alpha Lu(t, \cdot)\|_2 \\ &+ \int_0^t e^{-(c/2)(t-s)} \sum_{\substack{|\alpha|+\mu \leq M+\nu \\ \mu \leq \nu}} \|S^\mu \partial^\alpha Lu(s, \cdot)\|_2 ds. \quad (2.5) \end{aligned} \quad \square$$

2.2 Energy estimates

In this section, we gather the energy estimates which we shall require. The basic energy estimate is rather standard. In order to establish estimates for $\partial^\alpha u$, we shall use elliptic regularity as above. For energy estimates involving the scaling vector field, we use cutoff techniques from [15] and control the resulting commutator with estimates given

in Section 3.3. Finally, we handle the boundary terms that arise when studying $S^\mu Z^\alpha u$ using the KSS estimates given in the following section.

We begin with the standard energy inequality for the variable coefficient operator

$$(L_\gamma u)^I = (Lu)^I + \gamma^{IJ,jk} \partial_j \partial_k u^J. \quad (2.6)$$

We look at smooth solutions of

$$\begin{cases} L_\gamma u = F, & (t, \mathbf{x}) \in \mathbb{R}_+ \times \mathbb{R}^3 \setminus \mathcal{K}, \\ u|_{\partial \mathcal{K}} = 0, \\ u(0, \cdot) = f, \quad \partial_t u(0, \cdot) = g \end{cases} \quad (2.7)$$

assuming that

$$\gamma^{IJ,jk} = \gamma^{II,kj} \quad (2.8)$$

and

$$\sum_{I,J=1}^3 \sum_{j,k=1}^3 \|\gamma^{IJ,jk}(t, \cdot)\|_\infty \leq \frac{\delta}{1+t} \quad (2.9)$$

for $\delta > 0$ sufficiently small, depending on c_1 and c_2 .

We define the energy-momentum vector associated to L_γ ,

$$e_0[u] = \frac{1}{2} |\partial_t u|^2 + \frac{c_2^2}{2} |\nabla u|^2 + \frac{c_1^2 - c_2^2}{2} (\nabla \cdot u)^2 - \frac{1}{2} \gamma^{IJ,ij} \partial_i u^I \partial_j u^J, \quad (2.10)$$

$$e_k[u] = -c_2^2 \partial_t u^I \partial_k u^I - (c_1^2 - c_2^2) \partial_t u^k (\nabla \cdot u) + \gamma^{IJ,kj} \partial_j u^J \partial_t u^I, \quad k = 1, 2, 3. \quad (2.11)$$

With (2.8), it is easy to check that

$$\partial_0 e_0[u] + \partial_k e_k[u] = \partial_t u^I (L_\gamma u)^I - \frac{1}{2} (\partial_t \gamma^{IJ,ij}) \partial_i u^I \partial_j u^J + (\partial_k \gamma^{IJ,kj}) \partial_j u^J \partial_t u^I. \quad (2.12)$$

Setting

$$E_M(t) = \int \sum_{j=0}^M e_0(\partial_t^j u)(t, \mathbf{x}) \, dx,$$

the following results immediately from (2.12) and that ∂_t preserves the boundary conditions.

Lemma 2.4. Assume that γ satisfies (2.8) and (2.9), and let u be a smooth solution to (2.7) which vanishes for large $|x|$ for each t . Then

$$\partial_t E_M^{1/2}(t) \lesssim \sum_{j=0}^M \|L_\gamma \partial_t^j u(t, \cdot)\|_2 + \|\gamma'(t, \cdot)\|_\infty E_M^{1/2}(t) \quad (2.13)$$

for any fixed $M = 0, 1, 2, \dots$ ⁷ □

We next examine energy estimates involving the scaling vector field. To do so, we set

$$\tilde{S} = t\partial_t + \eta(x)r\partial_r - 1$$

where $\eta \in C^\infty(\mathbb{R}^3)$ with $\eta(x) \equiv 0$ for $x \in \mathcal{K}$ and $\eta(x) \equiv 1$ for $|x| > 1$. For

$$X_{\nu,j} = \int e_0(\tilde{S}^\nu \partial_t^j u)(t, x) dx,$$

we have the following.

Lemma 2.5. If u is a smooth solution to (2.7) and vanishes for large $|x|$ for each t , then

$$\begin{aligned} \partial_t X_{\nu,j} &\lesssim X_{\nu,j}^{1/2} \|\tilde{S}^\nu \partial_t^j L_\gamma u(t, \cdot)\|_2 + \|\gamma'(t, \cdot)\|_\infty X_{\nu,j} \\ &+ X_{\nu,j}^{1/2} \|\tilde{S}^\nu \partial_t^j [\gamma^{kj} \partial_k \partial_l] u(t, \cdot)\|_2 + X_{\nu,j}^{1/2} \sum_{\mu \leq \nu-1} \|S^\mu \partial_t^j Lu(t, \cdot)\|_2 \\ &+ X_{\nu,j}^{1/2} \sum_{\substack{|\alpha|+\mu \leq j+\nu \\ \mu \leq \nu-1}} \|S^\mu \partial^\alpha u'(t, \cdot)\|_{L^2(\{|x|<1\})} \end{aligned} \quad (2.14)$$

for any fixed ν, j .⁸ □

Proof of Lemma 2.5. As \tilde{S} and ∂_t preserve the boundary condition, it follows from (2.13) that

$$\partial_t X_{\nu,j} \lesssim X_{\nu,j}^{1/2} \|L_\gamma \tilde{S}^\nu \partial_t^j u(t, \cdot)\|_2 + \|\gamma'(t, \cdot)\|_\infty X_{\nu,j}. \quad (2.15)$$

⁷Here, we have set $\|\gamma'(t, \cdot)\|_\infty = \sum_{I,J=1}^3 \sum_{j,k=0}^3 \sum_{\beta=0}^3 \|\partial_\beta \gamma^{IJjk}(t, \cdot)\|_\infty$.

⁸For a differential operator $P = P(t, x, D_t, D_x)$, we set $[P, \gamma^{kl} \partial_k \partial_l] u = \sum_{I,J=1}^3 \sum_{k,l=1}^3 |[P, \gamma^{IJ,kl} \partial_k \partial_l] u^J|$.

We next examine the commutator

$$\begin{aligned}
|L_\gamma \tilde{S}^\nu \partial_t^j u| &\leq |\tilde{S}^\nu \partial_t^j L_\gamma u| + |[\tilde{S}^\nu \partial_t^j, \gamma^{kl} \partial_k \partial_l] u| + |[\tilde{S}^\nu, L] \partial_t^j u| \\
&\leq |\tilde{S}^\nu \partial_t^j L_\gamma u| + |[\tilde{S}^\nu \partial_t^j, \gamma^{kl} \partial_k \partial_l] u| + |[S^\nu, L] \partial_t^j u| + |[\tilde{S}^\nu - S^\nu, L] \partial_t^j u| \\
&\lesssim |\tilde{S}^\nu \partial_t^j L_\gamma u| + |[\tilde{S}^\nu \partial_t^j, \gamma^{kl} \partial_k \partial_l] u| + \sum_{\mu \leq \nu-1} |S^\mu \partial_t^j L u| + \mathbf{1}_{\{|x| < 1\}}(x) \sum_{\substack{|\alpha| + \mu \leq j + \nu \\ \mu \leq \nu-1}} |S^\mu \partial^\alpha u'|
\end{aligned}$$

Using this in (2.15) completes the proof. \blacksquare

Using the previous lemma and elliptic regularity, we shall prove.

Proposition 2.6. Assume (2.8) and (2.9) for δ sufficiently small, and let u be a smooth solution to (2.7). Suppose further that

$$\|\gamma'(t, \cdot)\|_\infty \leq \delta/(1+t) \quad (2.16)$$

and

$$\begin{aligned}
&\sum_{\substack{j+\mu \leq N+\nu \\ \mu \leq \nu}} \left(\|\tilde{S}^\mu \partial_t^j L_\gamma u(t, \cdot)\|_2 + \|[\tilde{S}^\mu \partial_t^j, \gamma^{kl} \partial_k \partial_l] u(t, \cdot)\|_2 \right) \\
&\leq \frac{\delta}{1+t} \sum_{\substack{j+\mu \leq N+\nu \\ \mu \leq \nu}} \|\tilde{S}^\mu \partial_t^j u'(t, \cdot)\|_2 + H_{\nu, N}(t) \quad (2.17)
\end{aligned}$$

for some fixed N and ν and some function $H_{\nu, N}(t)$.⁹ Then,

$$\begin{aligned}
&\sum_{\substack{|\alpha| + \mu \leq N+\nu \\ \mu \leq \nu}} \|S^\mu \partial^\alpha u'(t, \cdot)\|_2 \lesssim \sum_{\substack{|\alpha| + \mu \leq N+\nu-1 \\ \mu \leq \nu}} \|S^\mu \partial^\alpha L u(t, \cdot)\|_2 + (1+t)^{A\delta} \sum_{\substack{j+\mu \leq N+\nu \\ \mu \leq \nu}} X_{\mu, j}^{1/2}(0) \\
&+ (1+t)^{A\delta} \left(\int_0^t \sum_{\substack{|\alpha| + \mu \leq N+\nu-1 \\ \mu \leq \nu-1}} \|S^\mu \partial^\alpha L u(s, \cdot)\|_2 ds + \int_0^t H_{\nu, N}(s) ds \right) \\
&+ (1+t)^{A\delta} \int_0^t \sum_{\substack{|\alpha| + \mu \leq N+\nu \\ \mu \leq \nu-1}} \|S^\mu \partial^\alpha u'(s, \cdot)\|_{L^2(|x| < 1)} ds, \quad (2.18)
\end{aligned}$$

for some constant $A > 0$. \square

⁹In the sequel, $H_{\nu, N}(s)$ will involve $\| \langle x \rangle^{-1/2} S^\mu Z^\alpha u'(s, \cdot) \|_2^2$ for $|\alpha| + \mu \ll N + \nu$. As such, this term will be handled using the KSS estimates in the next section.

Proof of Proposition 2.6. For δ in (2.9) sufficiently small,

$$\sum_{\substack{j+\mu \leq N+\nu \\ \mu \leq \nu}} \|\tilde{S}^\mu \partial_t^j u'(t, \cdot)\|_2 \lesssim \sum_{\substack{j+\mu \leq N+\nu \\ \mu \leq \nu}} X_{\mu j}^{1/2}(t). \quad (2.19)$$

By (2.14), (2.16), and (2.17), it thus follows that

$$\begin{aligned} \partial_t \sum_{\substack{j+\mu \leq N+\nu \\ \mu \leq \nu}} X_{\mu j}^{1/2}(t) &\leq \frac{A\delta}{1+t} \sum_{\substack{j+\mu \leq N+\nu \\ \mu \leq \nu}} X_{\mu j}^{1/2}(t) + AH_{\nu, N}(t) \\ &+ A \sum_{\substack{j+\mu \leq N+\nu-1 \\ \mu \leq \nu-1}} \|S^\mu \partial_t^j Lu(t, \cdot)\|_2 + A \sum_{\substack{|\alpha|+\mu \leq N+\nu \\ \mu \leq \nu-1}} \|S^\mu \partial^\alpha u'(t, \cdot)\|_{L^2(|x|<1)}. \end{aligned} \quad (2.20)$$

By Gronwall's inequality, $\sum_{\substack{j+\mu \leq N+\nu \\ \mu \leq \nu}} X_{\mu j}^{1/2}$ satisfies the desired bound. Thus, by combining this estimate for $X_{\mu j}^{1/2}$, (2.19), and (2.3), the estimate (2.18) follows. \blacksquare

Finally, we establish the necessary energy estimates for $S^\mu Z^\alpha u$. To this end, we set

$$Y_{N, \nu}(t) = \sum_{\substack{|\alpha|+\mu \leq N+\nu \\ \mu \leq \nu}} \int e_0(S^\mu Z^\alpha u)(t, x) dx.$$

We then argue as in the proof of Lemma 2.4 with $E_M(t)$ replaced by $Y_{N, \nu}(t)$. The boundary terms that arise mesh well with the KSS estimates of the next section.

Proposition 2.7. Assume (2.8) and (2.9). Suppose that u solves (2.7) and vanishes for large $|x|$ for each t . Then,

$$\begin{aligned} \partial_t Y_{N, \nu} &\lesssim Y_{N, \nu}^{1/2} \sum_{\substack{|\alpha|+\mu \leq N+\nu \\ \mu \leq \nu}} \|L_\gamma S^\mu Z^\alpha u(t, \cdot)\|_2 + \|\gamma'(t, \cdot)\|_\infty Y_{N, \nu} \\ &+ \sum_{\substack{|\alpha|+\mu \leq N+\nu+1 \\ \mu \leq \nu}} \|S^\mu \partial^\alpha u'(t, \cdot)\|_{L^2(|x|<1)}^2. \end{aligned} \quad (2.21) \quad \square$$

Proof of Proposition 2.7. By arguing as in the proof of Lemma 2.4, it follows that

$$\begin{aligned} \partial_t Y_{N,\nu} \lesssim Y_{N,\nu}^{1/2} \sum_{\substack{|\alpha|+\mu \leq N+\nu \\ \mu \leq \nu}} \|L_\gamma S^\mu Z^\alpha u(t, \cdot)\|_2 + \|\gamma'(t, \cdot)\|_\infty Y_{N,\nu} \\ + \sum_{\substack{|\alpha|+\mu \leq N_0+\nu_0 \\ \mu \leq \nu_0}} \int_{\partial\mathcal{K}} |e_k(S^\mu Z^\alpha u)(t, y) n_k| d\sigma(y) \quad (2.22) \end{aligned}$$

where $n = (n_1, n_2, n_3)$ is the outward normal to \mathcal{K} at a point $x \in \partial\mathcal{K}$ and $e_k[\cdot]$ are as in (2.11).

Recalling that $\mathcal{K} \subset \{|x| < 1\}$,

$$\sum_{\substack{|\alpha|+\mu \leq N+\nu \\ \mu \leq \nu}} |S^\mu Z^\alpha u(t, x)| \lesssim \sum_{\substack{|\alpha|+\mu \leq N+\mu \\ \mu \leq \nu}} |S^\mu \partial^\alpha u(t, x)|, \quad x \in \partial\mathcal{K},$$

and thus, by a trace theorem, we have that the last term in (2.22) is

$$\lesssim \int_{\{x \in \mathbb{R}^3 \setminus \mathcal{K} : |x| < 1\}} \sum_{\substack{|\alpha|+\mu \leq N+\nu+1 \\ \mu \leq \nu}} |S^\mu \partial^\alpha u'(t, x)|^2 dx,$$

which completes the proof. ■

2.3 KSS estimates

In this section, we present a class of weighted, mixed-norm estimates, called KSS estimates, which are particularly useful for estimating the boundary term in (2.21) and for dealing with certain technicalities regarding the distribution of the occurrences of the scaling vector field in the proof of the main theorem. Such estimates were first used to study nonlinear problems in [8] and have played a fundamental role in previous studies of problems in exterior domains.

The estimates that we give are from [12], and we refer the interested reader to that article for detailed proofs. They are based on the boundaryless estimates for the wave equation from [25] and [17]. The corresponding estimates for elasticity in the boundaryless case follows from a Helmholtz–Hodge decomposition. In order to prove estimates in the exterior domain one relies on the decay of local energy when x is near the boundary of the obstacle and uses the boundaryless estimates when $|x|$ is large.

Proposition 2.8. Suppose that $\mathcal{K} \subset \{|x| < 1\} \subset \mathbb{R}^3$ is a nontrapping obstacle with smooth boundary. Suppose further that $u \in C^\infty$ satisfies $u|_{\partial\mathcal{K}} = 0$, $u(t, x) = 0$ for $t \leq 0$, and vanishes for large $|x|$ for every t . Then,

$$\begin{aligned} & (\log(2+T))^{-1/2} \sum_{\substack{|\alpha|+\mu \leq M+\nu \\ \mu \leq \nu}} \|\langle x \rangle^{-1/2} S^\mu \partial^\alpha u'\|_{L_t^2 L_x^2(S_T)} \\ & \lesssim \int_0^T \sum_{\substack{|\alpha|+\mu \leq M+\nu \\ \mu \leq \nu}} \|S^\mu \partial^\alpha Lu(s, \cdot)\|_2 ds + \sum_{\substack{|\alpha|+\mu \leq M+\nu-1 \\ \mu \leq \nu}} \|S^\mu \partial^\alpha Lu\|_{L_t^2 L_x^2(S_T)} \end{aligned} \quad (2.23)$$

and

$$\begin{aligned} & (\log(2+T))^{-1/2} \sum_{\substack{|\alpha|+\mu \leq M+\nu \\ \mu \leq \nu}} \|\langle x \rangle^{-1/2} S^\mu Z^\alpha u'\|_{L_t^2 L_x^2(S_T)} \\ & \lesssim \int_0^T \sum_{\substack{|\alpha|+\mu \leq M+\nu \\ \mu \leq \nu}} \|S^\mu Z^\alpha Lu(s, \cdot)\|_2 ds + \sum_{\substack{|\alpha|+\mu \leq M+\nu-1 \\ \mu \leq \nu}} \|S^\mu Z^\alpha Lu\|_{L_t^2 L_x^2(S_T)} \end{aligned} \quad (2.24)$$

for any fixed M, ν and $T > 0$. □

3 Pointwise estimates and boundary term estimates

In this section, we shall present our main pointwise decay estimates. These are analogues of the $L^1 - L^\infty$ estimates used in [9] for the wave equation. See also [15], [13, 14]. In the process of proving such estimates, we shall also prove the boundary term estimate which is required to handle the last term of (2.18). This is reminiscent of ideas from [15, 16].

We choose to employ the said $L^1 - L^\infty$ estimates rather than analogs of the $L^\infty - L^\infty$ estimates used in [6] and [1, 2]. First of all, the $L^1 - L^\infty$ estimates do not require the logarithmic factors that appear, e.g., in [1, Proposition 5.1] nor the technical arguments to show global existence with decay estimates involving a logarithmic loss. Secondly, as mentioned above, as a corollary to our proof, we obtain the boundary term estimates which we shall require in order to prove energy estimates involving the scaling vector field when there is no star-shapedness assumption. Finally, we mention that for quadratic nonlinearities, the $L^1 - L^\infty$ estimates mesh particularly well with the KSS estimates of the previous section. We do, however, note that wave equation analogs of

the $L^\infty - L^\infty$ estimates have previously been used to study multiple speed problems in exterior domains — see, e.g., [14] — and as such, it would not be surprising if these $L^\infty - L^\infty$ estimates could also be used in the current setting.

We begin by looking at the solution to the boundaryless problem

$$\begin{cases} Lv(t, x) = 0, \\ v(0, \cdot) = f, \quad \partial_t v(0, \cdot) = g. \end{cases} \quad (3.1)$$

By arguing using spherical means¹⁰, we have that

$$\begin{aligned} 4\pi v_I(t, x) &= \int_{S^2} \left[(2\delta_{IJ} - 4y_I y_J) f_J(x + c_2 t y) \right. \\ &\quad \left. + t(\delta_{IJ} - y_I y_J)(g_J(x + c_2 t y) + c_2 y_K \nabla_K f_J(x + c_2 t y)) \right] d\sigma(y) \\ &\quad + \int_{S^2} \left[(-\delta_{IJ} + 4y_I y_J) f_J(x + c_1 t y) \right. \\ &\quad \left. + t y_I y_J (g_J(x + c_1 t y) + c_1 y_K \nabla_K f_J(x + c_1 t y)) \right] d\sigma(y) \\ &\quad - \int_{c_2 t}^{c_1 t} \int_{S^2} r^{-1} (\delta_{IJ} - 3y_I y_J) (t g_J(x + r y) + f_J(x + r y)) d\sigma(y) dr. \end{aligned} \quad (3.2)$$

After a simple change of variables, this is

$$\begin{aligned} 4\pi v_I(t, x) &= \int_{S^2} \left[(2\delta_{IJ} - 4y_I y_J) f_J(x + c_2 t y) \right. \\ &\quad \left. + t(\delta_{IJ} - y_I y_J)(g_J(x + c_2 t y) + c_2 y_K \nabla_K f_J(x + c_2 t y)) \right] d\sigma(y) \\ &\quad + \int_{S^2} \left[(-\delta_{IJ} + 4y_I y_J) f_J(x + c_1 t y) \right. \\ &\quad \left. + t y_I y_J (g_J(x + c_1 t y) + c_1 y_K \nabla_K f_J(x + c_1 t y)) \right] d\sigma(y) \\ &\quad - \int_{c_2}^{c_1} \int_{S^2} c^{-1} (\delta_{IJ} - 3y_I y_J) (t g_J(x + c t y) + f_J(x + c t y)) d\sigma(y) dc. \end{aligned} \quad (3.3)$$

Using this representation, we shall be able to prove estimates on v using techniques similar to those for wave equations.

3.1 Pointwise estimates in \mathbb{R}^3

In [9], the authors adapted a $L^1 - L^\infty$ Hörmander-type estimate to eliminate the dependence on Lorentz invariance. As such, these estimates could be applied to study multiple

¹⁰See [6].

speed wave equations and certain boundary value problems. In this section, we prove analogs of these pointwise estimates.

The first of these estimates is for the homogeneous equation (3.1).

Lemma 3.1. Let v be a solution to (3.1). Then,

$$(1 + t + |x|)|v(t, x)| \lesssim \sum_{|\alpha| \leq 4} \|\langle x \rangle^{|\alpha|} \partial^\alpha f\|_2 + \sum_{|\alpha| \leq 3} \|\langle x \rangle^{1+|\alpha|} \partial^\alpha g\|_2. \quad (3.4) \quad \square$$

Proof of Lemma 3.1. Using the following estimates for the wave equation from [14]¹¹

$$(1 + t + |x|)t \int_{S^2} |h(x + ty)| d\sigma(y) \lesssim \sum_{\substack{|\alpha| + \mu \leq 3 \\ \mu \leq 1}} \int_{\mathbb{R}^3} |(|z| \partial_{|z|})^\mu Z^\alpha h(z)| \frac{dz}{\langle z \rangle} \quad (3.5)$$

and

$$(1 + t + |x|) \int_{S^2} |h(x + ty)| d\sigma(y) \lesssim \sum_{\substack{|\alpha| + \mu \leq 3 \\ \mu \leq 1}} \int_{\mathbb{R}^3} |(|z| \partial_{|z|})^\mu Z^\alpha h(z)| \frac{dz}{\langle z \rangle^2}, \quad (3.6)$$

(3.4) follows from (3.3) and the Schwarz inequality. ■

We now prove the corresponding result for the inhomogeneous equation.

Lemma 3.2. Let w solve $Lw = G$ for $(t, x) \in \mathbb{R}_+ \times \mathbb{R}^3$, and suppose that $w(t, x) = 0$ for $t \leq 0$. Then,

$$(1 + t + |x|)|w(t, x)| \lesssim \sum_{\substack{|\alpha| + \mu \leq 3 \\ \mu \leq 1}} \int_0^t \int_{\mathbb{R}^3} |S^\mu Z^\alpha G(s, y)| \frac{dy ds}{|y|}. \quad (3.7)$$

Suppose further that $G(s, y) = 0$ when $|y| > 20c_1 s$. Then,

$$(1 + t)|w(t, x)| \lesssim \sum_{\substack{|\alpha| + \mu \leq 3 \\ \mu \leq 1}} \int_{\theta t}^t \int_{\mathbb{R}^3} |S^\mu Z^\alpha G(s, y)| \frac{dy ds}{|y|}, \quad \text{for } |x| < c_2 t / 10, \quad (3.8)$$

where θ is some constant depending on c_2 . □

¹¹See the proof of Lemma 2.2.

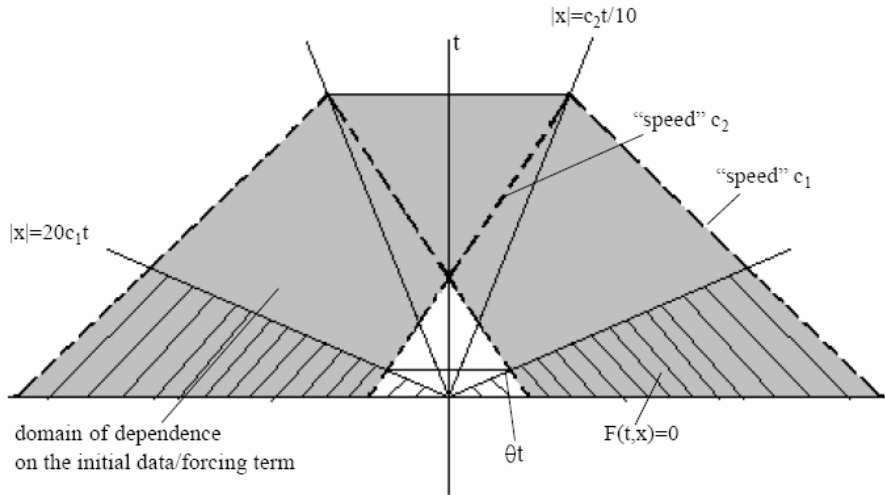


Figure 1 An illustration of the definition of θt .

Proof of Lemma 3.2. This follows as in the proof of the previous lemma, using the following bound from [9]¹²

$$\int_0^t \int_{S^2} (t-s) |F(s, x - (t-s)y)| d\sigma(y) ds \lesssim \sum_{\substack{|\alpha| + \mu \leq 3 \\ \mu \leq 1}} \int_0^t \int_{\mathbb{R}^3} |S^\mu Z^\alpha F(s, y)| \frac{dy ds}{|y|}.$$

The lemma then follows from (3.3) and Duhamel’s principle.

The second estimate follows from the appropriate Huygens’ principle for L . See Figure 1. ■

By Duhamel’s principle and the positivity of (3.3), we also have

¹²See the proof of Proposition 2.1.

Lemma 3.3. ¹³ Let v solve (3.1). Then,

$$\begin{aligned}
|x||v(t, x)| &\lesssim \sum_{i=1,2} \int_{\substack{|x|+c_i(t-s) \\ ||x|-c_i(t-s)|}}^{|x|+c_i(t-s)} \sup_{|\theta|=1} |f(\rho\theta)| d\rho + \int_{c_2}^{c_1} \int_{\substack{|x|+c(t-s) \\ ||x|-c(t-s)|}}^{|x|+c(t-s)} \sup_{|\theta|=1} |f(\rho\theta)| d\rho dc \\
&+ \sum_{i=1,2} \int_{\substack{|x|+c_i(t-s) \\ ||x|-c_i(t-s)|}}^{|x|+c_i(t-s)} \sup_{|\theta|=1} |\nabla f(\rho\theta)| \rho d\rho \\
&\quad + \sum_{i=1,2} \int_{\substack{|x|+c_i(t-s) \\ ||x|-c_i(t-s)|}}^{|x|+c_i(t-s)} \sup_{|\theta|=1} |g(\rho\theta)| \rho d\rho + \int_{c_2}^{c_1} \int_{\substack{|x|+c(t-s) \\ ||x|-c(t-s)|}}^{|x|+c(t-s)} \sup_{|\theta|=1} |g(\rho\theta)| \rho d\rho dc. \quad (3.9)
\end{aligned}$$

Moreover, let w solve $Lw = G$ for $(t, x) \in \mathbb{R}_+ \times \mathbb{R}^3$, and suppose that $w(t, x) = 0$ for $t \leq 0$. Then

$$\begin{aligned}
|x||w(t, x)| &\lesssim \sum_{i=1,2} \int_0^t \int_{\substack{|x|+c_i(t-s) \\ ||x|-c_i(t-s)|}}^{|x|+c_i(t-s)} \sup_{|\theta|=1} |G(s, \rho\theta)| \rho d\rho ds \\
&\quad + \int_{c_1}^{c_2} \int_0^t \int_{\substack{|x|+c(t-s) \\ ||x|-c(t-s)|}}^{|x|+c(t-s)} \sup_{|\theta|=1} |G(s, \rho\theta)| \rho d\rho ds dc. \quad (3.10)
\end{aligned}$$

□

3.2 Pointwise estimates in $\mathbb{R}^3 \setminus \mathcal{K}$

Here, we wish to establish the main decay estimates which shall be used in the sequel. At the first order of difficulty, these are exterior domain analogs of (3.4) and (3.7). For these estimates, we shall look at solutions to the following

$$\begin{cases} Lu(t, x) = F(t, x), & (t, x) \in \mathbb{R}_+ \times \mathbb{R}^3 \setminus \mathcal{K}, \\ u(t, x)|_{\partial\mathcal{K}} = 0, \\ u(0, x) = f(x), \quad \partial_t u(0, x) = g(x), \\ \text{supp } f, g \subset \{|x| \geq 6\}. \end{cases} \quad (3.11)$$

For such solutions, we shall have the following estimates.

¹³We note that a portion of this lemma has previously appeared in [26].

Theorem 3.4. Let $\mathcal{K} \subset \mathbb{R}^3$ be a bounded, nontrapping obstacle with smooth boundary. Then smooth solutions of (3.11) satisfy

$$\begin{aligned}
(1+t+|x|)|S^\nu Z^\alpha u(t,x)| &\lesssim \sum_{\substack{j+|\beta|+k \leq M+\nu+7 \\ j \leq 1}} \|\langle x \rangle^{j+|\beta|} \nabla^\beta \partial_t^{k+j} u(0, \cdot)\|_2 \\
&+ \int_0^t \int_{\mathbb{R}^3 \setminus \mathcal{K}} \sum_{\substack{|\beta|+\mu \leq M+\nu+6 \\ \mu \leq \nu+1}} |S^\mu Z^\beta F(s, y)| \frac{dy ds}{|y|} \\
&+ \int_0^t \sum_{\substack{|\beta|+\mu \leq M+\nu+3 \\ \mu \leq \nu+1}} \|S^\mu \partial^\beta F(s, \cdot)\|_{L^2(\{|x|<4\})} ds \quad (3.12)
\end{aligned}$$

for any $|\alpha| = M$ and any ν . □

Proof of Theorem 3.4. This proof resembles those in [9] and [15] for the wave equation quite closely. A portion of this argument was also given previously in [26].

As a first reduction, we prove that

$$\begin{aligned}
(1+t+|x|)|S^\nu Z^\alpha u(t,x)| &\lesssim \sum_{\substack{j+|\beta|+k \leq M+\nu+4 \\ j \leq 1}} \|\langle x \rangle^{j+|\beta|} \nabla^\beta \partial_t^{k+j} u(0, \cdot)\|_2 \\
&+ \int_0^t \int_{\mathbb{R}^3 \setminus \mathcal{K}} \sum_{\substack{|\beta|+\mu \leq |\alpha|+\nu+3 \\ \mu \leq \nu+1}} |S^\mu Z^\beta F(s, y)| \frac{dy ds}{|y|} \\
&+ \sup_{|y| \leq 2, 0 \leq s \leq t} (1+s) \sum_{\substack{|\beta|+\mu \leq |\alpha|+\nu \\ \mu \leq \nu}} \left(|S^\mu Z^\beta u'(s, y)| + |S^\mu Z^\beta u(s, y)| \right). \quad (3.13)
\end{aligned}$$

Proof of (3.13). To do this, we follow the proof of Lemma 4.2 of [9]. The estimate is obvious for $|x| < 2$. Thus, for the remainder of the proof, we shall assume that $|x| \geq 2$. As such, for a smooth ρ satisfying $\rho(r) \equiv 1$ for $r \geq 2$ and $\rho(r) \equiv 0$ for $r \leq 1$, it suffices to establish the estimate for $w(t, x) = \rho(|x|)S^\nu Z^\alpha u(t, x)$, which solves the boundaryless¹⁴ equation

$$\begin{aligned}
Lw &= \rho L S^\nu Z^\alpha u - 2c_2^2 \nabla \rho \cdot \nabla (S^\nu Z^\alpha u) - c_2^2 (\Delta \rho) S^\nu Z^\alpha u \\
&\quad - (c_1^2 - c_2^2) \nabla (\nabla \rho \cdot S^\nu Z^\alpha u) - (c_1^2 - c_2^2) \nabla \rho \nabla \cdot (S^\nu Z^\alpha u).
\end{aligned}$$

¹⁴Recall that we are assuming $\mathcal{K} \subset \{|x| < 1\}$.

We write $w = w_1 + w_2$ where $Lw_1 = \rho LS^\nu Z^\alpha u$ with $(w_1, \partial_t w_1)(0, \cdot) = (f, g)$. The estimate for w_1 follows from (3.4) and (3.7).

By (3.10), we see that

$$|w_2(t, x)| \lesssim \frac{1}{|x|} \sum_{i=1,2} \int_0^t \int_{|r-c_i(t-s)|}^{r+c_i(t-s)} \sup_{|\theta|=1} |G(s, \rho\theta)| \rho \, d\rho \, ds + \frac{1}{|x|} \int_{c_2}^{c_1} \int_0^t \int_{|r-c(t-s)|}^{r+c(t-s)} \sup_{|\theta|=1} |G(s, \rho\theta)| \rho \, d\rho \, ds \, dc \quad (3.14)$$

where

$$G(t, x) = Lw - \rho LS^\nu Z^\alpha u.$$

We shall now work with the last term in (3.14). The bound for the other term follows similarly. By noting that G vanishes for $|x| \geq 2$, it follows that we must have

$$-2 \leq |x| - c(t-s) \leq 2$$

for the $d\rho$ integral to be nonzero. That is, we must have

$$\frac{ct - |x| - 2}{c} \leq s \leq \frac{ct - |x| + 2}{c}.$$

Thus, we conclude that the last term in (3.14) is

$$\lesssim \frac{1}{|x|} \int_{c_2}^{c_1} \frac{1}{1 + |ct - |x||} \times \sum_{\substack{|\beta| + \mu \leq M + \nu \\ \mu \leq \nu}} \sup_{\substack{\frac{ct - |x| - 2}{c} \leq s \leq \frac{ct - |x| + 2}{c} \\ |y| \leq 2}} (1 + s) \left(|S^\mu Z^\beta u'(s, y)| + |S^\mu Z^\beta u(s, y)| \right) dc$$

from which the bound for $|w_2|$ follows.

In order to complete the proof of Theorem 3.4, it will suffice to show that

$$(1 + t) \sum_{\substack{|\beta| + \mu \leq M + \nu + 1 \\ \mu \leq \nu}} \sup_{|x| < 2} |S^\mu \partial^\beta u(t, x)|$$

is bounded by the right side of (3.12). Using smooth cutoffs, we may examine the following cases separately:

Case 1: u has vanishing data and $F(s, y)$ vanishes for $|y| > 4$

Case 2: $F(s, y) = 0$ for $|y| < 3$.

We use the following consequence of the Fundamental Theorem of Calculus and Sobolev embedding

$$\begin{aligned}
(1+t) \sum_{\substack{|\beta|+\mu \leq M+\nu+1 \\ \mu \leq \nu}} \sup_{|x|<2} |S^\mu \partial^\beta u(t, x)| & \\
& \lesssim \int_0^t \sum_{\substack{|\beta|+\mu \leq M+\nu+2 \\ \mu \leq \nu, j \leq 1}} \|(s \partial_s)^j S^\mu \partial^\beta u'(s, \cdot)\|_{L^2(\{|x|<4\})} ds \\
& \lesssim \int_0^t \sum_{\substack{|\beta|+\mu \leq M+\nu+3 \\ \mu \leq \nu+1}} \|S^\mu \partial^\beta u'(s, \cdot)\|_{L^2(\{|x|<4\})} ds.
\end{aligned} \tag{3.15}$$

Here, for the first inequality, we have also used the fact that the Dirichlet boundary conditions permit us to bound u locally by u' . The bound in *Case 1* now follows from an application of (2.5).

To complete the proof, we need only examine *Case 2*. Here, we write $u = w + u_r$ where w solves the boundaryless equation $Lw = F$ with $(w, \partial_t w)(0, \cdot) = (u, \partial_t u)(0, \cdot)$. We fix $\eta \in C^\infty(\mathbb{R}^3)$ with $\eta(x) \equiv 1$ for $|x| < 2$ and $\eta(x) \equiv 0$ for $|x| \geq 3$. Letting $\tilde{u} = \eta w + u_r$, we notice that $u = \tilde{u}$ for $|x| < 2$, that \tilde{u} solves

$$L\tilde{u} = -2c_2^2 \nabla \eta \cdot \nabla w - c_2^2 (\Delta \eta) w - (c_1^2 - c_2^2) \nabla (\nabla \eta \cdot w) - (c_1^2 - c_2^2) \nabla \eta \nabla \cdot w, \tag{3.16}$$

and that \tilde{u} has vanishing Cauchy data.

As the right side vanishes unless $2 \leq |x| \leq 3$, we may apply the result of the previous case to see that

$$(1+t) \sum_{\substack{|\beta|+\mu \leq M+\nu+1 \\ \mu \leq \nu}} \sup_{|x|<2} |S^\mu \partial^\beta \tilde{u}(t, x)| \lesssim \int_0^t \sum_{\substack{|\beta|+\mu \leq M+\nu+4 \\ \mu \leq \nu}} \|S^\mu \partial^\beta w(s, \cdot)\|_{L^\infty(2 \leq |x| \leq 3)} ds.$$

It remains only to bound the term on the right of the preceding equation. To do this, we apply (3.9), (3.10), and Sobolev's lemma on S^2 . For the forcing term, for example, we have

$$\sum_{\substack{|\beta|+\mu \leq M+\nu+6 \\ \mu \leq \nu}} \int_0^t \int_{c_2}^{c_1} \int_0^s \int_{|c(s-\tau)-|y|| \leq 4} |S^\mu Z^\beta F(\tau, y)| \frac{dy d\tau}{|y|} dc ds$$

corresponding to the last term in (3.10). For fixed c , we note that the sets $\Lambda_s = \{(\tau, y) : 0 \leq \tau \leq s, |c(s - \tau) - |y|| \leq 4\}$ have finite overlap, i.e., for $|s - s'| \geq 10$, $\Lambda_s \cap \Lambda_{s'} = \emptyset$. Thus, we see that the desired bound holds. Similar arguments can be used to estimate the remaining terms in (3.9) and (3.10), which completes the proof. \blacksquare

We will also require the following pointwise estimate which follows from arguments similar to those in [13].

Theorem 3.5. Suppose $\mathcal{K} \subset \{|x| < 1\} \subset \mathbb{R}^3$ is a nontrapping obstacle with smooth boundary. Let u solve

$$\begin{cases} Lu(t, x) = F(t, x), & (t, x) \in \mathbb{R}_+ \times \mathbb{R}^3 \setminus \mathcal{K}, \\ u|_{\partial\mathcal{K}} = 0, \\ u(t, x) = 0, & t \leq 0. \end{cases} \quad (3.17)$$

Suppose further that $F(t, x) = 0$ when $|x| > 20c_1 t$. Then, if $|x| < c_2 t/10$ and $t > 1$,

$$\begin{aligned} (1 + t + |x|)|S^\nu Z^\alpha u(t, x)| &\lesssim \sum_{\substack{|\beta| + \mu \leq |\alpha| + \nu + 6 \\ \mu \leq \nu + 1}} \int_{\theta t}^t \int_{\mathbb{R}^3 \setminus \mathcal{K}} |S^\mu Z^\beta F(s, y)| \frac{dy ds}{|y|} \\ &\quad + \sup_{0 \leq s \leq t} (1 + s) \sum_{\substack{|\beta| + \mu \leq |\alpha| + \nu + 2 \\ \mu \leq \nu}} \|S^\mu \partial^\beta F(s, \cdot)\|_2, \end{aligned} \quad (3.18)$$

for some constant θ depending on c_2 . \square

Proof of Theorem 3.5. By arguing as in (3.13) using (3.8), we see that for $|x| < c_2 t/10$,

$$\begin{aligned} (1 + t)|S^\nu Z^\alpha u(t, x)| &\lesssim \int_{\theta t}^t \int_{\mathbb{R}^3 \setminus \mathcal{K}} \sum_{\substack{|\beta| + \mu \leq |\alpha| + \nu + 3 \\ \mu \leq \nu + 1}} |S^\mu Z^\beta F(s, y)| \frac{dy ds}{|y|} \\ &\quad + \sup_{|y| \leq 2, 0 \leq s \leq t} (1 + s) \sum_{\substack{|\beta| + \mu \leq |\alpha| + \nu + 1 \\ \mu \leq \nu}} |S^\mu \partial^\beta u(s, y)|. \end{aligned} \quad (3.19)$$

Thus, we need only show that the last term is controlled by the right side of (3.18).

First suppose that $F(s, y) = 0$ if $|y| > 4$. Then, by a Sobolev estimate and (2.5), it follows that

$$(1 + t) \sup_{|y| < 2} \sum_{\substack{|\beta| + \mu \leq |\alpha| + \nu + 1 \\ \mu \leq \nu}} |S^\mu \partial^\beta u(t, x)| \lesssim \sup_{0 \leq s \leq t} \sum_{\substack{|\beta| + \mu \leq |\alpha| + \nu + 2 \\ \mu \leq \nu}} (1 + s) \|S^\mu \partial^\beta F(s, \cdot)\|_2. \quad (3.20)$$

Here we are also using that the Dirichlet boundary conditions allow us to control u locally by u' . Thus, for the remainder of the proof, we may assume that $F(s, y) = 0$ for $|y| \leq 3$.

Letting \tilde{u} be as in the proof of Theorem 3.4, we see from (3.16) and (3.20) that

$$(1+t) \sup_{|y| < 2} \sum_{\substack{|\beta| + \mu \leq |\alpha| + \nu + 1 \\ \mu \leq \nu}} |S^\mu \partial^\beta u(t, x)| \lesssim \sup_{0 \leq s \leq t} \sum_{\substack{|\beta| + \mu \leq |\alpha| + \nu + 3 \\ \mu \leq \nu}} (1+s) \|S^\mu \partial^\beta w(s, \cdot)\|_{L^\infty(2 \leq |x| \leq 3)}, \quad (3.21)$$

where w is a solution to the boundaryless equation $Lw = F$ with vanishing initial data. Provided that $s > 30/c_2$, the theorem follows from (3.8), (3.19), (3.20), and (3.21). Else, when $s \leq 30/c_2$, we may use (3.7) and finite propagation speed to bound this term by the second term in the right of (3.18), which completes the proof. \blacksquare

3.3 Boundary term estimates

In order to handle the boundary terms that arise in (2.18), we shall give an estimate which is in the spirit of the original ones of [15, 16] for the wave equation. This is the key estimate, along with the decay of local energy (2.2), that allows us to drop the star-shapedness assumption on \mathcal{K} . Here, we are merely isolating the proof of the bound for the right side of (3.15).

Lemma 3.6. Let \mathcal{K} be a bounded, nontrapping obstacle with smooth boundary. Suppose that u is a smooth solution of (3.17). Then,

$$\begin{aligned} \int_0^t \|S^\nu \partial^\alpha u'(s, \cdot)\|_{L^2(|x| < 1)} ds &\lesssim \sum_{\substack{|\beta| + \mu \leq M + \nu + 3 \\ \mu \leq \nu}} \int_0^t \int_{\mathbb{R}^3 \setminus \mathcal{K}} |S^\mu Z^\beta F(s, y)| \frac{dy ds}{|y|} \\ &+ \int_0^t \sum_{\substack{|\beta| + \mu \leq M + \nu \\ \mu \leq \nu}} \|S^\mu \partial^\beta F(s, \cdot)\|_{L^2(\{|x| < 4\})} ds \end{aligned} \quad (3.22)$$

for any $|\alpha| = M$ and ν . \square

4 Sobolev-type estimates and null form estimates

4.1 Weighted Sobolev estimates

In the sequel, we shall require the following rather standard weighted Sobolev estimate from [10]. This follows by applying Sobolev embedding on $\mathbb{R} \times S^2$. The decay results from the difference in the volume elements between $\mathbb{R} \times S^2$ and, say, \mathbb{R}^3 .

Lemma 4.1. Suppose that $h \in C^\infty(\mathbb{R}^3 \setminus \mathcal{K})$. Then, for $R \geq 1$,

$$\|h\|_{L^\infty(R/2 < |x| < R)} \lesssim R^{-1} \sum_{|\alpha|+|\beta| \leq 2} \|\tilde{\Omega}^\alpha \partial_x^\beta h\|_{L^2(R/4 < |x| < 2R)}. \quad (4.1)$$

□

4.2 Klainerman–Sideris decay estimates

In this section, we study another class of weighted Sobolev-type estimates. These follow from the boundaryless estimates of [21], which are closely related to estimates that originally appeared in [11]. The interested reader should also consult [23] for a unified approach to proving such estimates.

The main boundaryless estimate of [21]¹⁵ to consider states

$$\begin{aligned} & \| \langle c_1 t - r \rangle P_1 \partial \nabla h(t, \cdot) \|_2 + \| \langle c_2 t - r \rangle P_2 \partial \nabla h(t, \cdot) \|_2 \\ & \lesssim \sum_{|\alpha|+\mu \leq 1} \| S^\mu Z^\alpha h'(t, \cdot) \|_2 + t \| Lh(t, \cdot) \|_2. \end{aligned} \quad (4.2)$$

In order to establish similar bounds when there is a boundary, we use ideas from [13]. This yields

Lemma 4.2. Let $\mathcal{K} \subset \{|x| < 1\} \subset \mathbb{R}^3$ be an obstacle with smooth boundary. Suppose $u(t, x) \in C^\infty(\mathbb{R}_+ \times \mathbb{R}^3 \setminus \mathcal{K})$ and $u|_{\partial \mathcal{K}} = 0$. Then,

$$\begin{aligned} & \| \langle c_1 t - r \rangle P_1 \partial \nabla S^\nu Z^\alpha u(t, \cdot) \|_2 + \| \langle c_2 t - r \rangle P_2 \partial \nabla S^\nu Z^\alpha u(t, \cdot) \|_2 \\ & \lesssim \sum_{\substack{|\beta|+\mu \leq M+\nu+1 \\ \mu \leq \nu+1}} \| S^\mu Z^\alpha u'(t, \cdot) \|_2 + t \sum_{\substack{|\beta|+\mu \leq M+\nu \\ \mu \leq \nu}} \| S^\mu Z^\beta Lu(t, \cdot) \|_2 \\ & \quad + (1+t) \sum_{\mu \leq \nu} \| S^\mu u'(t, \cdot) \|_{L^2(|x| < 2)} \end{aligned} \quad (4.3)$$

for any fixed $|\alpha| = M$ and ν . □

Proof of Lemma 4.2. We first show that

$$\begin{aligned} & \sum_{\kappa=1,2} \| \langle c_\kappa t - r \rangle P_\kappa \partial \nabla S^\nu Z^\alpha u(t, \cdot) \|_2 \lesssim \sum_{\substack{|\beta|+\mu \leq M+\nu+1 \\ \mu \leq \nu+1}} \| S^\mu Z^\beta u'(t, \cdot) \|_2 \\ & \quad + \sum_{\substack{|\beta|+\mu \leq M+\nu \\ \mu \leq \nu}} t \| S^\mu Z^\beta Lu(t, \cdot) \|_2 + (1+t) \sum_{\substack{|\beta|+\mu \leq M+\nu+2 \\ \mu \leq \nu}} \| S^\mu \partial^\beta u(t, \cdot) \|_{L^2(|x| < 3/2)}. \end{aligned} \quad (4.4)$$

¹⁵See (3.24).

This estimate is trivial if the norm in the left is over $\{|x| < 3/2\}$ as the coefficients of Z are $O(1)$ on this set. To handle the case when the norms are over $\{|x| \geq 3/2\}$, we fix $\eta \in C^\infty(\mathbb{R}^3)$ with $\eta(x) \equiv 0$ for $|x| < 1$ and $\eta(x) \equiv 1$ for $|x| > 3/2$ and apply (4.2) to $h(t, x) = \eta(x)u(t, x)$ which solves the boundaryless equation

$$Lh = \eta(x)Lu - 2c_2^2 \nabla \eta \cdot \nabla u - c_2^2 (\Delta \eta) u - (c_1^2 - c_2^2) \nabla (\nabla \eta \cdot u) - (c_1^2 - c_2^2) \nabla \eta \nabla \cdot u.$$

This proves (4.4) as the last four terms are supported in $\{|x| \leq 3/2\}$.

It remains to show that the last term in (4.4) is bounded by the right side of (4.3). Here we simply apply (a trivial modification of) (2.4) and the differential inequality

$$\sum_{\kappa=1,2} |c_\kappa t - r| |P_\kappa \partial_t \nabla v(t, x)| \lesssim \sum_{|\alpha|+\mu \leq 1} |S^\mu Z^\alpha \nabla v(t, x)| + t |Lv(t, x)|,$$

which is also from [21].¹⁶ ■

We also have the following variants of the above.

Corollary 4.3. Let $\mathcal{K} \subset \{|x| < 1\} \subset \mathbb{R}^3$ be an obstacle with smooth boundary. Suppose $u(t, x) \in C^\infty(\mathbb{R}_+ \times \mathbb{R}^3 \setminus \mathcal{K})$ and $u|_{\partial \mathcal{K}} = 0$. Then,

$$\begin{aligned} |x| \langle c_\kappa t - r \rangle |P_\kappa \partial \nabla S^\nu Z^\alpha u(t, x)| &\lesssim \sum_{\substack{|\beta|+\mu \leq M+\nu+3 \\ \mu \leq \nu+1}} \|S^\mu Z^\beta u'(t, \cdot)\|_2 \\ &+ \sum_{\substack{|\beta|+\mu \leq M+\nu+2 \\ \mu \leq \nu}} t \|S^\mu Z^\beta Lu(t, \cdot)\|_2 + (1+t) \sum_{\mu \leq \nu} \|S^\mu u'(t, \cdot)\|_{L^2(|x|<2)}, \end{aligned} \quad (4.5)$$

$$\begin{aligned} \|\langle c_\kappa t - r \rangle P_\kappa \partial \nabla S^\nu Z^\alpha u(t, \cdot)\|_{L^2(|x|>t/4)} &\lesssim \sum_{\substack{|\beta|+\mu \leq M+\nu+1 \\ \mu \leq \nu+1}} \|S^\mu Z^\beta u'(t, \cdot)\|_{L^2(|x|>t/8)} \\ &+ (1+t) \sum_{\substack{|\beta|+\mu \leq M+\nu \\ \mu \leq \nu}} \|S^\mu Z^\beta Lu(t, \cdot)\|_{L^2(|x|>t/8)}, \end{aligned} \quad (4.6)$$

¹⁶See (3.10b).

and

$$\begin{aligned} \sup_{|x|>t/2} |x| \langle c_\kappa t - r \rangle |P_\kappa \partial \nabla S^\nu Z^\alpha u(t, \cdot)| &\lesssim \sum_{\substack{|\beta|+\mu \leq |\alpha|+\nu+3 \\ \mu \leq \nu+1}} \|S^\mu Z^\beta u'(t, \cdot)\|_{L^2(|x|>t/8)} \\ &+ (1+t) \sum_{\substack{|\beta|+\mu \leq |\alpha|+\nu+2 \\ \mu \leq \nu}} \|S^\mu Z^\beta Lu(t, \cdot)\|_{L^2(|x|>t/8)} \quad (4.7) \end{aligned}$$

for any $\kappa = 1, 2$ and fixed $|\alpha| = M$ and ν . \square

Proof of Corollary 4.3. Estimates (4.5) and (4.7) follow from (4.3) and (4.6), respectively, using (4.1). Estimate (4.6) follows from applying (4.2) to $h(t, x) = \eta(x/\langle t \rangle)u(t, x)$ where $\eta(z) \equiv 1$ for $|z| > 1/4$ and $\eta(z) \equiv 0$ for $|z| < 1/8$. Everything follows trivially unless two derivatives land on the cutoff. Here, our task would be complete if we could show the following variant of a Hardy inequality,

$$\frac{1}{t} \|v(t, \cdot)\|_{L^2(|x|>t/8)} \lesssim \|v'(t, \cdot)\|_{L^2(|x|>t/8)}.$$

This follows by integrating the ($n = 3$) identity

$$-\partial_j \left(\frac{x_j}{r^2} v^2 \right) = -\frac{1}{r^2} v^2 - \frac{2}{r} v \partial_r v$$

over $\{|x| \geq t/8\}$, which yields

$$0 < \int_{\{|x|=t/8\}} \frac{1}{r} v^2 d\sigma = - \int_{\{|x|>t/8\}} \frac{1}{r^2} v^2 dx - \int_{\{|x|>t/8\}} \frac{2}{r} v \partial_r v dx.$$

Applying the Schwarz inequality to the last term and bootstrapping yield the desired estimate. \blacksquare

We also have

Lemma 4.4. Let $\mathcal{K} \subset \{|x| < 1\} \subset \mathbb{R}^3$ be an obstacle with smooth boundary. Suppose $u(t, x) \in C^\infty(\mathbb{R}_+ \times \mathbb{R}^3 \setminus \mathcal{K})$ and $u|_{\partial \mathcal{K}} = 0$. Then,

$$\begin{aligned} \langle r \rangle \langle c_\kappa t - r \rangle^{1/2} |P_\kappa \partial S^\nu Z^\alpha u(t, x)| &\lesssim \sum_{\substack{|\beta|+\mu \leq M+\nu+2 \\ \mu \leq \nu+1}} \|S^\mu Z^\beta u'(t, \cdot)\|_2 \\ &+ (1+t) \sum_{\substack{|\beta|+\mu \leq M+\nu+1 \\ \mu \leq \nu}} \|S^\mu Z^\beta Lu(t, \cdot)\|_2 + (1+t) \sum_{\mu \leq \nu} \|S^\mu u'(t, \cdot)\|_{L^2(|x|<2)} \quad (4.8) \end{aligned}$$

and

$$\begin{aligned} \sup_{|x|>t/4} \langle r \rangle \langle c_\kappa t - r \rangle^{1/2} |P_\kappa \partial S^\nu Z^\alpha u(t, \mathbf{x})| &\lesssim \sum_{\substack{|\beta|+\mu \leq M+\nu+2 \\ \mu \leq \nu+1}} \|S^\mu Z^\beta u'(t, \cdot)\|_{L^2(|x|>t/8)} \\ &+ (1+t) \sum_{\substack{|\beta|+\mu \leq M+\nu+1 \\ \mu \leq \nu}} \|S^\mu Z^\beta Lu(t, \cdot)\|_{L^2(|x|>t/8)} \quad (4.9) \end{aligned}$$

for any $\kappa = 1, 2$ and fixed $|\alpha| = M$ and ν . \square

Proof of Lemma 4.4. Using the arguments of the proof of Lemma 4.2, this follows from the boundaryless estimate

$$\langle r \rangle \langle c_\kappa t - r \rangle^{1/2} |P_\kappa \partial h(t, \mathbf{x})| \lesssim \sum_{\substack{|\beta|+\mu \leq 2 \\ \mu \leq 1}} \|S^\mu Z^\beta h'(t, \cdot)\|_2 + \langle t \rangle \sum_{|\beta| \leq 1} \|Z^\beta Lh(t, \cdot)\|_2$$

of [21].¹⁷ \blacksquare

4.3 Null form estimates

The additional decay afforded to us by the null condition is encapsulated in the following estimate of [21].¹⁸ This is closely related to those for multiple speed systems of wave equations used, e.g., in [24].

Lemma 4.5. Assume that Q satisfies the null condition (1.4). Let $\mathcal{N} = \{(\alpha, \beta, \gamma) \neq (1, 1, 1), (2, 2, 2)\}$ be the set of nonresonant indices. Then,

$$\begin{aligned} &|\langle u, Q(v, w) \rangle| \\ &\lesssim \frac{1}{r} |u| \sum_{|\alpha| \leq 1} \left[|\nabla \tilde{Q}^\alpha v| |\nabla w| + |\nabla \tilde{Q}^\alpha w| |\nabla v| + |\nabla^2 v| |\tilde{Q}^\alpha w| + |\nabla^2 w| |\tilde{Q}^\alpha v| \right] \\ &\quad + \sum_{\mathcal{N}} |P_\alpha u| \left[|P_\beta \nabla^2 v| |P_\gamma \nabla w| + |P_\beta \nabla^2 w| |P_\gamma \nabla v| \right], \quad (4.10) \end{aligned}$$

for sufficiently regular u, v, w . \square

¹⁷See (3.20c) and (3.24).

¹⁸See Proposition 3.2.

5 Proof of global existence

In this section, we prove Theorem 1.1. We take $N = 112$ in (1.6), but this is far from optimal. By scaling in the t variable, we may take $c_1 = 1$ without loss of generality.

We begin by making a reduction that allows us to avoid technicalities related to the compatibility conditions. While the reduction from [9]¹⁹ works when we have truncated at the quadratic level, the necessary scaling breaks down when the higher order terms are present. We instead use a reduction that is more reminiscent of that from [14].

We begin by noting that if ε in (1.6) is sufficiently small, then there is a constant C_0 for which

$$\sup_{0 \leq t \leq 2} \sum_{|\alpha| \leq 112} \|\partial^\alpha u(t, \cdot)\|_{L^2(|x| \leq 25)} \leq C_0 \varepsilon. \quad (5.1)$$

This follows from well-known local existence theory. See, e.g., [7], which is only stated for diagonal wave equations but as the proofs only rely on energy estimates the results carry over to the current setting.

On the other hand, over $\{|x| > 5(t+1)\}$, u corresponds to a boundaryless solution.²⁰ Thus by the estimates which correspond to those that follow for the boundaryless problem²¹, we have

$$\begin{aligned} & \sup_{0 \leq t < \infty} \sum_{|\alpha| + \mu \leq 111} \|S^\mu Z^\alpha u'(t, \cdot)\|_{L^2(|x| \geq 5(t+1))} \\ & \quad + \sup_{\substack{|x| \geq 5(t+1) \\ 0 \leq t < \infty}} (1+t+|x|) \sum_{|\alpha| + \mu \leq 108} |S^\mu Z^\alpha u'(t, x)| \leq C_1 \varepsilon. \end{aligned} \quad (5.2)$$

We fix a smooth cutoff function η with $\eta(t, x) \equiv 1$ if $t \leq 3/2$ and $|x| \leq 20$, $\eta(t, \cdot) \equiv 0$ for $t > 2$, and $\eta(\cdot, x) \equiv 0$ for $|x| > 25$. Setting $u_0 = \eta u$, which by local existence theory is known and satisfies (5.1), it follows that u solves (1.5) for $0 < t < T$ if and only if

¹⁹See also [15] and [13].

²⁰Recall that $\mathcal{K} \subset \{|x| < 1\}$ by assumption.

²¹See, e.g., [21] and [14].

$w = u - u_0$ solves

$$\begin{cases} Lw = (1 - \eta)Q(\nabla u, \nabla^2 u) - [L, \eta]u, & (t, \mathbf{x}) \in \mathbb{R}_+ \times \mathbb{R}^3 \setminus \mathcal{K}, \\ w|_{\partial\mathcal{K}} = 0, \\ w(0, \cdot) = (1 - \eta)(0, \cdot)f, \\ \partial_t w(0, \cdot) = (1 - \eta)(0, \cdot)g - \eta_t(0, \cdot)f \end{cases} \quad (5.3)$$

over the same interval.

We now fix another smooth cutoff β with $\beta(z) \equiv 1$ for $z \geq 10$ and $\beta(z) \equiv 0$ for $z \leq 6$. Then, let v solve the linear equation

$$\begin{cases} Lv = \beta\left(\frac{|\mathbf{x}|}{t+1}\right)(1 - \eta)Q(\nabla u, \nabla^2 u) - [L, \eta]u, \\ v|_{\partial\mathcal{K}} = 0, \\ v(0, \cdot) = (1 - \eta)(0, \cdot)f, \\ \partial_t v(0, \cdot) = (1 - \eta)(0, \cdot)g - \eta_t(0, \cdot)f. \end{cases} \quad (5.4)$$

We shall show that

$$\begin{aligned} (1 + t + |\mathbf{x}|) \sum_{|\alpha|+\mu \leq 102} |S^\mu Z^\alpha v(t, \mathbf{x})| + \sum_{|\alpha|+\mu \leq 100} \|S^\mu Z^\alpha v'(t, \cdot)\|_2 \\ + (\log(2 + t))^{-1} \sum_{|\alpha|+\mu \leq 98} \|\langle \mathbf{x} \rangle^{-1/2} S^\mu Z^\alpha v'\|_{L_t^2 L_x^2(S_t)} \leq C_2 \varepsilon \end{aligned} \quad (5.5)$$

for some absolute constant $C_2 > 0$ and for any $t > 0$.

Proof of (5.5). We first notice that by (3.12) and (1.6) the first term in the left side of (5.5) is

$$\begin{aligned} \lesssim \varepsilon + \int_0^t \int_{\mathbb{R}^3 \setminus \mathcal{K}} \sum_{|\alpha|+\mu \leq 108} |S^\mu Z^\alpha \beta\left(\frac{|\mathbf{x}|}{t+1}\right)(1 - \eta)Q(\nabla u, \nabla^2 u)| \frac{dy ds}{|y|} \\ + \int_0^t \int_{\mathbb{R}^3 \setminus \mathcal{K}} \sum_{|\alpha|+\mu \leq 108} |S^\mu Z^\alpha [L, \eta]u| \frac{dy ds}{|y|}. \end{aligned} \quad (5.6)$$

Here, we have also used that Sobolev's lemma and the assumption that $0 \in \mathcal{K}$ allow us to control the last term in the right of (3.12) by the one which precedes it. Since $[L, \eta]u$ vanishes unless $t \leq 2$ and $|\mathbf{x}| \leq 25$, the last term is clearly $O(\varepsilon)$ by (5.1). The second term

in (5.6) is

$$\lesssim \int_0^t \int_{|y| \geq 6(s+1)} \sum_{|\beta| + \mu \leq 108} |S^\mu Z^\beta \nabla u(s, y)| \sum_{|\beta| + \mu \leq 108} |S^\mu Z^\beta \nabla^2 u(s, y)| \frac{dy ds}{|y|}. \quad (5.7)$$

Over $|x| > 6(t+1)$, we note that by (a trivial modification of) (4.7) and (4.9)

$$\begin{aligned} & \sup_{|x| > 6(t+1)} \left(\langle c_\kappa t - r \rangle \langle r \rangle |P_\kappa \nabla^2 S^\mu Z^\beta u(t, x)| + \langle r \rangle \langle c_\kappa t - r \rangle^{1/2} |P_\kappa \nabla S^\mu Z^\beta u(t, x)| \right) \\ & \lesssim \sum_{|\alpha| + \nu \leq 111} \|S^\nu Z^\alpha u'(t, \cdot)\|_{L^2(|x| > 5(t+1))} + \sum_{|\alpha| + \nu \leq 110} \langle t \rangle \|S^\nu Z^\alpha Q(\nabla u, \nabla^2 u)\|_{L^2(|x| > 5(t+1))} \end{aligned} \quad (5.8)$$

for $\kappa = 1, 2$ and for $|\beta| + \mu \leq 108$. The right side of (5.8) is easily seen to be $O(\varepsilon)$ from (5.2).

Applying this bound to the terms of (5.7) and recalling that we are assuming $c_1 = 1$, it follows that (5.7) is

$$\lesssim \varepsilon^2 \int_0^t \frac{1}{s^{(3/2)^-}} \int_{\{y \in \mathbb{R}^3 \setminus \mathcal{K} : |y| \geq 6(s+1)\}} \frac{1}{|y|^{3+}} dy ds \lesssim \varepsilon^2$$

as desired.

For the second term in the left side of (5.5), we use the standard energy integral method and see that

$$\begin{aligned} & \partial_t \sum_{|\alpha| + \nu \leq 100} \|(S^\nu Z^\alpha v)'(t, \cdot)\|_2^2 + \partial_t \sum_{|\alpha| + \nu \leq 100} \|\nabla \cdot (S^\nu Z^\alpha v)(t, \cdot)\|_2^2 \\ & \lesssim \left(\sum_{|\alpha| + \nu \leq 100} \|S^\nu Z^\alpha v'(t, \cdot)\|_2 \right) \left(\sum_{|\alpha| + \nu \leq 100} \|S^\nu Z^\alpha Lv(t, \cdot)\|_2 \right) \\ & \quad + \sum_{|\alpha| + \nu \leq 100} \left| \int_{\partial \mathcal{K}} \partial_t S^\nu Z^\alpha v(t, \cdot) \nabla S^\nu Z^\alpha v(t, \cdot) \cdot n d\sigma \right| \end{aligned}$$

where n is the outward unit normal to \mathcal{K} at $x \in \partial\mathcal{K}$. Since $\mathcal{K} \subset \{|x| < 1\}$, we may use (1.6), (5.4), and a trace theorem to see that

$$\begin{aligned}
& \sum_{|\alpha|+|\nu| \leq 100} \|S^\nu Z^\alpha v'(t, \cdot)\|_2^2 \lesssim \varepsilon^2 \\
& + \left(\int_0^t \sum_{|\alpha|+|\nu| \leq 100} \|S^\nu Z^\alpha \beta(|x|/(s+1))(1-\eta)(s, \cdot) Q(\nabla u, \nabla^2 u)(s, \cdot)\|_2 ds \right)^2 \\
& + \left(\int_0^t \sum_{|\alpha|+|\nu| \leq 100} \|S^\nu Z^\alpha [L, \eta] u(s, \cdot)\|_2 ds \right)^2 \\
& + \int_0^t \sum_{|\alpha|+|\nu| \leq 101} \|S^\nu \partial^\alpha v'(s, \cdot)\|_{L^2(|x|<1)}^2 ds. \tag{5.9}
\end{aligned}$$

Using the bound for the first term in (5.5), it is easy to see that the last term is $O(\varepsilon^2)$. By the support properties of η and (5.1), the preceding term satisfies the same bound. The desired bound for the second term in the right of (5.9) follows from (5.8) and (5.2) as above.

It remains only to show that the last term in the left of (5.5) is $O(\varepsilon)$. Here, it suffices to show that

$$\sum_{|\alpha|+\mu \leq 98} \int_0^t \|\langle x \rangle^{-1/2} S^\mu Z^\alpha v'(s, \cdot)\|_{L^2(|x| \leq c_2 s/2)}^2 ds \tag{5.10}$$

is $O(\varepsilon^2(\log(2+t))^2)$ as the estimate follows trivially from the bound for the second term in (5.5) when $|x| > c_2 s/2$.

For this, we note that by (4.8)

$$\begin{aligned}
\langle r \rangle \langle c_\kappa t - r \rangle^{1/2} \sum_{|\alpha|+\mu \leq 98} |P_\kappa \partial S^\mu Z^\alpha v(t, x)| & \lesssim \sum_{|\alpha|+\mu \leq 100} \|S^\mu Z^\alpha v'(t, \cdot)\|_2 \\
& + \sum_{|\alpha|+\mu \leq 99} (1+t) \|S^\mu Z^\alpha L v(t, \cdot)\|_2 + (1+t) \sum_{\mu \leq 98} \|S^\mu v'(t, \cdot)\|_{L^\infty(|x|<2)}. \tag{5.11}
\end{aligned}$$

The first and last terms in the right are $O(\varepsilon)$ by the bounds for the preceding terms in (5.5). The second term in the right is

$$\begin{aligned}
& \lesssim \sum_{|\alpha|+\mu \leq 99} (1+t) \|S^\mu Z^\alpha \beta(|\cdot|/(t+1))(1-\eta)(t, \cdot) Q(\nabla u, \nabla^2 u)\|_2 \\
& + \sum_{|\alpha|+\mu \leq 99} (1+t) \|S^\mu Z^\alpha [L, \eta] u(t, \cdot)\|_2.
\end{aligned}$$

Each of these terms are $O(\varepsilon)$ using (5.2) and (5.1), respectively.

Using the $O(\varepsilon)$ bound for the right side of (5.11), it follows that (5.10) is

$$\lesssim \varepsilon^2 \int_0^t \frac{1}{1+s} \|\langle x \rangle^{-3/2}\|_{L^2(|x| \leq c_2 s/2)}^2 ds \lesssim \varepsilon^2 (\log(2+t))^2$$

which completes the proof. \blacksquare

The estimate (5.5) allows us in most instances to restrict our attention to $w - v$ which solves

$$\begin{cases} L(w - v) = (1 - \beta) \left(\frac{|x|}{t+1} \right) (1 - \eta)(t, x) Q(\nabla u, \nabla^2 u), & (t, x) \in \mathbb{R}_+ \times \mathbb{R}^3 \setminus \mathcal{K}, \\ (w - v)|_{\partial \mathcal{K}} = 0, \\ (w - v)(t, x) = 0, & t \leq 0. \end{cases} \quad (5.12)$$

This equation meets the requirements of our estimates in the previous sections. In particular, it has vanishing Cauchy data and a forcing term which is supported in $|x| \leq 20c_1 t$. Depending on the estimates being utilized, we shall use certain L^2 and L^∞ bounds for u while at other times we shall use them for $w - v$ or w . Since $u = (w - v) + v + u_0$ and since u_0 and v satisfy (5.1) and (5.5), respectively, it will always be the case that a bound for $w - v$, w , or u will imply the same bound for the others.

We now set up a continuity argument. From this point forward, the argument resembles that of [13] quite closely. For $\varepsilon > 0$ as above, we assume that we have a solution to (1.5) on $0 \leq t \leq T$ satisfying

$$\sum_{\substack{|\alpha|+\nu \leq 52 \\ \nu \leq 1}} \|S^\nu Z^\alpha w'(t, \cdot)\|_2 \leq A_0 \varepsilon, \quad (5.13)$$

$$(1 + t + r) \sum_{|\alpha| \leq 40} |Z^\alpha w'(t, x)| \leq A_1 \varepsilon, \quad (5.14)$$

$$(1 + t + r) \sum_{\substack{|\alpha|+\nu \leq 55 \\ \nu \leq 2}} |S^\nu Z^\alpha (w - v)(t, x)| \leq B_1 \varepsilon^2 (1 + t)^{1/10} \log(2 + t), \quad (5.15)$$

$$\sum_{|\alpha| \leq 100} \|\partial^\alpha u'(t, \cdot)\|_2 \leq B_2 \varepsilon (1 + t)^{1/40}, \quad (5.16)$$

$$\sum_{\substack{|\alpha|+\nu \leq 72 \\ \nu \leq 3}} \|S^\nu Z^\alpha u'(t, \cdot)\|_2 \leq B_3 \varepsilon (1 + t)^{1/20}, \quad (5.17)$$

$$\sum_{\substack{|\alpha|+\nu \leq 70 \\ \nu \leq 3}} \|\langle x \rangle^{-1/2} S^\nu Z^\alpha u'\|_{L_t^2 L_x^2(S_t)} \leq B_4 \varepsilon (1 + t)^{1/20} (\log(2 + t))^{1/2}. \quad (5.18)$$

We may take $A_0 = A_1 = 4C_2$ in the first two estimates, where C_2 is as in (5.5). For sufficiently small ε , all of the above estimates hold for, say, $T = 2$ by the local existence theory. Thus, for ε sufficiently small, we prove

- (i) (5.13) is valid with A_0 replaced by $A_0/2$,
- (ii) Assuming (i), (5.14) is valid with A_1 replaced by $A_1/2$,
- (iii) (5.15)–(5.18) follow from (5.13) and (5.14) for fixed constants B_i .

For such an ε , this then proves that a solution exists for all $t > 0$.

5.1 Preliminary results

We begin with the following preliminary results assuming (5.13)–(5.18):

$$r\langle c_\kappa t - r \rangle |P_\kappa \partial \nabla S^\nu Z^\alpha u(t, x)| \lesssim \varepsilon(1+t)^{3/20} \log(2+t) \quad (5.19)$$

$$\langle r \rangle \langle c_\kappa t - r \rangle^{1/2} |P_\kappa \partial S^\nu Z^\alpha u(t, x)| \lesssim \varepsilon(1+t)^{3/20} \log(2+t), \quad (5.20)$$

and

$$(1+t) \|S^\nu Z^\alpha Lu(t, \cdot)\|_2 \lesssim \varepsilon(1+t)^{3/20} \log(2+t) \quad (5.21)$$

for $\kappa = 1, 2$, $\nu \leq 1$ and $|\alpha| + \nu \leq 63$.

Using (4.5) and (4.8), the first two estimates above follow from (5.21) via (5.1), (5.5), (5.15), and (5.17). To show the latter estimate, we expand the left side to see that it is

$$\lesssim (1+t) \left\| \sum_{\substack{|\alpha|+\mu \leq 32 \\ \mu \leq 1}} |S^\mu Z^\alpha u'(t, \cdot)| \sum_{\substack{|\alpha|+\mu \leq 64 \\ \mu \leq 1}} |S^\mu Z^\alpha u'(t, \cdot)| \right\|_2.$$

By (5.15) and (5.17), as well as (5.1) and (5.5), the preceding equation is clearly bounded by the right side of (5.21). Notice that the same proof shows that (5.19) and (5.20) hold with u replaced by $w - v$.

5.2 Proof of (i.)

As the better estimate (5.5) holds for v , it suffices to show that

$$\sum_{\substack{|\alpha|+\nu \leq 52 \\ \nu \leq 1}} \|S^\nu Z^\alpha (w - v)'(t, \cdot)\|_2^2 \lesssim \varepsilon^3. \quad (5.22)$$

Using (2.12) with $\gamma^{ljk} \equiv 0$ and with $e_\mu[u]$ replaced by $e_\mu[S^\nu Z^\alpha(w-v)]$, we see that the left side of (5.22) is

$$\begin{aligned} &\lesssim \sum_{\substack{|\alpha|+\nu \leq 52 \\ \nu \leq 1}} \int_0^t \int_{\mathbb{R}^3 \setminus \mathcal{K}} |\langle \partial_t S^\nu Z^\alpha(w-v), LS^\nu Z^\alpha(w-v) \rangle| dy ds \\ &\quad + \sum_{\substack{|\alpha|+\nu \leq 52 \\ \nu \leq 1}} \left| \int_0^t \int_{\partial \mathcal{K}} \partial_t S^\nu Z^\alpha(w-v) \partial_a S^\nu Z^\alpha(w-v) n_a d\sigma ds \right|. \end{aligned} \quad (5.23)$$

Recalling that $\mathcal{K} \subset \{|x| < 1\}$, and thus that the coefficients of Z are $O(1)$ on $\partial \mathcal{K}$, it follows that the last term is bounded by

$$\sum_{\substack{|\alpha|+\nu \leq 53 \\ \nu \leq 1}} \int_0^t \int_{\{|x| < 1\}} |S^\nu \partial^\alpha(w-v)'(s, y)|^2 dy ds.$$

This boundary term is seen to be $O(\varepsilon^4)$ using (5.15).

By (1.7), (5.12), and the fact that the vector fields preserve the null structure²², we may apply (4.10), and thus it suffices to dominate

$$\begin{aligned} &\int_0^t \int_{\mathbb{R}^3 \setminus \mathcal{K}} \frac{1}{|Y|} \sum_{\substack{|\alpha|+\nu \leq 53 \\ \nu \leq 1}} |S^\nu Z^\alpha u| \sum_{\substack{|\alpha|+\nu \leq 53 \\ \nu \leq 1}} |S^\nu Z^\alpha \nabla u| \sum_{\substack{|\alpha|+\nu \leq 52 \\ \nu \leq 1}} |S^\nu Z^\alpha \partial_t(w-v)| dy ds \\ &\quad + \int_0^t \int_{\mathbb{R}^3 \setminus \mathcal{K}} \sum_{\substack{1 \leq I, J, K \leq 2 \\ (I, J) \neq (J, K)}} \sum_{\substack{|\alpha|+\nu \leq 52 \\ \nu \leq 1}} |P_K \partial S^\nu Z^\alpha(w-v)| \sum_{\substack{|\alpha|+\nu \leq 52 \\ \nu \leq 1}} |P_I \nabla S^\nu Z^\alpha u| \\ &\quad \times \sum_{\substack{|\alpha|+\nu \leq 53 \\ \nu \leq 1}} |P_J \nabla S^\nu Z^\alpha u| dy ds. \end{aligned} \quad (5.24)$$

For the first term in (5.24), we apply (5.1), (5.5), and (5.15) to see that it is

$$\begin{aligned} &\lesssim \varepsilon \int_0^t \frac{\log(2+s)}{(1+s)^{9/10}} \sum_{\substack{|\alpha|+\nu \leq 53 \\ \nu \leq 1}} \|\langle Y \rangle^{-1/2} S^\nu Z^\alpha u'(s, \cdot)\|_{L_y^2} \\ &\quad \times \sum_{\substack{|\alpha|+\nu \leq 52 \\ \nu \leq 1}} \|\langle Y \rangle^{-1/2} S^\nu Z^\alpha(w-v)'(s, \cdot)\|_{L_y^2} ds. \end{aligned}$$

By the Schwarz inequality, (5.1), (5.5), and (5.18), this term's contribution is $O(\varepsilon^3)$.

²²See [21], Proposition 3.1

To be concrete, we will focus on the case that $I \neq K$, $I = J$ for the second term in (5.24). A symmetric argument will then complete the proof. For $\delta < (c_1 - c_2)/2$,

$$\{y \in \mathbb{R}^3 \setminus \mathcal{K} : |y| \in [(1-\delta)c_I s, (1+\delta)c_I s]\} \cap \{y \in \mathbb{R}^3 \setminus \mathcal{K} : |y| \in [(1-\delta)c_K s, (1+\delta)c_K s]\} = \emptyset.$$

Thus, it suffices to show the estimate over the complements of each of these sets separately.

By (5.20), it follows that over $\{|y| \notin [(1-\delta)c_K s, (1+\delta)c_K s]\}$ the second term in (5.24) is dominated by

$$\varepsilon \int_0^t \frac{\log(2+s)}{(1+s)^{7/20}} \int_{\{|y| \notin [(1-\delta)c_K s, (1+\delta)c_K s]\}} \frac{1}{\langle y \rangle} \sum_{\substack{|\alpha|+\mu \leq 53 \\ \mu \leq 2}} |\partial S^\mu Z^\alpha u|^2 dy ds.$$

The required $O(\varepsilon^3)$ bound follows from (5.18). A symmetric argument then completes the proof of (i).

5.3 Proof of (ii.)

The estimate follows trivially from (4.1) and (5.22) when $|x| > c_2 t/10$. By (5.5), it again suffices to show (5.14) when w is replaced by $w - v$. With this substitution, it follows from (3.18) that the left side of (5.14) is dominated by

$$\begin{aligned} & \sum_{\substack{|\alpha|+\nu \leq 46 \\ \nu \leq 1}} \int_{\theta t}^t \int_{\mathbb{R}^3 \setminus \mathcal{K}} |S^\nu Z^\alpha Q(\nabla u, \nabla^2 u)| \frac{dy ds}{|y|} \\ & + \sup_{0 \leq s \leq t} (1+s) \sum_{|\alpha| \leq 42} \|\partial^\alpha Q(\nabla u, \nabla^2 u)(s, \cdot)\|_2. \end{aligned} \quad (5.25)$$

For $|y| \geq c_2 s/2$, the first term is trivially $O(\varepsilon^2)$ by (5.13). Thus, it remains to control

$$\begin{aligned} & \int_{\theta t}^t \int_{\{|y| \leq c_2 s/2\}} \sum_{\substack{|\alpha|+\nu \leq 46 \\ \nu \leq 1}} |S^\nu Z^\alpha Q(\nabla u, \nabla^2 u)| \frac{dy ds}{|y|} \\ & \lesssim \sum_{\kappa_1, \kappa_2=1,2} \int_{\theta t}^t \int_{\{|y| \leq c_2 s/2\}} \sum_{\substack{|\beta|+\mu \leq 46 \\ \mu \leq 1}} |r \langle c_{\kappa_1} s - r \rangle^{1/2} P_{\kappa_1} \nabla S^\mu Z^\beta u| \\ & \times \sum_{\substack{|\alpha|+\nu \leq 46 \\ \nu \leq 1}} |r \langle c_{\kappa_2} s - r \rangle P_{\kappa_2} \nabla^2 S^\nu Z^\alpha u| \frac{dy ds}{|y|^3 \langle s \rangle^{3/2}}, \end{aligned}$$

which is $O(\varepsilon^2)$ by (5.19) and (5.20).

Since the second term in (5.25) is easily seen to be $O(\varepsilon^2)$ using (5.13) and (5.14), as well as (5.1), the proof of (ii.) is complete.

5.4 Proof of (iii.)

The remainder of the proof of Theorem 1.1 follows from the proof given in [15], which we sketch for completeness. In this section, we will apply the results of Section 2.2 with

$$\gamma^{IJ,lm} = -2B_{lmn}^{IJK} \partial_n u^K.$$

The necessary hypotheses (2.8) and (2.9) follow from (1.3) and (5.14), respectively. The additional hypothesis (2.16) also follows from (5.14).

We begin by proving (5.16). We first examine the case $\partial^\alpha u' = \partial_t^{|\alpha|} u'$. In the sequel, we shall use elliptic regularity to prove the general result.

For $M \leq 100$, we apply (2.13) and (5.14) to see that

$$\partial_t E_M^{1/2}(u)(t) \lesssim \sum_{j=0}^M \|L_\gamma \partial_t^j u(t, \cdot)\|_2 + \frac{\varepsilon}{1+t} E_M^{1/2}(u)(t). \quad (5.26)$$

For fixed $M \leq 100$, we have

$$\begin{aligned} \sum_{j \leq M} |L_\gamma \partial_t^j u| &\lesssim \sum_{|\alpha| \leq 40} |\partial^\alpha u'| \sum_{j \leq M-1} |\partial_t^j \partial^2 u| + \sum_{|\alpha| \leq M-41} |\partial^\alpha u'| \sum_{40 \leq |\alpha| \leq M/2} |\partial^\alpha u'| \\ &\lesssim \frac{\varepsilon}{1+t} \sum_{j \leq M-1} |\partial_t^j \partial^2 u| + \sum_{|\alpha| \leq M-41} |\partial^\alpha u'| \sum_{|\alpha| \leq M/2} |\partial^\alpha u'|, \end{aligned}$$

where, in the last line, we have applied (5.1) and (5.14). By (2.3) and a similar argument, we have

$$\begin{aligned} \sum_{j \leq M-1} \|\partial_t^j \partial^2 u(t, \cdot)\|_2 &\lesssim \sum_{j \leq M} \|\partial_t^j u'(t, \cdot)\|_2 + \frac{\varepsilon}{1+t} \sum_{j \leq M-1} \|\partial_t^j \partial^2 u(t, \cdot)\|_2 \\ &+ \sum_{\substack{|\alpha| \leq M-41 \\ |\beta| \leq M/2}} \|\partial^\alpha u'(t, \cdot) \partial^\beta u'(t, \cdot)\|_2. \end{aligned}$$

For ε sufficiently small, the second term can be bootstrapped. Thus, by combining the previous two estimates with (5.26) and using that $\sum_{j \leq M} \|\partial_t^j u'(t, \cdot)\|_2 \approx E_M^{1/2}(u)(t)$ for ε

small, we have

$$\partial_t E_M^{1/2}(u)(t) \lesssim \frac{\varepsilon}{1+t} E_M^{1/2}(u)(t) + \sum_{\substack{|\alpha| \leq M-41 \\ |\beta| \leq M/2}} \|\partial^\alpha u'(t, \cdot) \partial^\beta u'(t, \cdot)\|_2. \quad (5.27)$$

For $M \leq 40$, the last term drops out. Thus, by (1.6) and Gronwall's inequality, we have

$$\sum_{j \leq 40} \|\partial_t^j u'(t, \cdot)\|_2 \lesssim \varepsilon (1+t)^{C\varepsilon},$$

which provides the same bound for

$$\sum_{|\alpha| \leq 40} \|\partial^\alpha u'(t, \cdot)\|_2$$

by elliptic regularity and (5.14).

In order to study the case $M > 40$, we shall use the following lemma.

Lemma 5.1. Suppose that (5.1), (5.5), (5.13), and (5.14) hold. Moreover, for $\mu \leq 3$ fixed, assume that

$$\begin{aligned} & \sum_{\substack{|\alpha| + \nu \leq 100 - 8(\mu-1) \\ \nu \leq \mu-1}} \|S^\nu \partial^\alpha u'(t, \cdot)\|_2 + \sum_{\substack{|\alpha| + \nu \leq 97 - 8(\mu-1) \\ \nu \leq \mu-1}} \|\langle x \rangle^{-1/2} S^\nu \partial^\alpha u'\|_{L_t^2 L_x^2(S_t)} \\ & + \sum_{\substack{|\alpha| + \nu \leq 96 - 8(\mu-1) \\ \nu \leq \mu-1}} \|S^\nu Z^\alpha u'(t, \cdot)\|_2 + \sum_{\substack{|\alpha| + \nu \leq 94 - 8(\mu-1) \\ \nu \leq \mu-1}} \|\langle x \rangle^{-1/2} S^\nu Z^\alpha u'\|_{L_t^2 L_x^2(S_t)} \\ & \lesssim \varepsilon (1+t)^{C\varepsilon + \sigma} \end{aligned} \quad (5.28)$$

and, for $M \leq 100 - 8\mu$,

$$\begin{aligned} & \sum_{\substack{|\alpha| + \nu \leq M \\ \nu \leq \mu}} \|S^\nu \partial^\alpha u'(t, \cdot)\|_2 + \sum_{\substack{|\alpha| + \nu \leq M-3 \\ \nu \leq \mu}} \|\langle x \rangle^{-1/2} S^\nu \partial^\alpha u'\|_{L_t^2 L_x^2(S_t)} \\ & + \sum_{\substack{|\alpha| + \nu \leq M-4 \\ \nu \leq \mu}} \|S^\nu Z^\alpha u'(t, \cdot)\|_2 + \sum_{\substack{|\alpha| + \nu \leq M-6 \\ \nu \leq \mu}} \|\langle x \rangle^{-1/2} S^\nu Z^\alpha u'\|_{L_t^2 L_x^2(S_t)} \lesssim \varepsilon (1+t)^{C\varepsilon + \sigma} \end{aligned} \quad (5.29)$$

for $C, \sigma > 0$. Then there is a constant $C' > 0$ so that

$$\begin{aligned} & \sum_{\substack{|\alpha|+\nu \leq M-2 \\ \nu \leq \mu}} \|\langle \mathbf{x} \rangle^{-1/2} S^\nu \partial^\alpha \mathbf{u}'\|_{L_t^2 L_x^2(S_t)} + \sum_{\substack{|\alpha|+\nu \leq M-3 \\ \nu \leq \mu}} \|S^\nu Z^\alpha \mathbf{u}'(t, \cdot)\|_2 \\ & + \sum_{\substack{|\alpha|+\nu \leq M-5 \\ \nu \leq \mu}} \|\langle \mathbf{x} \rangle^{-1/2} S^\nu Z^\alpha \mathbf{u}'\|_{L_t^2 L_x^2(S_t)} \lesssim \varepsilon (1+t)^{C'\varepsilon + C'\sigma}. \end{aligned} \quad (5.30) \quad \square$$

Proof of Lemma 5.1. We begin by bounding the first term in the left of (5.30). By (5.1) and (5.5), it suffices to bound this term when u is replaced by $w - v$. By (2.23), we thus need to establish control of

$$\sum_{\substack{|\alpha|+\nu \leq M-2 \\ \nu \leq \mu}} \int_0^t \|S^\nu \partial^\alpha L u(s, \cdot)\|_2 ds + \sum_{\substack{|\alpha|+\nu \leq M-3 \\ \nu \leq \mu}} \|S^\mu \partial^\alpha L u\|_{L_t^2 L_x^2(S_t)}.$$

We will show the bound for the first term. The estimate for the second term follows from straightforward modifications of the argument.

For $\mu = 0$, we have, as above, that

$$\sum_{|\alpha| \leq M-2} |\partial^\alpha L u| \lesssim \sum_{|\alpha| \leq 40} |\partial^\alpha \mathbf{u}'| + \sum_{|\alpha| \leq M-1} |\partial^\alpha \mathbf{u}'| + \sum_{|\alpha| \leq M-41} |\partial^\alpha \mathbf{u}'| + \sum_{|\alpha| \leq M/2} |\partial^\alpha \mathbf{u}'|.$$

When $M < 41$, the last term vanishes and, in this case, we have that

$$\sum_{|\alpha| \leq M-2} \int_0^t \|\partial^\alpha L u(s, \cdot)\|_2 ds \lesssim \int_0^t \frac{\varepsilon}{1+s} \sum_{|\alpha| \leq M-1} \|\partial^\alpha \mathbf{u}'(s, \cdot)\|_2 ds.$$

The desired bound then follows from (5.29). When $M \geq 41$, we apply (4.1) to see that

$$\begin{aligned} & \sum_{|\alpha| \leq M-2} \int_0^t \|\partial^\alpha L u(s, \cdot)\|_2 ds \lesssim \int_0^t \frac{\varepsilon}{1+s} \sum_{|\alpha| \leq M-1} \|\partial^\alpha \mathbf{u}'(s, \cdot)\|_2 ds \\ & + \sum_{|\alpha| \leq \max(M-38, M/2)} \|\langle \mathbf{x} \rangle^{-1/2} Z^\alpha \mathbf{u}'\|_{L_t^2 L_x^2(S_t)}, \end{aligned}$$

and (5.30) again follows from (5.29).

For $\mu > 0$, we have the similar bound

$$\begin{aligned}
\sum_{\substack{|\alpha|+\nu \leq M-2 \\ \nu \leq \mu}} |S^\nu \partial^\alpha Lu| &\lesssim \sum_{|\alpha| \leq 40} |\partial^\alpha u'| \sum_{\substack{|\alpha|+\nu \leq M-1 \\ \nu \leq \mu}} |S^\nu \partial^\alpha u'| \\
&+ \sum_{\substack{|\alpha|+\nu \leq M-41 \\ \nu \leq \mu}} |S^\nu \partial^\alpha u'| \sum_{|\alpha| \leq M-1} |\partial^\alpha u'| + \sum_{\substack{|\alpha|+\nu \leq M/2 \\ \nu \leq \mu-1}} |S^\nu \partial^\alpha u'| \sum_{\substack{|\alpha|+\nu \leq M-1 \\ \nu \leq \mu-1}} |S^\nu \partial^\alpha u'|. \quad (5.31)
\end{aligned}$$

We may apply (5.14) and (4.1) to see that

$$\begin{aligned}
\sum_{\substack{|\alpha|+\nu \leq M-2 \\ \nu \leq \mu}} \int_0^t \|S^\nu \partial^\alpha Lu(s, \cdot)\|_2 ds &\lesssim \int_0^t \frac{\varepsilon}{1+s} \sum_{\substack{|\alpha|+\nu \leq M-1 \\ \nu \leq \mu}} \|S^\nu \partial^\alpha u'(s, \cdot)\|_2 ds \\
&+ \sum_{\substack{|\alpha|+\nu \leq M-39 \\ \nu \leq \mu}} \|\langle x \rangle^{-1/2} S^\nu Z^\alpha u'\|_{L_t^2 L_x^2(S_t)} \sum_{|\alpha| \leq M-1} \|\langle x \rangle^{-1/2} \partial^\alpha u'\|_{L_t^2 L_x^2(S_t)} \\
&+ \sum_{\substack{|\alpha|+\nu \leq \max(M-1, M/2+2) \\ \nu \leq \mu-1}} \|\langle x \rangle^{-1/2} S^\nu Z^\alpha u'\|_{L_t^2 L_x^2(S_t)}^2,
\end{aligned}$$

and the desired bounds follow from (5.28) and (5.29).

Once the estimate (5.30) is established for the second term in the left, the bound for the third term in the left follows from the arguments above using (2.24).

It remains to bound the second term in the left of (5.30). To this end, we shall apply (2.21). We first examine the $\mu = 0$ case. Notice that

$$\begin{aligned}
\sum_{|\alpha| \leq M-3} \|L_\gamma Z^\alpha u(t, \cdot)\|_2 &\lesssim \sum_{|\alpha| \leq 40} \|Z^\alpha u'(t, \cdot)\|_\infty \sum_{|\alpha| \leq M-3} \|Z^\alpha u'(t, \cdot)\|_2 \\
&+ \sum_{|\alpha| \leq M-43, |\beta| \leq M/2} \|Z^\alpha u'(t, \cdot) Z^\beta u'(t, \cdot)\|_2 \\
&\lesssim \frac{\varepsilon}{1+t} Y_{M-3,0}^{1/2}(t) + \sum_{|\alpha| \leq \max(M-41, M/2+2)} \|\langle x \rangle^{-1/2} Z^\alpha u'(t, \cdot)\|_2^2,
\end{aligned}$$

and moreover notice that the last term is unnecessary when $M < 43$. By (5.14), when $M < 43$, it follows from (2.21) and Gronwall's inequality that

$$\sum_{|\alpha| \leq M-3} \|Z^\alpha u'(t, \cdot)\|_2^2 \lesssim (1+t)^{C\varepsilon} \left(\varepsilon^2 + \sum_{|\alpha| \leq M-2} \|\partial^\alpha u'\|_{L_t^2 L_x^2([0,t] \times \{|x| < 1\})}^2 \right).$$

The desired bound then follows from that for the first term in the left of (5.30). When $M \geq 43$, we similarly apply (2.21), (5.14), and Gronwall's inequality to see that

$$\begin{aligned} & \sum_{|\alpha| \leq M-3} \|Z^\alpha u'(t, \cdot)\|_2^2 \\ & \lesssim (1+t)^{C\varepsilon} \left(\varepsilon^2 + \sum_{|\alpha| \leq \max(M-41, M/2+2)} \|\langle x \rangle^{-1/2} Z^\alpha u'\|_{L_t^2 L_x^2(S_t)}^2 \sup_{0 < s < t} Y_{M-3,0}^{1/2}(s) \right. \\ & \qquad \qquad \qquad \left. + \sum_{|\alpha| \leq M-2} \|\partial^\alpha u'\|_{L_t^2 L_x^2([0,t] \times \{|x| < 1\})}^2 \right). \end{aligned}$$

Thus, by additionally applying (5.29) and bootstrapping the $Y_{M-3,0}^{1/2}$ term, the estimate follows.

For $\mu > 0$, we argue in a fashion similar to (5.31). Indeed, we have

$$\begin{aligned} & \sum_{\substack{|\alpha| + \nu \leq M-3 \\ \nu \leq \mu}} \|L_\gamma S^\nu Z^\alpha u(t, \cdot)\|_2 \lesssim \frac{\varepsilon}{1+t} Y_{M-3-\mu, \mu}^{1/2}(t) \\ & + \sum_{\substack{|\alpha| + \nu \leq M-41 \\ \nu \leq \mu}} \|\langle x \rangle^{-1/2} S^\nu Z^\alpha u'(t, \cdot)\|_2 \sum_{|\alpha| \leq M-3} \|\langle x \rangle^{-1/2} Z^\alpha u'(t, \cdot)\|_2 \\ & \qquad \qquad \qquad + \sum_{\substack{|\alpha| + \nu \leq \max(M-3, M/2+1) \\ \nu \leq \mu-1}} \|\langle x \rangle^{-1/2} S^\nu Z^\alpha u'(t, \cdot)\|_2^2. \end{aligned}$$

Thus, by (1.6), (2.21), (5.14), and Gronwall's inequality, it follows that

$$\begin{aligned} & \sum_{\substack{|\alpha| + \nu \leq M-3 \\ \nu \leq \mu}} \|S^\nu Z^\alpha u'(t, \cdot)\|_2^2 \\ & \lesssim (1+t)^{C\varepsilon} \left[\varepsilon^2 + \sup_{0 < s < t} Y_{M-3-\mu, \mu}^{1/2}(s) \left(\sum_{\substack{|\alpha| + \nu \leq \max(M-3, M/2+1) \\ \nu \leq \mu-1}} \|\langle x \rangle^{-1/2} S^\nu Z^\alpha u'\|_{L_t^2 L_x^2(S_t)}^2 \right. \right. \\ & + \sum_{\substack{|\alpha| + \nu \leq M-41 \\ \nu \leq \mu}} \|\langle x \rangle^{-1/2} S^\nu Z^\alpha u'\|_{L_t^2 L_x^2(S_t)} \sum_{|\alpha| \leq M-3} \|\langle x \rangle^{-1/2} Z^\alpha u'\|_{L_t^2 L_x^2(S_t)} \left. \right. \\ & \qquad \qquad \qquad \left. + \sum_{\substack{|\alpha| + \nu \leq M-2 \\ \nu \leq \mu}} \|S^\nu \partial^\alpha u'\|_{L_t^2 L_x^2([0,t] \times \{|x| < 1\})}^2 \right]. \end{aligned}$$

Applying (5.28) and (5.29) completes the proof. \blacksquare

Reverting back to (5.27) and applying (4.1) and Gronwall's inequality, it follows that

$$E_M^{1/2}(\mathbf{u})(t) \lesssim (1+t)^{C\varepsilon} \left(\varepsilon + \sum_{|\alpha| \leq \max(M-39, M/2)} \|\langle \mathbf{x} \rangle^{-1/2} \mathbf{Z}^\alpha \mathbf{u}'\|_{L_t^2 L_x^2(S_t)}^2 \right)$$

for $40 \leq M \leq 100$. Using this in an induction argument based on the previous lemma gives

$$\sum_{j \leq 100} \|\partial_t^j \mathbf{u}'(t, \cdot)\|_2 \lesssim \varepsilon (1+t)^{C\varepsilon + \sigma},$$

and coupled with elliptic regularity, this yields (5.16). Additionally, from Lemma 5.1, we obtain

$$\begin{aligned} \sum_{|\alpha| \leq 98} \|\langle \mathbf{x} \rangle^{-1/2} \partial^\alpha \mathbf{u}'\|_{L_t^2 L_x^2(S_t)} + \sum_{|\alpha| \leq 97} \|\mathbf{Z}^\alpha \mathbf{u}'(t, \cdot)\|_2 \\ + \sum_{|\alpha| \leq 95} \|\langle \mathbf{x} \rangle^{-1/2} \mathbf{Z}^\alpha \mathbf{u}'\|_{L_t^2 L_x^2(S_t)} \lesssim \varepsilon (1+t)^{C'\varepsilon + C'\sigma}, \end{aligned} \quad (5.32)$$

which gives the $\nu = 0$ versions of (5.17) and (5.18) and will provide the hypothesis (5.28) for further applications of the lemma corresponding to a single occurrence of S .

We now approach the proofs of the estimates involving the scaling vector field S following a similar strategy. For a fixed $\mu = 1, 2, 3$, we shall assume (5.28), and the main step is to show that

$$\sum_{\substack{|\alpha| + \nu \leq 100 - 8\mu \\ \nu \leq \mu}} \|\mathbf{S}^\nu \partial^\alpha \mathbf{u}'(t, \cdot)\|_2 \lesssim \varepsilon (1+t)^{\tilde{C}(\varepsilon + \sigma)} \quad (5.33)$$

for some $\tilde{C} > 0$.

We need first to establish (2.17). Noticing that, for $M \leq 100 - 8\mu$,

$$\begin{aligned} \sum_{\substack{j + \nu \leq M \\ \nu \leq \mu}} \left(|\tilde{\mathbf{S}}^\nu \partial_t^j L_\gamma \mathbf{u}| + |[\tilde{\mathbf{S}}^\nu \partial_t^j, \gamma^{kl} \partial_k \partial_l] \mathbf{u}| \right) \\ \lesssim \sum_{|\alpha| \leq 40} |\partial^\alpha \mathbf{u}'| \sum_{\substack{j + \nu \leq M - 1 \\ \nu \leq \mu}} |\tilde{\mathbf{S}}^\nu \partial_t^j \partial^2 \mathbf{u}| + \sum_{\substack{|\alpha| + \nu \leq M - 40 \\ \nu \leq \mu}} |\mathbf{S}^\nu \partial^\alpha \mathbf{u}'| \sum_{\substack{|\alpha| + \nu \leq M \\ \nu \leq \mu - 1}} |\mathbf{S}^\nu \partial^\alpha \mathbf{u}'|, \end{aligned}$$

it follows from elliptic regularity, (4.1), and (5.14) that

$$\begin{aligned} \sum_{\substack{j+\nu \leq M \\ \nu \leq \mu}} \left(\|\tilde{S}^\nu \partial_t^j L_\gamma u(t, \cdot)\|_2 + \|\tilde{S}^\nu \partial_t^j \gamma^{kl} \partial_k \partial_l u(t, \cdot)\|_2 \right) &\lesssim \frac{\varepsilon}{1+t} \sum_{\substack{j+\nu \leq M \\ \nu \leq \mu}} \|\tilde{S}^\nu \partial_t^j u'(t, \cdot)\|_2 \\ &+ \sum_{\substack{|\alpha|+\nu \leq M-38 \\ \nu \leq \mu}} \|\langle x \rangle^{-1/2} S^\nu Z^\alpha u'(t, \cdot)\|_2^2 + \sum_{\substack{|\alpha|+\nu \leq M \\ \nu \leq \mu-1}} \|\langle x \rangle^{-1/2} S^\nu \partial^\alpha u'(t, \cdot)\|_2^2. \end{aligned}$$

Setting the sum of the last two terms above to be $H_{\nu, M-\nu}(t)$, it follows from (2.18), (5.14), and (5.28) that

$$\begin{aligned} \sum_{\substack{|\alpha|+\nu \leq M \\ \nu \leq \mu}} \|S^\mu \partial^\alpha u'(t, \cdot)\|_2 &\lesssim (1+t)^{A(\varepsilon+\sigma)} \varepsilon \\ &+ (1+t)^{A\varepsilon} \sum_{\substack{|\alpha|+\nu \leq M-38 \\ \nu \leq \mu}} \|\langle x \rangle^{-1/2} S^\nu Z^\alpha u'\|_{L_t^2 L_x^2(S_t)}^2 \\ &+ (1+t)^{A\varepsilon} \sum_{\substack{|\alpha|+\nu \leq M \\ \nu \leq \mu-1}} \int_0^t \|S^\nu \partial^\alpha u'(s, \cdot)\|_{L^2(\{|x|<1\})} ds. \end{aligned} \quad (5.34)$$

To control the last term in (5.34), we apply (5.1), (5.5), and (3.22) to see that it is

$$\lesssim \varepsilon \log(2+t)(1+t)^{A\varepsilon} + (1+t)^{A\varepsilon} \sum_{\substack{|\alpha|+\nu \leq M+3 \\ \nu \leq \mu-1}} \int_0^t \int |S^\nu Z^\alpha Lu(s, y)| \frac{dy ds}{|y|},$$

where we have also applied Sobolev's lemma to bound the second term in the right of (3.22) by the preceding one. Since the second term above is

$$\leq \sum_{\substack{|\alpha|+\nu \leq M+4 \\ \nu \leq \mu-1}} \|\langle x \rangle^{-1/2} S^\nu Z^\alpha u'\|_{L_t^2 L_x^2(S_t)}^2,$$

it follows from (5.28) that

$$\begin{aligned} \sum_{\substack{|\alpha|+\nu \leq M \\ \nu \leq \mu}} \|S^\mu \partial^\alpha u'(t, \cdot)\|_2 &\lesssim (1+t)^{A(\varepsilon+\sigma)} \varepsilon \\ &+ (1+t)^{A\varepsilon} \sum_{\substack{|\alpha|+\nu \leq M-38 \\ \nu \leq \mu}} \|\langle x \rangle^{-1/2} S^\nu Z^\alpha u'\|_{L_t^2 L_x^2(S_t)}^2. \end{aligned}$$

For $M \leq 40$, this yields (5.33). For $M \geq 40$, we again argue inductively using Lemma 5.1. Arguing as such proves

$$\begin{aligned}
& \sum_{\substack{|\alpha|+\nu \leq 100-8\mu \\ \nu \leq \mu}} \|S^\nu \partial^\alpha u'(t, \cdot)\|_2 + \sum_{\substack{|\alpha|+\nu \leq 97-8\mu \\ \nu \leq \mu}} \|\langle x \rangle^{-1/2} S^\nu \partial^\alpha u'\|_{L_t^2 L_x^2(S_t)} \\
& + \sum_{\substack{|\alpha|+\nu \leq 96-8\mu \\ \nu \leq \mu}} \|S^\nu Z^\alpha u'(t, \cdot)\|_2 + \sum_{\substack{|\alpha|+\nu \leq 94-8\mu \\ \nu \leq \mu}} \|\langle x \rangle^{-1/2} S^\nu Z^\alpha u'\|_{L_t^2 L_x^2(S_t)} \\
& \lesssim \varepsilon(1+t)^{C\varepsilon+\sigma}, \quad \mu = 0, 1, 2, 3,
\end{aligned}$$

which implies (5.17) and (5.18).

It remains only to prove (5.15). Using (3.12), it is easy to see that the left side of (5.15) is dominated by the square of the left side of (5.18), from which (5.15) follows. This completes the proof of Theorem 1.1. \blacksquare

Acknowledgments

The authors are particularly grateful to T. Sideris and C. Sogge for introducing us to this problem and for numerous related discussions and collaborations. The authors were supported by the NSF.

References

- [1] Agemi, R. "Global existence of nonlinear elastic waves." *Invent. Math.* 142 (2000): 225–250.
- [2] Agemi, R. "Weighted uniform decay estimates for solutions to isotropic elastic wave equations." *Adv. Math. Sci. Appl.* 12 (2002): 151–173.
- [3] Ikawa, M. "Decay of solutions of the wave equation in the exterior of two convex bodies." *Osaka J. Math.* 19 (1982): 459–509.
- [4] Ikawa, M. "Decay of solutions of the wave equation in the exterior of several convex bodies." *Ann. Inst. Fourier (Grenoble)* 38 (1998): 113–146.
- [5] John, F. *Formation of singularities in elastic waves. Lec. Notes in Physics*, Vol. 195, edited by P. G. Ciarlet and M. Rousseau eds., Springer, New York (1984): 190–214.
- [6] John, F. "Almost global existence of elastic waves of finite amplitude arising from small initial disturbances." *Comm. Pure Appl. Math.* 41 (1988): 615–666.
- [7] Keel, M., H. Smith, and C. D. Sogge. "Global existence for a quasilinear wave equation outside of star-shaped domains." *J. Funct. Anal.* 189 (2002): 155–226.
- [8] Keel, M., H. Smith, and C. D. Sogge. "Almost global existence for some semilinear wave equations." *J. Anal. Math.* 87 (2002): 265–279.

- [9] Keel, M., H. Smith, and C. D. Sogge. "Almost global existence for quasilinear wave equations in three space dimensions." *J. Amer. Math. Soc.* 17 (2004): 109–153.
- [10] Klainerman, S. "The null condition and global existence to nonlinear wave equations." *Lect. Appl. Math.* 23 (1986): 293–326.
- [11] Klainerman, S. and T. Sideris. "On almost global existence for nonrelativistic wave equations in 3d." *Comm. Pure Appl. Math.* 49 (1996): 307–321.
- [12] Metcalfe, J. "Elastic waves in exterior domains, Part I: Almost global Global Existence." *Int. Math. Res. Not.* (2006): 1–41.
- [13] Metcalfe, J., M. Nakamura, and C. D. Sogge. "Global existence of solutions to multiple speed systems of quasilinear wave equations in exterior domains." *Forum. Math.* 17 (2005): 133–168.
- [14] Metcalfe, J., M. Nakamura, and C. D. Sogge. "Global existence of quasilinear, nonrelativistic wave equations satisfying the null condition." *Japan. J. Math.* 31 (2005): 391–472.
- [15] Metcalfe, J. and C. D. Sogge. "Hyperbolic trapped rays and global existence of quasilinear wave equations." *Invent. Math.* 159 (2005): 75–117.
- [16] Metcalfe, J. and C. D. Sogge. "Global existence for Dirichlet-wave equations with quadratic nonlinearities in high dimensions." *Math. Ann.* 336 (2006): 391–420.
- [17] Metcalfe, J. and C. D. Sogge. "Long time existence of quasilinear wave equations exterior to star-shaped obstacles via energy methods." *SIAM J. Math. Anal.* 38 (2006): 188–209.
- [18] Metcalfe, J. and C. D. Sogge. "Global existence of null form wave equations in exterior domains." *Math. Z.* 256 (2007): 521–549.
- [19] Morawetz, C. S., J. Ralston, and W. Strauss. "Decay of solutions of the wave equations outside nontrapping obstacles." *Comm. Pure Appl. Math.* 30 (1977): 447–508.
- [20] Sideris, T. C. "The null condition and global existence of nonlinear elastic waves." *Invent. Math.* 123 (1996): 323–342.
- [21] Sideris, T. C. "Nonresonance and global existence of prestressed nonlinear elastic waves." *Ann. of Math.* 151 (2000): 849–874.
- [22] Sideris, T. C. *Nonlinear hyperbolic systems and elastodynamics, Phase Space Analysis of PDEs, Scuole Normale Superiore Pisa II* (2004): 451–485.
- [23] Sideris, T. C. and B. Thomases. "Local energy decay for solutions of multi-dimensional isotropic symmetric hyperbolic systems." *J. Hyperbolic Differ. Equ.* 3 (2006): 673–690.
- [24] Sideris, T. C. and S. Y. Tu. "Global existence for systems of nonlinear wave equations in 3D with multiple speeds." *SIAM J. Math. Anal.* 33 (2001): 477–488.
- [25] Sterbenz, J. "Angular regularity and Strichartz estimates for the wave equation with an appendix by I. Rodnianski." *Int. Math. Res. Not.* (2005): 187–231.
- [26] Xin, J. and T. Qin. "Almost global existence of Dirichlet initial-boundary value problem for nonlinear elastodynamic system outside a star-shaped domain." *Chinese Ann. Math. Ser. B.* 25 (2004): 433–454.
- [27] Yamamoto, K. "Exponential energy decay of solutions of elastic wave equations with the Dirichlet condition." *Math. Scand.* 65 (1989): 206–220.