

Lecture 7: Rigid body dynamics (and Lie algebras) $G = SO(3)$, $\mathfrak{g} := T_e G := so(3)$.

For Poisson brackets:
 (1) (\mathbb{R}^3, \times) is Poisson, (2) $(\mathfrak{g} = so(3), [\cdot, \cdot])$

Prop: $(\mathbb{R}^3, \times) \cong (so(3), [\cdot, \cdot])$.
 The isomorphism is given by $(x_1, x_2, x_3) \in \mathbb{R}^3 \xrightarrow{\varphi} A(x) = \begin{pmatrix} 0 & -x_3 & x_2 \\ x_3 & 0 & -x_1 \\ -x_2 & x_1 & 0 \end{pmatrix}$.

We are saying $(v \times w) = [\varphi(v), \varphi(w)]$
 Rank: $A(v) \cdot w = v \times w$.

in general, \mathfrak{g}^* always has a canonical Poisson bracket (from $[\cdot, \cdot]$ Lie bracket)

Ex 1. Rigid body setup: $b \in \mathfrak{B}$, $w(t) = Q(t) \cdot b$, $Q(t) \in SO(3)$.

To understand the kinetic term in Hamiltonian we need $w(t)$ and dual momentum.

$w(t) = \dot{Q}(t) \cdot b = \dot{Q}(t) \cdot Q^{-1}(t) \cdot w(t)$

EX: Rotating \mathfrak{B} along z-axis gives $T_{id} so(3) \ni \omega$
 $Q(t) = \begin{pmatrix} \cos(t) & \sin(t) & 0 \\ -\sin(t) & \cos(t) & 0 \\ 0 & 0 & 1 \end{pmatrix}$
 $Q'(t) = \begin{pmatrix} -\sin(t) & \cos(t) & 0 \\ -\cos(t) & -\sin(t) & 0 \\ 0 & 0 & 0 \end{pmatrix}$
 $Q^{-1}(t) = Q(-t) = \begin{pmatrix} \cos(t) & -\sin(t) & 0 \\ \sin(t) & \cos(t) & 0 \\ 0 & 0 & 1 \end{pmatrix}$

$Q'(t) \cdot Q^{-1}(t) = \begin{pmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \in so(3)$, i.e. antisymmetric.

Kinetic energy & Hamiltonian in a rigid body $\exists! x$ s.t. $w = x \times w$

(1) $K_b = \frac{m_b}{2} \|w(t)\|^2 = \frac{m_b}{2} \|x(t) \times w(t)\|^2 = \frac{m_b}{2} \|\dot{Q}(t) \cdot (x \times w(t))\|^2 = \dot{X} \cdot \left(\frac{1}{I} \right) \cdot \dot{X}$

This is at a point $b \in \mathfrak{B}$: $K = \int_{\mathfrak{B}} K_b \cdot \Pi \cdot X$

The dual momentum is given by $M = \Pi \cdot \dot{X} = (M_1, M_2, M_3) \in so(3)^*$

Rank: we'll assume $\Pi = \text{diag}(I_1, I_2, I_3)$. So $H(M_1, M_2, M_3) = \frac{1}{2} \left(\frac{M_1^2}{I_1} + \frac{M_2^2}{I_2} + \frac{M_3^2}{I_3} \right)$

Eq of motion: $\begin{cases} \dot{M}_1 = \{M_1, H\} \\ \dot{M}_2 = \{M_2, H\} \\ \dot{M}_3 = \{M_3, H\} \end{cases}$

e.g. $\{M_1, H\} = (M_1, M_2, M_3) \cdot \left((1, 0, 0) \times \left(\frac{M_2}{I_1}, \frac{M_2}{I_1}, \frac{M_3}{I_3} \right) \right)$
 $= (M_1, M_2, M_3) \cdot \left(0, -\frac{M_3}{I_3}, \frac{M_2}{I_1} \right) = \frac{M_2 M_3}{I_2 I_3} (I_3 - I_2)$

Ex 2. Symmetries: $M = \mathfrak{g}^* = so(3)^*$, $H(M) = \frac{1}{2} M^t \Pi^{-1} \cdot M$

(i) The energy is a conserved quantity, and the level sets of H are ellipsoids.

$H = \frac{1}{2} \left(\frac{M_1^2}{I_1} + \frac{M_2^2}{I_2} + \frac{M_3^2}{I_3} \right)$

(ii) A 2nd symmetry exists: $J = M_1^2 + M_2^2 + M_3^2$

Proof: we want $\{J, H\} = 0$.
 $\{M_1^2, H\} = 2M_1 \cdot \{M_1, H\} = 2M_1 M_2 M_3 \cdot (I_3 - I_2) / I_2 I_3$
 $\{M_2^2, H\} = 2M_2 M_3 M_1 \cdot (I_1 - I_3) / I_1 I_3$
 $\{M_3^2, H\} = 2M_3 M_1 M_2 \cdot (I_2 - I_1) / I_2 I_1$

$\{J, H\} = 2M_1 M_2 M_3 \left(\frac{I_3 - I_2}{I_2 I_3} + \frac{I_1 - I_3}{I_1 I_3} + \frac{I_2 - I_1}{I_2 I_1} \right) = 0$