

MAT 271: Applied & Computational Harmonic Analysis

Supplementary Notes I by Naoki Saito

The Fourier Inversion Theorem

- The Fourier transform \mathcal{F} was defined initially on $L^1(\mathbb{R})$, a space of integrable functions, and $\mathcal{F} : L^1(\mathbb{R}) \rightarrow BC(\mathbb{R}) = C(\mathbb{R}) \cap L^\infty(\mathbb{R})$.
- However, \hat{f} , the Fourier transform of $f \in L^1$, may not be in L^1 .
An example: $f(x) = \chi_{(-\frac{1}{2}, \frac{1}{2})}(x) \Rightarrow \hat{f}(\xi) = \text{sinc}(\xi) = \frac{\sin \pi \xi}{\pi \xi} \notin L^1$.
- **The Inverse Fourier Transform:** For $f \in L^1$, $\check{f}(x) = \int_{-\infty}^{\infty} f(\xi) e^{2\pi i \xi x} d\xi$.
- **[The Fourier Inversion Theorem]** If both f and \hat{f} are in L^1 , then $(\hat{\check{f}}) = (\check{\hat{f}}) = f$ almost everywhere.
- There are many functions in L^1 whose Fourier transforms are also in L^1 ; one needs only a little *smoothness* of f for necessary *decay* of \hat{f} as $|\xi| \rightarrow \infty$.
An example: If $f \in C^2(\mathbb{R})$, f' and f'' are both in L^1 , then $\mathcal{F}\{f''\}(\xi) = -(2\pi\xi)^2 \hat{f}(\xi) \in BC(\mathbb{R})$. This boundedness implies that $|\hat{f}(\xi)| \leq C/(1 + \xi^2)$. This, in turn, implies that $\hat{f} \in L^1$.

The Fourier Transforms on L^2

- The previous remark leads to the L^2 theory of the Fourier transforms. In general, simply assuming $f \in L^2$ is not enough; $\int_{-\infty}^{\infty} f(x) e^{-2\pi i \xi x} dx$ may not converge.
An example: $f(x) = \text{sinc}(x) = \frac{\sin \pi x}{\pi x} \in L^2$, but not in L^1 .
- We will overcome this problem as follows. Define a subspace of L^1 , $\mathcal{X} \triangleq \{f \in L^1 \mid \hat{f} \in L^1\}$. We first note that for such functions, we can have the Parseval equality: $\langle f, g \rangle = \langle \hat{f}, \hat{g} \rangle$ as well as the Plancherel equality. Also, for any $f \in \mathcal{X}$, $f, \hat{f} \in BC(\mathbb{R})$ as the remark after the Fourier inversion theorem. This implies that both f and \hat{f} are also in L^2 ; i.e., $\mathcal{X} \subset L^2$ (because $f \in L^1 \cap BC$ implies $f \in L^2$ thanks to the theorem: $L^p \cap L^r \subset L^q$ for $0 < p < q < r \leq \infty$, which in turn can be proved by Hölder's inequality). Now, the point is that \mathcal{X} is also *dense* in L^2 .
- We can proceed as follows: for any $f \in L^2$, we can find a sequence $\{f_n\} \subset \mathcal{X}$ such that $\|f_n - f\|_2 \rightarrow 0$ as $n \rightarrow \infty$. $\{f_n\} \subset \mathcal{X}$ means that $\{\hat{f}_n\} \subset \mathcal{X}$. Now using the Plancherel equality to this sequence, we can see $\|\hat{f}_n - \hat{f}_m\|_2 = \|f_n - f_m\|_2 \rightarrow 0$ as $m, n \rightarrow \infty$. In other words, $\{\hat{f}_n\}$ is a *Cauchy sequence* in L^2 . Since L^2 is *complete*, there exists the limit of \hat{f}_n in L^2 , and we *define* this limit as \hat{f} , the Fourier transform of $f \in L^2$.
- **[The Plancherel Theorem]** For any $f, g \in L^2$, $\langle f, g \rangle = \langle \hat{f}, \hat{g} \rangle$ and $\|f\|_2 = \|\hat{f}\|_2$.
- Finally, we can use all these facts for computing the Fourier transform of L^2 functions as follows: Suppose we set $\phi(x) = \hat{f}(x)$ where $\hat{f} \in L^2$. Then, $\hat{\phi}(\xi) = f(-\xi)$. An example: $\phi(x) = \text{sinc}(x) \in L^2$. Then, $\hat{\phi}(\xi) = \chi_{(-\frac{1}{2}, \frac{1}{2})}(\xi)$.