MAT 271: Applied & Computational Harmonic Analysis Lecture 7: Discrete Fourier Transform (DFT)

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Outline

- Definitions
- The Reciprocity Relations
- The Vector-Matrix Notation of DFT
- 4 Pictorial View of W_N^*
- Different Definitions of DFT
- 6 References

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- The DFT can be viewed as either an approximation to the Fourier transform or an approximation to the Fourier series coefficients.
- Suppose $f \in L^2[-A/2, A/2]$, and f(x) = 0 for |x| > A/2. That is, f is a *space-limited*, square integrable function, which is a reasonable assumption in practice.
- Then, we can invoke the *dual* version of *the Sampling Theorem* in the *frequency* domain, which can also be stated in terms of Fourier coefficients of the Fourier series (Recall that *the Fourier transform of the periodic functions gives the line spectrum in the frequency domain).*
- In fact, we have the following relationship

$$\hat{f}(k/A) = \int_{-A/2}^{A/2} f(x) e^{-2\pi i kx/A} dx = \left\langle f, e^{2\pi i k \cdot /A} \right\rangle = \sqrt{A} \alpha_k, \quad k \in \mathbb{Z}. \quad (1)$$

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- In general, $f \in L^2[-A/2, A/2]$ is **not** a band-limited function; Recall the Uncertainty Principle!
- Therefore, to have a finite length vector representing the frequency samples (or the Fourier coefficients), we need to truncate the frequency information for $|\xi| > \Omega/2$ for some $\Omega > 0$.
- This is the first source of error of DFT approximation to FT/FS.
- This truncation allows us to consider only k with $|k| \le A\Omega/2$

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- We now need to approximate the integration in (1) numerically. We
 use the trapezoid rule. Here is the second source of the error of DFT.
- Let's divide the interval [-A/2,A/2] into N (positive even integer¹) subintervals of equal length of $\Delta x = A/N$. Let $x_{\ell} = \ell \Delta x$, $\ell = (-N/2):(N/2)$ be the points used in the trapezoid rule. Let $g(x) = f(x) \mathrm{e}^{-2\pi \mathrm{i} k x/A}$. Then we have

$$\hat{f}(k/A) \approx \frac{\Delta x}{2} \left\{ g(-A/2) + 2 \sum_{\ell=-N/2+1}^{N/2-1} g(x_{\ell}) + g(A/2) \right\}$$

• If we assume f(-A/2) = f(A/2) (which we can do by extending f by reflection, windowing, or zero-padding followed by redefining A), ther the above approximation is simplified:

$$\hat{f}(k/A) \approx \Delta x \sum_{\ell=-N/2}^{N/2-1} g(x_{\ell}) = \frac{A}{N} \sum_{\ell=-N/2}^{N/2-1} f(\ell A/N) e^{-2\pi i k \ell/N},$$

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 $^{^{1}}$ All the subsequent matrix representations assume this. See [2, Sec. 3.1] for N being positive odd integer as well as the other cases, e.g., different starting and ending indices.

$$F_k := \frac{1}{\sqrt{N}} \sum_{\ell=-N/2}^{N/2-1} f_\ell e^{-2\pi i k \ell/N}, \quad k = -\frac{N}{2}, \dots, \frac{N}{2} - 1.$$
 (2)

- The factor $1/\sqrt{N}$ is to make DFT a unitary transformation (i.e., ℓ^2 -norm (energy) preserving transformation, so that the Parseval & Plancherel equalities holds.)²
- We now have the following relationship

$$\hat{f}(k/A) = \sqrt{A}\alpha_k \approx \frac{A}{\sqrt{N}}F_k.$$

The N-point inverse DFT is defined, as you can imagine, as follows:

$$f_{\ell} := \frac{1}{\sqrt{N}} \sum_{k=-N/2}^{N/2-1} F_k e^{2\pi i k \ell/N}, \quad \ell = -\frac{N}{2}, \dots, \frac{N}{2} - 1.$$

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• Let $\Delta \xi$ be a sampling rate in the frequency domain, i.e., $\Delta \xi = 1/A$. Since we know $\Delta x = A/N$, and $k/A = \Omega/2$ at k = N/2 (highest frequency in consideration), we have the following fundamental relations:

$$\Delta x \Delta \xi = \frac{1}{N}, \quad A\Omega = N.$$

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 - For fixed N: $A \uparrow \Longrightarrow \Delta x \uparrow \Omega \downarrow \Delta \xi \downarrow$ (coarser space sampling leads to finer frequency sampling, but the frequency bandwidth also decreases,
 - For fixed A: $N \uparrow \Longrightarrow \Delta x \downarrow \Omega \uparrow \Delta \xi \equiv const. = 1/A$ (finer space sampling leads to the increase of the frequency bandwidth).

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- We can gain great insights by expressing DFT using the *vector-matrix notation*. To do this, we need to define a couple of things.
- Let $\omega_N := \mathrm{e}^{2\pi\mathrm{i}/N}$, i.e., the *Nth root of unity*.
- Note that $\overline{\omega}_N = \omega_N^{-1}$; $\omega_N^0 = \omega_N^N = 1$; $\omega_N^{N/2} = -1$; and $\omega_N^{k+N} = \omega_N^k$ for any $k \in \mathbb{Z}$.
- Then, define a column vector:

$$\boldsymbol{w}_N^k := \frac{1}{\sqrt{N}} \left(\omega_N^{k \cdot 0}, \omega_N^{k \cdot 1}, \dots, \omega_N^{k \cdot \frac{N}{2}}, \dots, \omega_N^{k \cdot (N-1)} \right)^{\mathsf{T}}, \quad k = 0, \dots, N-1.$$

We also define another column vector:

$$\widetilde{\boldsymbol{w}}_{N}^{k} := \frac{1}{\sqrt{N}} \left(\boldsymbol{\omega}_{N}^{k \cdot (-\frac{N}{2})}, \, \boldsymbol{\omega}_{N}^{k \cdot (-\frac{N}{2}+1)}, \dots, \boldsymbol{\omega}_{N}^{k \cdot 0}, \dots, \boldsymbol{\omega}_{N}^{k \cdot (\frac{N}{2}-1)} \right)^{\mathsf{T}}, \quad k = -\frac{N}{2}, \dots, \frac{N}{2} - 1.$$

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ullet Using the properties of ω_N listed above, one can easily show that

$$\widetilde{\boldsymbol{w}}_{N}^{k}=S_{N}\boldsymbol{w}_{N}^{k},$$

where S_N is equivalent to **fftshift** in Julia/MATLAB:

$$S_N := \begin{bmatrix} O_{\frac{N}{2}} & I_{\frac{N}{2}} \\ I_{\frac{N}{2}} & O_{\frac{N}{2}} \end{bmatrix},$$

i.e.,
$$S_N(a_1,...,a_N)^{\mathsf{T}} = \left(a_{\frac{N}{2}+1},...,a_N,a_1,...,a_{\frac{N}{2}}\right)^{\mathsf{T}}$$
.

• Note that $S_N^{\mathsf{T}} = S_N^{-1} = S_N$ if N is an *even* integer, which is our assumption here. But you have to be careful for odd integer cases. Hence, in general, it is safer to use **ifftshift** in Julia/MATLAB to undo fftshift.

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- Let $f = \left(f_{-\frac{N}{2}}, \dots, f_{\frac{N}{2}-1}\right)^{\mathsf{T}}$ be a vector of sampled points $f_{\ell} = f(\ell \Delta x)$.
- Now DFT can be written as follows

$$F_k = \left\langle \boldsymbol{f}, \widetilde{\boldsymbol{w}}_N^k \right\rangle, \quad k = -\frac{N}{2}, \dots, \frac{N}{2} - 1.$$

$$W_N := \left[oldsymbol{w}_N^0 \;\middle|\; oldsymbol{w}_N^1 \;\middle|\; \cdots \;\middle|\; oldsymbol{w}_N^{N-1}
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• Let
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• Then, the N-point DFT/IDFT can be conveniently written as

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, $f = \widetilde{W}_N F$,

where \widetilde{W}_N^* is an hermitian conjugate (transposition followed by element-wise complex conjugation) of \widetilde{W}_N , and also often written as \widetilde{W}_N^* in literature.

- In fact, $\widetilde{W}_N^* = (S_N W_N S_N^\mathsf{T})^* = S_N W_N^* S_N^\mathsf{T}.$
- We also denote $\mathscr{D}_N[f] := \widetilde{W}_N^* f$.

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Both W_N and \widetilde{W}_N are N-by-N unitary matrix. In other words, both $\{\boldsymbol{w}_N^k\}_{k=0}^{N-1}$ and $\{\widetilde{\boldsymbol{w}}_N^k\}_{k=-\frac{N}{2}}^{\frac{N}{2}-1}$ are orthonormal bases of \mathbb{C}^N .

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(Proof) See e.g., $[1,\ 2,\ 3]$. Note that from this theorem we have $W_N^4=\widetilde W_N^4=I_N.$

Theorem (McClellan-Parks [3])

The multiplicities of the eigenvalues of W_N are summarized as

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4m + 2	m+1	m+1	m	m
4m + 3	m+1	m+1	m+1	m

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Outline

- Definitions
- 2 The Reciprocity Relations
- 3 The Vector-Matrix Notation of DFT
- 4 Pictorial View of W_N^*
- 5 Different Definitions of DFT
- 6 References

Using the properties of ω_N , in particular the periodicity with period N, we have:

$$\begin{aligned} \boldsymbol{W}_{N}^{*} &= \begin{bmatrix} (\boldsymbol{w}_{N}^{0})^{*} \\ (\boldsymbol{w}_{N}^{0})^{*} \\ (\boldsymbol{w}_{N}^{0})^{*} \\ (\boldsymbol{w}_{N}^{N})^{*} \\ \vdots \\ (\boldsymbol{w}_{N}^{N/2})^{*} \\ \vdots \\ (\boldsymbol{w}_{N}^{N-1})^{*} \end{bmatrix} = \frac{1}{\sqrt{N}} \begin{bmatrix} 1 & 1 & 1 & \dots & 1 \\ 1 & \overline{\omega}_{N}^{1} & \overline{\omega}_{N}^{2} & \dots & \overline{\omega}_{N}^{N-1} \\ 1 & \overline{\omega}_{N}^{2} & \overline{\omega}_{N}^{4} & \dots & \overline{\omega}_{N}^{2(N-1)} \\ \vdots & \vdots & \vdots & & \vdots \\ 1 & \overline{\omega}_{N}^{N/2} & \overline{\omega}_{N}^{2N/2} & \dots & \overline{\omega}_{N}^{(N-1)N/2} \\ \vdots & \vdots & \vdots & & \vdots \\ 1 & \overline{\omega}_{N}^{N-1} & \overline{\omega}_{N}^{2(N-1)} & \dots & \overline{\omega}_{N}^{(N-1)(N-1)} \\ \end{bmatrix} \\ &= \frac{1}{\sqrt{N}} \begin{bmatrix} 1 & 1 & 1 & \dots & 1 \\ 1 & \omega_{N}^{-1} & \omega_{N}^{-2} & \dots & \omega_{N}^{-2(N-1)} \\ 1 & \omega_{N}^{-2} & \omega_{N}^{-4} & \dots & \omega_{N}^{-2(N-1)} \\ \vdots & \vdots & \vdots & \dots & \vdots \\ 1 & \omega_{N}^{N/2-1} & \omega_{N}^{2(N/2+1)} & \dots & \omega_{N}^{(N-1)(N/2-1)} \\ 1 & \omega_{N}^{N/2-1} & \omega_{N}^{2(N/2-1)} & \dots & \omega_{N}^{(N-1)(N/2-2)} \\ \vdots & \vdots & \ddots & \ddots & \vdots \\ 1 & \omega_{N}^{N/2-2} & \omega_{N}^{2(N/2-2)} & \dots & \omega_{N}^{N-1} \end{bmatrix} \end{aligned}$$

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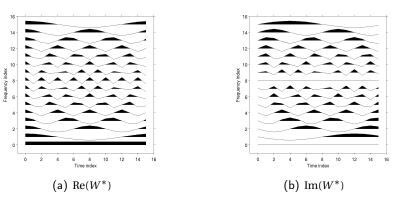
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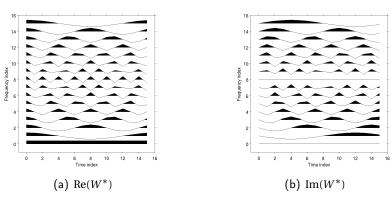
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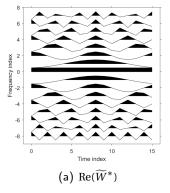


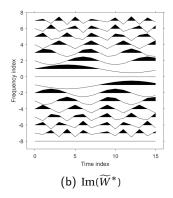
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Now, how about \widetilde{W}_N^* ?

Note the change of the locations of the basis vectors as well as symmetry $(W_N^*)^T = W_N^*, \ (\widetilde{W}_N^*)^T = \widetilde{W}_N^*.$

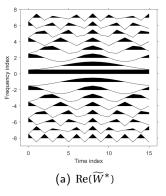
Now, how about \widetilde{W}_N^* ?

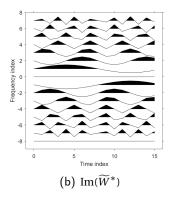




Note the change of the locations of the basis vectors as well as symmetry $(W_N^*)^{\mathsf{T}} = W_N^*$, $(\widetilde{W}_N^*)^{\mathsf{T}} = \widetilde{W}_N^*$.

Now, how about \widetilde{W}_N^* ?





Note the change of the locations of the basis vectors as well as symmetry $(W_N^*)^\mathsf{T} = W_N^*$, $(\widetilde{W}_N^*)^\mathsf{T} = \widetilde{W}_N^*$.

Outline

- Definitions
- 2 The Reciprocity Relations
- 3 The Vector-Matrix Notation of DFT
- 4 Pictorial View of W_N^*
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$$\begin{split} \text{MATLAB, Julia}^3, & \text{ R, S-Plus: } F_k = \sum_{\ell=1}^N f_\ell \, \mathrm{e}^{-2\pi \mathrm{i}(k-1)(\ell-1)/N} \, \, \text{for } k=1:N. \\ \text{Mathematica: } & F_k = \frac{1}{\sqrt{N}} \sum_{\ell=1}^N f_\ell \, \mathrm{e}^{2\pi \mathrm{i}(k-1)(\ell-1)/N} \, \, \text{for } k=1:N. \\ \text{Maple: } & F_{k+1} = \frac{1}{N} \sum_{\ell=0}^{N-1} f_\ell \, \mathrm{e}^{-2\pi \mathrm{i}k\ell/N} \, \, \text{for } k=0:(N-1). \\ \text{MathCad: } & F_k = \frac{1}{\sqrt{N}} \sum_{\ell=0}^{N-1} f_\ell \, \mathrm{e}^{2\pi \mathrm{i}k\ell/N} \, \, \text{for } k=0:(N-1). \end{split}$$

• Hence, the DFT we defined in this lecture, i.e., $F = \widetilde{W}_N^* f$, can be realized by the following MATLAB/Julia command (assuming that f is a 1D vector):

³Need to add FFTW.jl package via Pkg.add("FFTW"); using FFT

MATLAB, Julia³, R, S-Plus:
$$F_k = \sum_{\ell=1}^{N} f_{\ell} e^{-2\pi i (k-1)(\ell-1)/N}$$
 for $k = 1:N$.

Mathematica:
$$F_k = \frac{1}{\sqrt{N}} \sum_{\ell=1}^N f_\ell \, \mathrm{e}^{2\pi \mathrm{i}(k-1)(\ell-1)/N}$$
 for $k=1:N$ Maple: $F_{k+1} = \frac{1}{N} \sum_{\ell=0}^{N-1} f_\ell \, \mathrm{e}^{-2\pi \mathrm{i}k\ell/N}$ for $k=0:(N-1)$. MathCad: $F_k = \frac{1}{\sqrt{N}} \sum_{\ell=0}^{N-1} f_\ell \, \mathrm{e}^{2\pi \mathrm{i}k\ell/N}$ for $k=0:(N-1)$.

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Maple:
$$F_{k+1} = \frac{1}{N} \sum_{\ell=0}^{N-1} f_{\ell} e^{-2\pi i k \ell/N}$$
 for $k = 0 : (N-1)$.

MathCad:
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 for $k = 0$: $(N-1)$.

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• Hence, the DFT we defined in this lecture, i.e., $F = \widetilde{W}_N^* f$, can be realized by the following MATLAB/Julia command (assuming that f is a 1D vector):

F=fftshift(fft(ifftshift(f)))/sqrt(length(f));

DFT

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MATLAB, Julia³, R, S-Plus:
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$$F_k = \frac{1}{\sqrt{N}} \sum_{\ell=1}^{N} f_{\ell} e^{2\pi i (k-1)(\ell-1)/N}$$
 for $k=1:N$.

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$$F_k = \frac{1}{\sqrt{N}} \sum_{\ell=0}^{N-1} f_{\ell} e^{2\pi i k \ell/N}$$
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• Hence, the DFT we defined in this lecture, i.e., $F = \widetilde{W}_N^* f$, can be realized by the following MATLAB/Julia command (assuming that f is a 1D vector):

 $^{^3}$ Need to add FFTW.jl package via Pkg.add("FFTW"); using FFTW

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 for $k=1:N$.

Mathematica:
$$F_k = \frac{1}{\sqrt{N}} \sum_{\ell=1}^N f_\ell \, \mathrm{e}^{2\pi \mathrm{i}(k-1)(\ell-1)/N}$$
 for $k=1:N$. Maple: $F_{k+1} = \frac{1}{N} \sum_{\ell=0}^{N-1} f_\ell \, \mathrm{e}^{-2\pi \mathrm{i}k\ell/N}$ for $k=0:(N-1)$.

$$\mathsf{MathCad:}\ F_k = \frac{1}{\sqrt{N}} \sum_{\ell=0}^{N-1} f_\ell \, \mathrm{e}^{2\pi \mathrm{i} k\ell/N} \ \text{for} \ k = 0 : (N-1).$$

• Hence, the DFT we defined in this lecture, i.e., $F = \widetilde{W}_N^* f$, can be realized by the following MATLAB/Julia command (assuming that f is a 1D vector):

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Further caution:

- If an input argument to the DFT/FFT function is a matrix (or multidimensional array), then MATLAB applies DFT on each column vector for a matrix (or the first non-singleton dimension for a 3D or higher dimensional array.
- On the other hand, the DFT functions in the other packages including Julia perform the *multidimensional DFT* on the input.

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References

For more information about the DFT including higher-dimensional versions, see [2].

Also, the DFT matrix has more *profound* properties. See the challenging and deep paper by [1].

- [1] L. Auslander and R. Tolimieri, *Is computing with the finite Fourier transform pure or applied mathematics?*, Bull. Amer. Math. Soc., 1 (1979), pp. 847–897.
- [2] W. L. Briggs and V. E. Henson, The DFT: An Owner's Manual for the Discrete Fourier Transform, SIAM, Philadelphia, PA, 1995.
- [3] J. H. McClellan and T. W. Parks, *Eigenvalue and eigenvector decomposition of the discrete Fourier transform*, IEEE Trans. Audio Electacoust., AU-20 (1972), pp. 66–74.

 See also comments appeared in AU-21, pp. 65, 1973.