

Lecture 13: Wavelet Transforms

Note Title

* Problems of STFT & Gabor systems

- If a window is too large (wide), then cannot localize around sharp transitions in an input signal.
- If a window is too small (narrow), then cannot detect low freq. oscillations.
- The Balian-Low Thm: \nexists "nice" Gabor ONBs

* Key idea of wavelets:

Use translations and dilations of a single fun to analyze a given signal at different resolutions.

Def. A wavelet is a fcn $\psi \in L^2(\mathbb{R})$ s.t.

often
called
a
"mother"
wavelet

$$\cdot \int_{-\infty}^{\infty} \psi(x) dx = 0;$$

• Normalized to have $\|\psi\| = 1$; and

• Centered around $x = 0$.

Let's generate a family of TF atoms:

$$\psi_{a,b}(x) = \frac{1}{\sqrt{a}} \psi\left(\frac{x-b}{a}\right)$$

$$= \tau_b \delta_a \psi$$

$$a > 0$$

$$b \in \mathbb{R}$$

Note $\|\psi_{a,b}\|_2 = 1$.

The **wavelet transform** of $f \in L^2(\mathbb{R})$ is defined as

$$Wf(a, b) = W_\psi f(a, b) := \langle f, \psi_{a,b} \rangle$$

often called the "continuous" wavelet

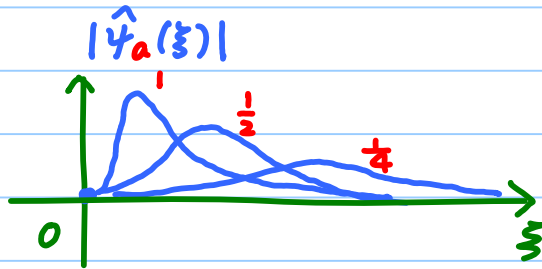
$$= \int_{-\infty}^{\infty} f(x) \frac{1}{\sqrt{a}} \overline{\psi\left(\frac{x-b}{a}\right)} dx$$

Can be viewed as a linear filtering:

transf. (CWT) $\int_{-\infty}^{\infty} f(x) \frac{1}{\sqrt{a}} \overline{\psi\left(\frac{x-b}{a}\right)} dx = f * \tilde{\psi}_a(b)$

$$\tilde{\psi}_a(x) := \frac{1}{\sqrt{a}} \psi\left(\frac{-x}{a}\right) \xrightarrow{\mathcal{F}} \hat{\tilde{\psi}}_a(\xi) = \sqrt{a} \overline{\hat{\psi}(a\xi)} = \delta \frac{1}{a} \hat{\psi}(\xi)$$

mirror of $\psi_a(x)$



Types of wavelets :

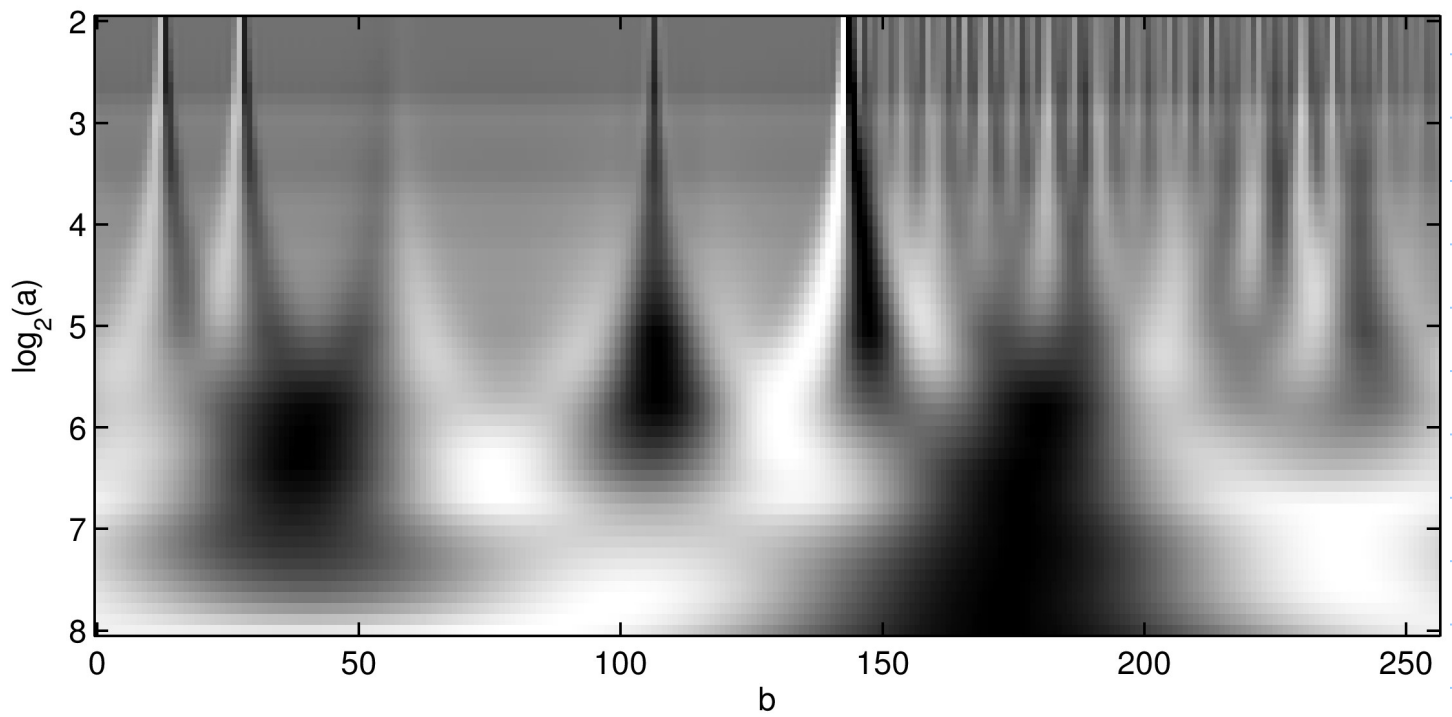
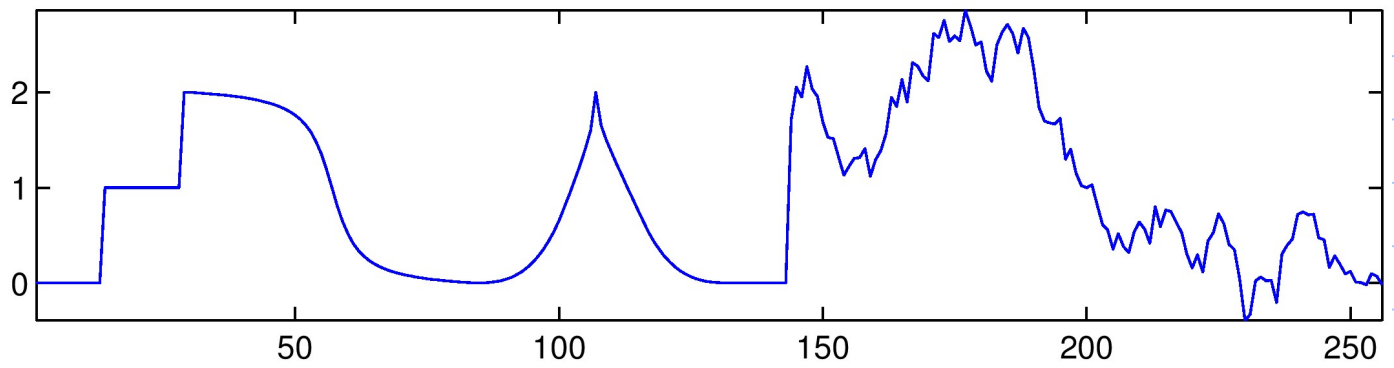
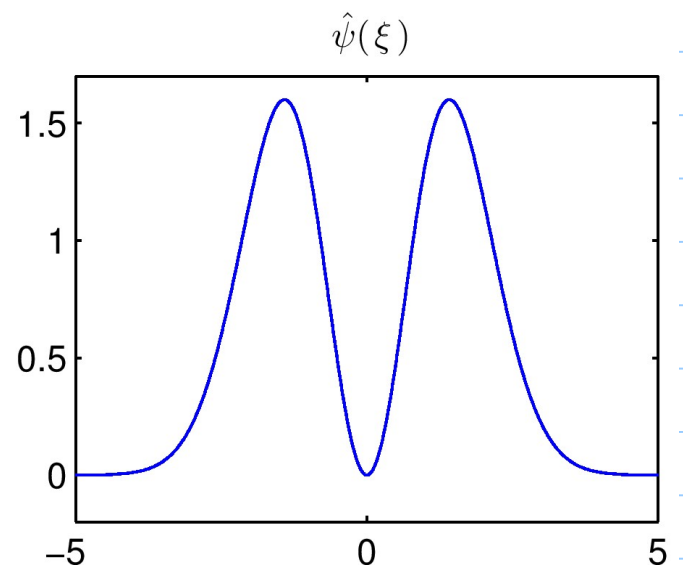
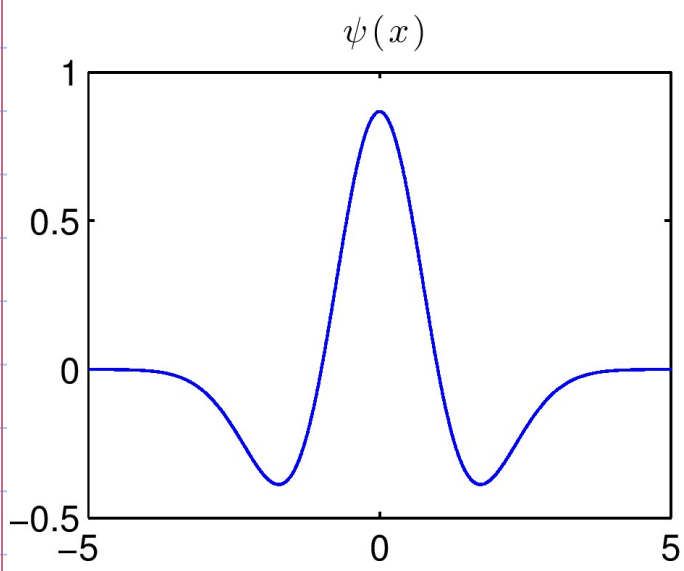
- Real wavelets \rightarrow Good for **edges**
 - Analytic (or complex) wavelets \rightarrow Can detect **phases**
- For the time being, let's focus on real wavelets.

Example : Mexican hat fcn or a.k.a. Laplacian of Gaussian (LOG)

$$\begin{cases} \psi(x) = \frac{2}{\pi^{\frac{1}{4}} \sqrt{3}\sigma} \left(1 - \frac{x^2}{\sigma^2}\right) e^{-x^2/2\sigma^2} \\ \hat{\psi}(\xi) = 8 \sqrt{\frac{2}{3}} \pi^{\frac{9}{4}} \sigma^{\frac{5}{2}} \xi^2 e^{-2\pi^2 \sigma^2 \xi^2} \end{cases}$$

$$\hat{\psi}(0) = 0, \quad \hat{\psi}(\xi) \sim \xi^2 \text{ around } \xi = 0$$

often called a **pseudo differential op.** \rightarrow approx. to $\frac{d^2}{dx^2}$



Inverse Wavelet Transform

1964

1984

Thm (Calderón - Grossmann - Morlet)

Let $\psi \in L^2(\mathbb{R})$, $\psi \in \mathbb{R}$ s.t.

$$C_\psi := \int_0^\infty \frac{|\hat{\psi}(\xi)|^2}{\xi} d\xi < +\infty$$

Then any $f \in L^2(\mathbb{R})$ satisfies

$$(*) \quad f(x) = \frac{1}{C_\psi} \int_0^\infty \int_{-\infty}^{+\infty} Wf(a, b) \psi_{a,b}(x) db \frac{da}{a^2}$$

$$\text{and } \|f\|_2^2 = \frac{1}{C_\psi} \int_0^\infty \int_{-\infty}^\infty |Wf(a, b)|^2 db \frac{da}{a^2}.$$

$$(Pf) \quad Wf(a, b) = f * \tilde{\psi}_a(b)$$

$$\text{RHS of } (*) = \frac{1}{C_\psi} \int_0^\infty (Wf(a, \cdot) * \psi_{a, \cdot})(x) \frac{da}{a^2}$$

$$= \frac{1}{C_\psi} \int_0^\infty (f * \tilde{\psi}_a * \psi_a)(x) \frac{da}{a^2}$$

\mathcal{F}

$$= \frac{1}{C_\psi} \int_0^\infty \hat{f}(\xi) \sqrt{a} \overline{\hat{\psi}(a\xi)} \sqrt{a} \hat{\psi}(a\xi) \frac{da}{a^2}$$

$$= \frac{\hat{f}(\xi)}{C_\psi} \int_0^\infty \frac{|\hat{\psi}(a\xi)|^2}{a} da$$

$a\xi = \eta$

$$= \frac{\hat{f}(\xi)}{C_\psi} \int_0^\infty \frac{|\hat{\psi}(\eta)|^2}{\eta} d\eta = \hat{f}(\xi) \quad \quad \quad = C_\psi$$

$C_\psi < +\infty$ is called the **admissibility condition**,
 (*) is called **Calderón's reproducing formula**.

$$f(x) = \frac{1}{C_\psi} \int_0^\infty f * \tilde{\psi}_a * \psi_a(x) \frac{da}{a^2}$$

↳ also called the **resolution of identity**.

To guarantee $C_\psi < \infty$, we need

$$\hat{\psi}(0) = 0 \iff \int_{-\infty}^{\infty} \psi(x) dx = 0$$

So, ψ must be oscillatory
with \pm values

also need decay on ψ

e.g., $\int_{-\infty}^{\infty} (1+|x|) \psi(x) dx < \infty$.

★ Reproducing Kernel

CWT = a **redundant** representation

$$(**) \quad Wf(a,b) = \int_{-\infty}^{\infty} \left(\frac{1}{C_\psi} \int_0^{\infty} \int_{-\infty}^{\infty} Wf(a',b') \psi_{a',b'}(x) db' \frac{da'}{a'^2} \right) \overline{\psi_{a,b}(x)} dx$$

$$= \frac{1}{C_\psi} \int_0^{\infty} \int_{-\infty}^{\infty} K(a,a',b,b') Wf(a',b') db' \frac{da'}{a'^2}$$

where $K(a,a',b,b') := \langle \psi_{a,b}, \psi_{a',b'} \rangle$

↳ measuring the **correlation** between $\psi_{a,b}$ & $\psi_{a',b'}$

If $K(a,a',b,b') = \delta(a-a') \delta(b-b')$

then **no redundancy!**

Prop. A function $\Phi(a,b) \in L^2(\mathbb{R}_+ \times \mathbb{R})$
is a wavelet transform of some $f \in L^2(\mathbb{R})$
⇔ $\Phi(a,b)$ satisfies (**).

* Scaling Function ("Father" Wavelet)

Reconstruction formula requires all values of scale $0 < a < +\infty$

If we only know $Wf(a, b)$ for $a < a_0$, then we need complementary info.

for $a > a_0$ provided by the **scaling function** (**father wavelet**) $\phi(x)$ s.t.

$$|\hat{\phi}(\xi)|^2 := \int_1^\infty |\hat{\psi}(a\xi)|^2 \frac{da}{a}$$
$$= \int_\xi^\infty \frac{|\hat{\psi}(\eta)|^2}{\eta} d\eta$$

The phase of ϕ can be arbitrary chosen.

• $\lim_{\xi \rightarrow 0} |\hat{\phi}(\xi)|^2 = C_\psi$

• $\|\phi\|_2 = 1 \leftarrow$ Exercise, use the def.

So, the low freq. approx. of f at scale a can be written as

$$L f(a, x) := \langle f, \underbrace{\phi_a}_{\equiv \phi_a} \rangle = f * \tilde{\phi}_a(x)$$

$$\Rightarrow f(x) = \frac{1}{C_\psi} \int_0^{a_0} (Wf(a, \cdot) * \psi_a)(x) \frac{da}{a^2}$$
$$+ \underbrace{\frac{1}{C_\psi a_0} (L f(a_0, \cdot) * \phi_{a_0})(x)}_{\text{Complementary info}}$$

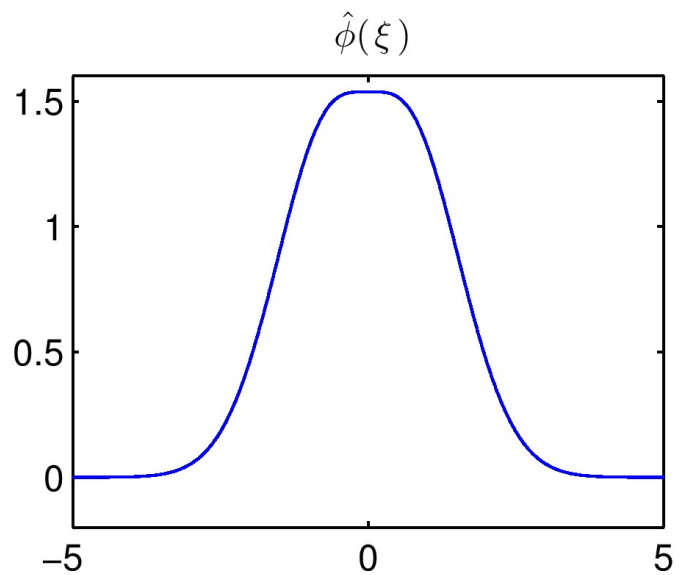
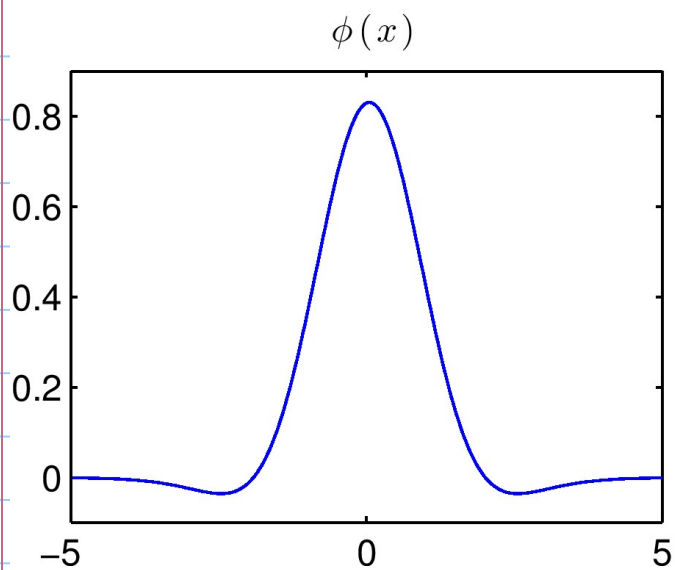
Ex. $\psi(x) = \frac{2}{\pi^{1/4} \sqrt{3\sigma}} \left(1 - \frac{x^2}{\sigma^2}\right) e^{-x^2/2\sigma^2}$

$$\hat{\psi}(\xi) = 8 \sqrt{\frac{2}{3}} \pi^{3/4} \sigma^{5/2} \xi^2 e^{-2\pi^2 \sigma^2 \xi^2}$$

$$\Rightarrow |\hat{\phi}(\xi)|^2 = \frac{4\sigma}{3\sqrt{\pi}} (1 + 4\pi^2 \sigma^2 \xi^2) e^{-4\pi^2 \sigma^2 \xi^2}$$

$$\Rightarrow \hat{\phi}(\xi) = 2 \sqrt{\frac{\sigma}{3\pi}} \sqrt{1 + 4\pi^2 \sigma^2 \xi^2} e^{-2\pi^2 \sigma^2 \xi^2}$$

↳ choose a simple phase factor. //

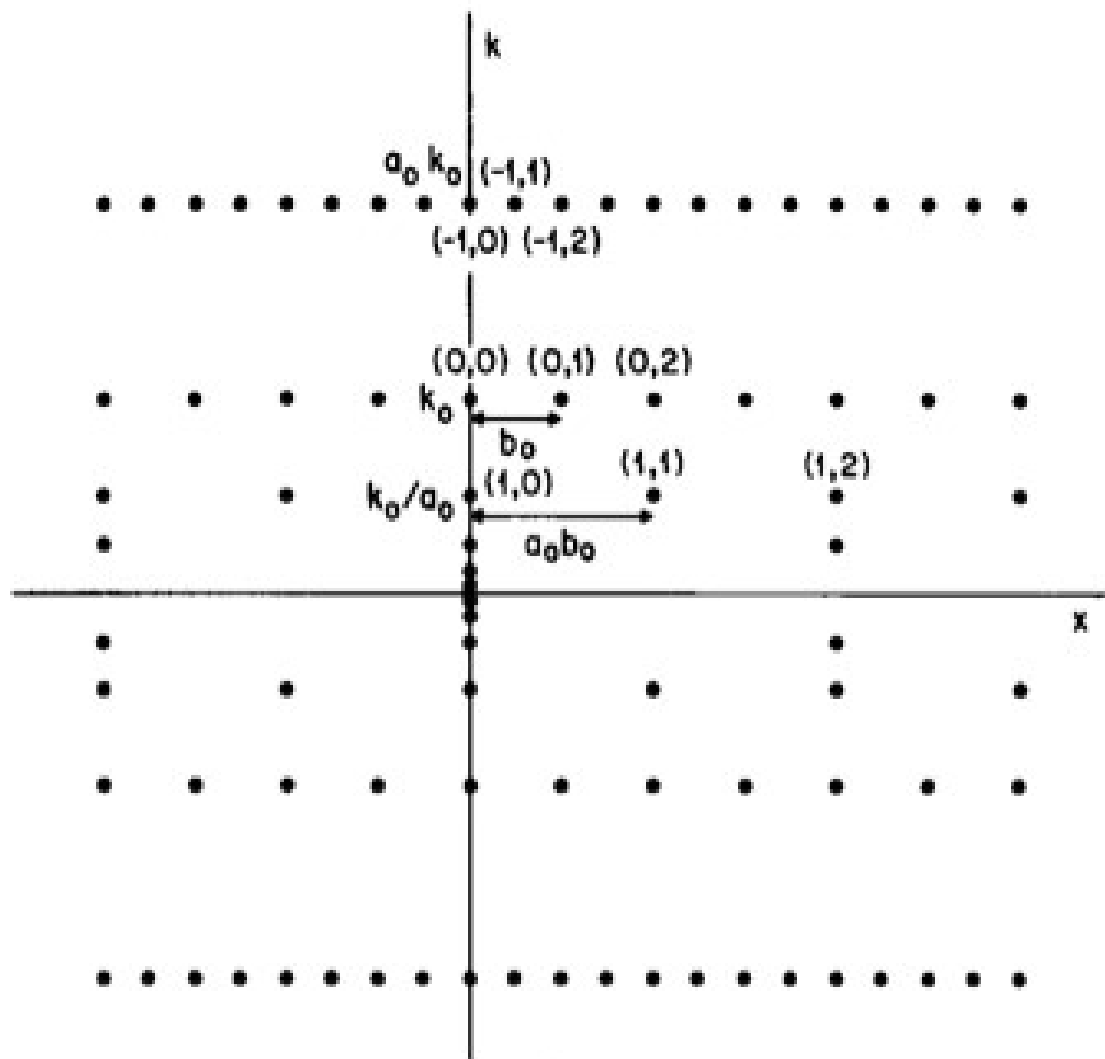


★ Discrete Wavelet Transforms

How to **sample** $Wf(a,b)$??

⇒ Another great insight by J. Morlet
"regular hyperbolic grid"

$$(a, b) = (a_0^m, n a_0^m b_0), \quad m, n \in \mathbb{Z}$$



Thm (Regular sampling thm, Daubechies'90)

Let ψ be a real-valued L^2 -function.

For fixed a_0, b_0 , define

$$\psi_{m,n}(x) := a_0^{-m/2} \psi(a_0^{-m}x - nb_0), \quad m, n \in \mathbb{Z}$$
$$= \frac{1}{\sqrt{a_0^m}} \psi\left(\frac{x - na_0^m b_0}{a_0^m}\right)$$

(1) If $\{\psi_{m,n}\}_{(m,n) \in \mathbb{Z}^2}$ is a **frame** of $L^2(\mathbb{R})$ with the frame bounds A, B , then we must have

$$A \leq \frac{1}{b_0} \sum_{-\infty}^{\infty} |\hat{\psi}(a_0^m \xi)|^2 \leq B \quad \text{for } \xi \in \mathbb{R} \text{ a.e.}$$

In particular, ψ satisfies the admissibility cond.

$$C_\psi = \int_0^\infty |\hat{\psi}(\xi)|^2 \frac{d\xi}{\xi} < +\infty$$

(2) If, for some $\epsilon > 0$, ψ satisfies

$|x|^{\frac{1}{2}+\epsilon} \psi \in L^2$, $|\xi|^\epsilon \hat{\psi} \in L^2$ and $\int \psi(x) dx = 0$, then ψ satisfies:

$$(*) \left\{ \begin{array}{l} \text{ess inf } \sum_{m \in \mathbb{Z}} |\hat{\psi}(a_0^m \xi)|^2 > 0 \\ \text{ess sup } \sum_{m \in \mathbb{Z}} |\hat{\psi}(a_0^m \xi)|^2 < +\infty \end{array} \right\} \text{ for any } a_0 \text{ close enough to } 1.$$

(i.e., $\exists \alpha = \alpha(\psi) > 1$ s.t. $(*)$ is satisfied $\forall a_0 \in (1, \alpha)$.)

Moreover, if b_0 is close enough to 0 (i.e., $\exists \beta = \beta(a_0, \psi)$ s.t. $(*)$ is satisfied $\forall b_0 \in (0, \beta)$),

then $\{\psi_{m,n}\}_{(m,n) \in \mathbb{Z}^2}$ constitute a **frame**!

EX. $\psi(x)$ = the Mexican hat fcn
 $a_0 = 2, b_0 = 1/4$.

$\Rightarrow \{\psi_{m,n}\}_{(m,n) \in \mathbb{Z}^2}$ forms a frame
(called a **wavelet frame**)

$A = 13.09, B = 14.18$ i.e., almost **tight**!

Dual Frame: Wavelet frame operator U
commutes with dilations $S_{a_0^m}$,

but **not** with translations $T_{na_0^m b_0}$.

\Rightarrow dual frame $\left\{ (U^*U)^{-1} T_{na_0^m b_0} S_{a_0^m} \psi \right\}_{(m,n)}$
is in general **not** a wavelet system $\in \mathbb{Z}^2$
(unlike the Gabor frame).