

## Lecture 7 : Rigid body dynamics (and Lie algebras)

For Poisson brackets:

$$(1) (R^3, x) \text{ is Poisson, } (2) (g = SO(3), [ , ])$$

$$\text{Prop: } (R^3, x) \xrightarrow{\text{Poisson}} (SO(3), [ , ]) .$$

The isomorphism is given by

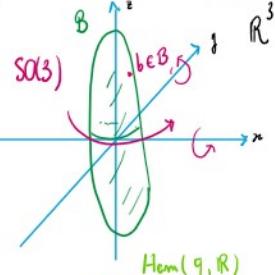
$$(x_1, x_2, x_3) \in R^3 \xrightarrow{\text{map}} A(x) = \begin{pmatrix} 0 & -x_3 & x_2 \\ x_3 & 0 & -x_1 \\ -x_2 & x_1 & 0 \end{pmatrix}.$$

$$\text{We are saying } \varphi(v \times w) = [\varphi(v), \varphi(w)]$$

$$\text{Remark: } A(v) \cdot w = v \times w .$$

out of commutation  
 $[A, B] := AB - BA$

$$ACM_3(A): AA^t = I \quad ACM_3(R): A + A^t = 0$$

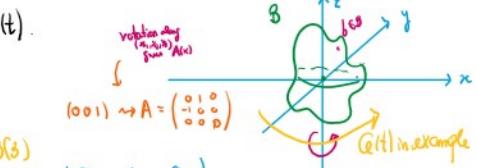


$\Rightarrow$  in general,  $g^*$  always has  
a canonical Poisson bracket  
(from  $[ , ]$  Lie bracket)

## Q 2. Rigid body setup: $b \in B$ , $w(t) = Q(t) \cdot b$ , $Q(t) \in SO(3)$ .

To understand the kinetic term in Hamiltonian we need  $w(t)$  and dual momentum.

$$\dot{w}(t) = \dot{Q}(t) \cdot b = Q(t) \cdot \dot{Q}(t) \cdot w(t).$$



Ex: Rotating B along z-axis gives

$$Q(t) = \begin{pmatrix} \cos(t) & 0 & 0 \\ -\sin(t) & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\dot{Q}(t) = \begin{pmatrix} -\sin(t) & 0 & 0 \\ -\cos(t) & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$T_{SO(3)}_{Q(t)}$$

$$Q(t) \cdot \dot{Q}(t) = \begin{pmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \in so(3), \text{ i.e. antisymmetric.}$$

$$\tilde{Q}(t) = Q(-t) = \begin{pmatrix} \cos(t) & -\sin(t) & 0 \\ \sin(t) & \cos(t) & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

## Q 2. Symmetries: $M = g^* = SO(3)^*$ , $H(M) = \frac{1}{2} M^t \cdot \Pi^{-1} \cdot M$ ,

(i) The energy is a conserved quantity, since the  
level sets of H are ellipsoids:

$$H = \frac{1}{2} \left( \frac{M_1^2}{I_1} + \frac{M_2^2}{I_2} + \frac{M_3^2}{I_3} \right)$$

(ii) A  $\mathbb{Z}_2$  symmetry exists:  $J = M_1^2 + M_2^2 + M_3^2$ .

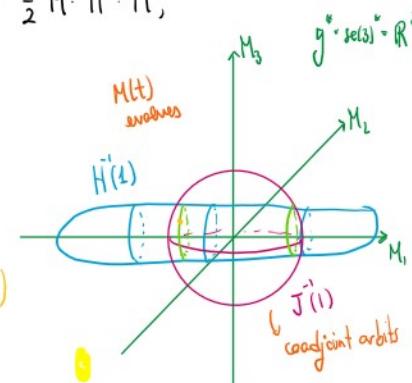
Proof: we want  $\{J, H\} = 0$ .

$$\{M_i^2, H\} = 2M_i M_j \cdot M_i = 2M_i M_2 M_3 \cdot (I_3 - I_2) / I_2 I_3$$

$$\{M_2^2, H\} = 2M_2 M_3 \left( \frac{I_1 - I_3}{I_2 I_3} \right)$$

$$\{M_3^2, H\} = 2M_3 M_1 \left( \frac{I_2 - I_1}{I_2 I_3} \right)$$

$$\{J, H\} = 2 \cdot M_1 M_2 M_3 \left( \frac{I_1(I_2 - I_3) + \dots}{I_1 I_2 I_3} \right) = 0. \quad \square$$



$$(1) K_b = \frac{m_b}{2} \|w(t)\|^2 = \frac{m_b}{2} \|x(t) \times w(t)\|^2 = \frac{m_b}{2} \|Q(t) \cdot (x \times w(t))\|^2 = \frac{m_b}{2} \|Q(t) \cdot (\Pi^t M)\|^2 = \frac{m_b}{2} \|M\|^2 \quad \text{out of rot.}$$

$$\text{This is at a point } b \in B: K = \int_B K_b dV = \int_B \frac{m_b}{2} \|M\|^2 dV \quad M = \Pi \cdot \bar{M} \quad \bar{M} \in SO(3)$$

$$\text{The dual momentum is given by } M = \Pi \cdot \bar{M} = (M_1, M_2, M_3) \in SO(3)^* \quad \Rightarrow \quad H: SO(3)^* \rightarrow \mathbb{R}, \quad H(M_1, M_2, M_3) = \frac{1}{2} M^t \cdot \Pi^{-1} \cdot M.$$

$$\text{Remark: we'll assume } \Pi = \text{diag}(I_1, I_2, I_3) \quad \text{moment of inertia. So } H(M_1, M_2, M_3) = \frac{1}{2} \left( \frac{M_1^2}{I_1} + \frac{M_2^2}{I_2} + \frac{M_3^2}{I_3} \right).$$

$$\Rightarrow \text{Eq of motion } \left\{ \begin{array}{l} M_1 = \frac{d}{dt} M_1, H_f \\ M_2 = \frac{d}{dt} M_2, H_f \\ M_3 = \frac{d}{dt} M_3, H_f \end{array} \right\}$$

$$\text{e.g. } \{M_1, H\} = (M_1, M_2, M_3) \cdot ((1, 0, 0) \times \left( \frac{M_2}{I_1}, \frac{M_3}{I_1}, \frac{M_1}{I_3} \right)) = M_2 M_3 \left( \frac{I_1 - I_3}{I_1 I_3} \right)$$

$$\text{same in } x = (M_1, M_2, M_3) \cdot \left( 0, -\frac{M_3}{I_2}, \frac{M_2}{I_2} \right) = \frac{M_2 M_3}{I_2 I_3} (I_3 - I_2)$$