# MAT 280: Harmonic Analysis on Graphs \& Networks Lecture 10: Distances on Graphs 

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## Outline

(1) Need of a Variety of Graph Distances
(2) The (Shortest) Path Distance
(3) Resistance Distance

4 Commute-Time Distance

## Outline

(1) Need of a Variety of Graph Distances

## (2) The (Shortest) Path Distance

(3) Resistance Distance

## (4) Commute-Time Distance

## Need of a Variety of Graph Distances

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- Let $w(e)=w_{e} \geq 0$ be an edge weight of the edge $e \in E(G)$, which we also write $a_{i j}$, i.e., $(i, j)$ th entry of the weighted adjacent matrix $A$ if $e=\left(v_{i}, v_{j}\right)$.
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- Such applications include: search engines; data mining; social network analysis; pattern recognition; image processing, ...


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- If $G$ is directed, then $d_{\text {path }}(u, v) \neq d_{\text {path }}(v, u)$ in general.
- Efficient algorithms to compute $d_{\text {path }}(\cdot, \cdot)$ exist, e.g., the $A^{*}$ algorithm of Hart-Nilsson-Raphael (1968), etc. Yet, in general, its complexity is at least polynomial time w.r.t. $n . \Rightarrow$ See the excellent Wikipedia page!


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- Hence, it is reasonable that the "distance" between $v_{i}$ and $v_{j}$ should be smaller in the left than in the right.


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- Then, the voltage (potential) difference is the effective resistance between $\nu_{i}$ and $v_{j}$ by Ohm's law $(V=I \cdot R)$, which is called the resistance distance between $\nu_{i}$ and $\nu_{j}$ and written as $r\left(\nu_{i}, v_{j}\right)=r_{i j}$.


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a_{i j}= \begin{cases}1 / r_{i j} & \text { if } v_{i} \sim v_{j} \text { and } r_{i j} \nsupseteq 0 ; \\ 0 & \text { otherwise. }\end{cases}
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- But, how can we compute $\left\{r_{i j}\right\}$ if $A$ is given? Note that $r_{i j}$ should be defined even if $\nu_{i} \nsim \nu_{j}$.


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- The amazing fact is the following relationship between the resistance distance and the (Moore-Penrose) pseudoinverse $L^{\dagger}(G)$ of the unnormalized graph Laplacian $L(G)$ !

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\begin{aligned}
r_{i j} & =\left(L^{\dagger}\right)_{i i}+\left(L^{\dagger}\right)_{j j}-\left(L^{\dagger}\right)_{i j}-\left(L^{\dagger}\right)_{j i} \\
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\boldsymbol{y}_{i}=\left(I-\mathbf{1}_{n} \mathbf{1}_{n}^{\top} / n\right) \boldsymbol{e}_{i}\left[\text { Note } \operatorname{null}(L(G))=\operatorname{span}\left\{\mathbf{1}_{n}\right\} \text { if } G\right. \text { is connected.] }
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(2) Find a solution $\widehat{\ell_{i}}$ of $L \boldsymbol{\ell}=\boldsymbol{y}_{i}$ where the Cholesky factorization of $L$ should be utilized.
(3) Project the result on the row space of $L$ (which is the same as the column space thanks to $L^{\top}=L$ ) to compute the $i$ th column vector $\boldsymbol{\ell}_{i}^{\dagger}=\left(I-\mathbf{1}_{n} \mathbf{1}_{n}^{\top} / n\right) \widehat{\boldsymbol{\ell}_{i}}$.


## Intermezzo: The (Moore-Penrose) Pseudoinverse

## Definition

The pseudoinverse $A^{\dagger} \in \mathbb{C}^{n \times m}$ of a general matrix $A \in \mathbb{C}^{m \times n}$ is defined to be the unique matrix $X \in \mathbb{C}^{n \times m}$ that satisfies the following Moore-Penrose conditions: i) $A X A=A$; ii) $X A X=X$; iii) $(A X)^{*}=A X$; iv) $(X A)^{*}=X A$, where $A^{*}$ is the Hermitian transposition of $A$.

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Using the SVD of $A=U \Sigma V^{*}$ where $U \in \mathbb{C}^{m \times m}, V \in \mathbb{C}^{n \times n}$ are unitary and $\Sigma=\operatorname{diag}\left(\sigma_{1}, \cdots, \sigma_{r}, 0, \cdots, 0\right) \in \mathbb{R}^{m \times n}, r=\operatorname{rank}(A) \leq \min (m, n), A^{\dagger}$ can be expressed as

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A^{\dagger}=V \Sigma^{\dagger} U^{*}, \quad \Sigma^{\dagger}:=\operatorname{diag}\left(1 / \sigma_{1}, \cdots, 1 / \sigma_{r}, 0, \cdots, 0\right) \in \mathbb{R}^{n \times m}
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Note that $A A^{\dagger}$ and $A^{\dagger} A$ are the orthogonal projectors onto range $(A)$ and range $\left(A^{*}\right)$, respectively.

## Intermezzo: The EP Matrices

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If $A \in \mathbb{C}^{m \times m}$ commutes with its pseudoinverse, i.e., $A^{\dagger} A=A A^{\dagger}$, then it is called an EP matrix (EP for Equal Projection).

The properties of an EP matrix:

- If $\left(\lambda_{i} \neq 0, \boldsymbol{\phi}_{i}\right)$ is an eigenpair of $A$, then $\left(1 / \lambda_{i}, \boldsymbol{\phi}_{i}\right)$ is an eigenpair of $A^{\dagger}$.


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- If $\left(\lambda_{i}=0, \boldsymbol{\phi}_{i}\right)$ is an eigenpair of $A$, then $\left(\lambda_{i}=0, \boldsymbol{\phi}_{i}\right)$ is also an eigenpair of $A^{\dagger}$.


## Properties of the Pseudoinverse $L^{\dagger}(G)$

- $\left(L^{\dagger}\right)^{\top}=L^{\dagger}$
- $\operatorname{rank}\left(L^{\dagger}\right)=n-1$
- $\operatorname{null}(L)=\operatorname{null}\left(L^{\dagger}\right)=\operatorname{span}\left\{\mathbf{1}_{n}\right\}$
- $L^{\dagger}$ is doubly centered just like $L$, i.e., its column sum and row sum are zeros.
- $L^{\dagger}$ is positive semidefinite just like $L$ is.
- $L^{\dagger}=\left(L-\mathbf{1}_{n} \mathbf{1}_{n}^{\top} / n\right)^{-1}+\mathbf{1}_{n} \mathbf{1}_{n}^{\top} / n$.
- Since $L$ represents the local properties of $G$ (e.g., connectivities, etc.), $L^{\dagger}$ represents the global properties of $G \Longrightarrow$ What is the relationship between $L^{\dagger}$ and the integral operator commuting with $L$ in Lecture 2?


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## Commute-Time Distance

- Is quite similar to the resistance distance.
- Is intimately related to random walks on G.
- The commute time between $v_{i}$ and $v_{j}$ is the expected time it takes - The commute time $c\left(v_{i}, v_{j}\right)$ is intimately related to the resistance distance $r\left(v_{i}, v_{j}\right)$


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- The commute time $c\left(v_{i}, v_{j}\right)$ is intimately related to the resistance distance $r\left(\nu_{i}, v_{j}\right)$ :

$$
c\left(v_{i}, v_{j}\right)=c_{i j}=\operatorname{vol}(V(G)) \cdot r\left(v_{i}, v_{j}\right)=\operatorname{vol}(V(G)) \cdot r_{i j}
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## Intermezzo 2: Random Walks on Graphs

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- Since $G$ is connected, the Markov chain is irreducible, i.e., every state can be reached from any other state.
- Let $\pi(t)=\left[\pi_{1}(t), \ldots, \pi_{n}(t)\right]^{\top}$ where $\pi_{i}(t):=\operatorname{Pr}\left(s(t)=v_{i}\right)$, and let $P=\left(p_{i j}\right)=D^{-1} A$ be the transition matrix.


## Intermezzo 2: Random Walks on Graphs ...

- Then, the evolution of the Markov chain is characterized by

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\boldsymbol{\pi}(t+1)=P^{\top} \boldsymbol{\pi}(t) \Leftrightarrow \boldsymbol{\pi}^{\top}(t+1)=\boldsymbol{\pi}^{\top}(t) P .
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Proof: Since $L=D-A=D(I-P), P^{\top}=I-L D^{-1}$. Now, we have
$P^{\top} D \mathbf{1}_{n}=D \mathbf{1}_{n}-L \mathbf{1}_{n}=D \mathbf{1}_{n}$. Hence, $D \mathbf{1}_{n}=\left[d_{1}, \ldots, d_{n}\right]^{\top}$ is the eigenvector of $P^{\top}$ corresponding to the eigenvalue 1 modulo normalization constants. To make it as a probability distribution over $V$, we need to normalize it by $\operatorname{vol}(V)=\sum_{j} d_{j}$, i.e., $\boldsymbol{\pi}^{\infty}=D \mathbf{1}_{n} / \operatorname{vol}(V)$.

## Intermezzo 2: Random Walks on Graphs ...

There is a formal equivalence between Ncut and transition probabilities of the random walk:

## Proposition (Meila and Shi (2001))

Let $G$ be connected and non bipartite. Assume that we run the random walk $s(t)$ starting with $s(0)$ in the stationary distribution $\boldsymbol{\pi}^{\infty}$. For disjoint subsets $X, Y \subset V$, denote $\operatorname{Pr}(Y \mid X):=\operatorname{Pr}(s(1) \in Y \mid s(0) \in X)$. Then

$$
\operatorname{Ncut}\left(X, X^{c}\right)=\operatorname{Pr}\left(X^{c} \mid X\right)+\operatorname{Pr}\left(X \mid X^{c}\right) .
$$

## Intermezzo 2: Random Walks on Graphs ...

Proof: First of all, observe that

$$
\begin{aligned}
\operatorname{Pr}(s(0) \in X, s(1) \in Y) & =\sum_{x \in X, y \in Y} \operatorname{Pr}(s(0)=x, s(1)=y)=\sum_{x \in X, y \in Y} \pi_{x}^{\infty} p_{x y} \\
& =\sum_{x \in X, y \in Y} \frac{d_{x}}{\operatorname{vol}(V)} \cdot \frac{a_{x y}}{d_{x}}=\frac{1}{\operatorname{vol}(V)} \sum_{x \in X, y \in Y} a_{x y} .
\end{aligned}
$$

Using this, we obtain

$$
\begin{aligned}
\operatorname{Pr}(s(1) \in Y \mid s(0) \in X) & =\frac{\operatorname{Pr}(s(0) \in X, s(1) \in Y)}{\operatorname{Pr}(s(0) \in X)} \\
& =\frac{1}{\operatorname{vol}(V)}\left(\sum_{x \in X, y \in Y} a_{x y}\right) \cdot\left(\frac{\operatorname{vol}(X)}{\operatorname{vol}(V)}\right)^{-1} \\
& =\frac{\sum_{x \in X, y \in Y} a_{x y}}{\operatorname{vol}(X)} .
\end{aligned}
$$

Now, it is clear that the proposition holds.

## Average First-Passage Time/Cost

- The average first-passage time $m\left(v_{k} \mid v_{i}\right)=m(k \mid i)$ is defined as the average number of steps that a random walker, starting in state $v_{i} \neq v_{k}$, will take to enter state $v_{k}$ for the first time. More precisely, we define the minimum time until hitting state $\nu_{k}$, when staring from state $v_{i}$, as $T_{i k}:=\min \left(t \geq 0 \mid s(t)=v_{k} ; s(0)=v_{i}\right)$ for one realization of the stochastic process.


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- The random walker may pass through $\nu_{k}$ repeatedly; the minimum time corresponds to the first passage.
- The average first-passage time is the expected first-passage time: $m(k \mid i):=E\left[T_{i k} \mid s(0)=v_{i}\right]$.
- Similarly, the average first-passage cost $o\left(v_{k} \mid v_{i}\right)=o(k \mid i)$ is the average cost incurred by the random walker starting from state $v_{i}$ to reach state $v_{k}$ for the first time. Let $\gamma(j \mid i)$ be the cost of transition from state $v_{i}$ to $v_{j}$. For example, $\gamma(j \mid i)=1 / a_{i j}$ (if $i \sim j$ ) is a possibility.


## Average First-Passage Time \& Average Commute Time

- The recurrence relations computing $m(k \mid i)$ and $o(k \mid i)$ can easily be obtained by first-step analysis:

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\begin{gathered}
m(k \mid i)= \begin{cases}0 & \text { if } i=k ; \\
1+\sum_{j=1}^{n} p_{i j} m(k \mid j) & \text { otherwise. }\end{cases} \\
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## Commute-Time Distance \& Resistance Distance

- Hence, $c(i, j)$ is called the commute-time distance between $v_{i}$ and $v_{j}$. That amazing property of the resistance distance carries over to the commute-time distance: with a weighting matrix $L^{\dagger}$ and $\operatorname{vol}(V) L^{\dagger}$, respectively.


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\begin{aligned}
c_{i j} & =\operatorname{vol}(V) \cdot r_{i j} \\
& =\operatorname{vol}(V)\left\langle\boldsymbol{e}_{i}-\boldsymbol{e}_{j}, L^{\dagger}\left(\boldsymbol{e}_{i}-\boldsymbol{e}_{j}\right)\right\rangle \\
& =\operatorname{vol}(V)\left(\left(L^{\dagger}\right)_{i i}+\left(L^{\dagger}\right)_{j j}-2\left(L^{\dagger}\right)_{i j}\right) .
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- Hence, both $\sqrt{r_{i j}}$ and $\sqrt{c_{i j}}$ are nothing but a Mahalanobis distance with a weighting matrix $L^{\dagger}$ and $\operatorname{vol}(V) L^{\dagger}$, respectively.


## Commute-Time Distance \& Resistance Distance . . .

- The above formula also allows us to interpret $\sqrt{r_{i j}}$ and $\sqrt{c_{i j}}$ as Euclidean distances on $V(G)$, i.e., to embed $\nu_{i} \in V(G)$ on a point $z_{i} \in \mathbb{R}^{n}$.



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- Let $L=\Phi \Lambda \Phi^{\top}$ be the eigendecomposition of $L$. Then, $L^{\dagger}=\Phi \Lambda^{\dagger} \Phi^{\top}$ where $\Lambda^{\dagger}=\operatorname{diag}\left(\lambda_{0}^{\dagger}, \ldots, \lambda_{n-1}^{\dagger}\right)$ is defined as

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\lambda_{i}^{\dagger}:= \begin{cases}1 / \lambda_{i} & \text { if } \lambda_{i} \neq 0 ; \\ 0 & \text { otherwise. }\end{cases}
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- Hence, $z_{i}^{\top}=\left(\Phi \sqrt{\Lambda^{\dagger}}\right)_{i, 1: n}$ for the resistance distance and $z_{i}^{\top}=\left(\Phi \sqrt{\operatorname{vol}(V) \Lambda^{\dagger}}\right)_{i, 1: n}$ for the commute-time distance.


## Commute-Time Distance \& Resistance Distance ...

- Compare these embeddings with the one used in the spectral clustering ( with $L$ ) in Lecture 7, i.e., $\boldsymbol{y}_{i}^{\top}=\Phi_{i, 1: k}, k \leq n$.


## Commute-Time Distance \& Resistance Distance ...

- Compare these embeddings with the one used in the spectral clustering (with $L$ ) in Lecture 7, i.e., $\boldsymbol{y}_{i}^{\top}=\Phi_{i, 1: k}, k \leq n$.
- $\left\{\boldsymbol{z}_{i}\right\}$ and $\left\{\boldsymbol{y}_{i}\right\}$ could be considerably different. For example, in the optimal case where $G$ consists of $k$ disconnected components, the first $k$ eigenvalues of $L$ are zeros and the corresponding columns of $\Phi$ are the indicator vectors of these $k$ components. However, the first $k$ columns of the matrix $\Phi \sqrt{\Lambda^{\dagger}}$ are zero vectors.


## Commute-Time Distance \& Resistance Distance ...

- Compare these embeddings with the one used in the spectral clustering (with $L$ ) in Lecture 7, i.e., $\boldsymbol{y}_{i}^{\top}=\Phi_{i, 1: k}, k \leq n$.
- $\left\{\boldsymbol{z}_{i}\right\}$ and $\left\{\boldsymbol{y}_{i}\right\}$ could be considerably different. For example, in the optimal case where $G$ consists of $k$ disconnected components, the first $k$ eigenvalues of $L$ are zeros and the corresponding columns of $\Phi$ are the indicator vectors of these $k$ components. However, the first $k$ columns of the matrix $\Phi \sqrt{\Lambda^{\dagger}}$ are zero vectors.
- On the other hand, if $G$ is connected, $z_{i}$ 's got more influenced by the eigenvectors corresponding to the small eigenvalues because $\lambda_{i}^{\dagger}=1 / \lambda_{i}$ if $\lambda_{i} \neq 0$.

