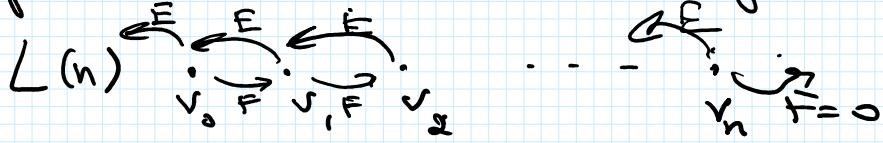


We constructed irreducible

representations of the Lie algebra $\mathfrak{sl}_2(\mathbb{C})$:



and for the Lie group $SL(2, \mathbb{C})$:

$$\begin{aligned} V_n &= \{ \text{degree } n \text{ homogeneous polynomials } \\ &\quad \text{in } x, y \} \\ &= \langle x^n, x^{n-1}y, x^{n-2}y^2, \dots, y^n \rangle \end{aligned}$$

This is the same basis up to scaling, $V_n \cong L(n)$.

Any complex representation of $\mathfrak{sl}_2(\mathbb{C}) / SL(2, \mathbb{C})$ is a direct sum of such representations.

Rank $SU_2 \otimes \mathbb{C} = \mathfrak{sl}_2(\mathbb{C})$, so all irreducible complex representations of SU_2 and $SU(2)$ are isomorphic to V_n for some n .

Fact ① Irreducible representations of $SO(3)$ are isomorphic to $L(n)$

② A representation $L(n)$ lifts to a representation of Lie group $SO(3)$ if and only if n is even.

Proof ① $SO(3) \cong SU(2)$ as Lie algebras, they have same representations.

② $SO(3) \cong SU(2) / \{\pm I\}$, in particular

we have a homomorphism $\Phi: SU(2) \rightarrow SO(3)$

Any representation $p: SO(3) \rightarrow GL(V)$ yields a

representation of $SU(2)$: $SU(2) \xrightarrow{\Phi} SO(3) \xrightarrow{f} GL(V)$

Conversely, given a representation of $SU(2)$, we get a representation of $SO(3)$ iff $\pm I$ acts trivially.

$$V_n = \langle x^n, x^{n-1}y, \dots, y^n \rangle$$

$-I \in SU(2)$ acts on V_n diagonally: $x \mapsto -x, y \mapsto -y$

$$x^n \mapsto (-1)^n x^n$$

$$x^{n-1}y \mapsto (-1)^n x^{n-1}y$$

$$-I \text{ acts by } (-1)^n I.$$

$\Rightarrow -I$ acts trivially iff n even. □

For Irred. representations of $SO(3; \mathbb{R})$ are

$$V_0 = \langle 1 \rangle \text{ trivial rep. } V_2 = \langle x^2, xy, y^2 \rangle = \text{defining representation of } SO(3).$$

$\dim V_4 = 5, \dim V_6 = 7, \dots$

Rmk Related to spectrum of hydrogen/quantum mechanics.

Operations on representations

$G = \text{Lie group}$ $\mathfrak{g} = \text{Lie}(G)$ Lie algebra

$U, V = \text{representations}$

• Direct sum: $U \oplus V$, $\dim U \oplus V = \dim U + \dim V$

• Tensor product: $U \otimes V$ tensor product of underlying vector spaces

$$\dim U \otimes V = \dim U \cdot \dim V$$

$A \in G$ acts on $u \otimes v \in U \otimes V$

$$\text{by } A(u \otimes v) = (Au) \otimes (Av).$$

— blunder +

by $\text{Tr}(u \otimes v) = (\text{Tr} u) \otimes (\text{Tr} v)$.

$X \in g$ acts on $u \otimes v$:

$$X(u \otimes v) = (Xu) \otimes v + u \otimes (Xv)$$

product rule for derivatives.

In particular, if $Xu = \lambda u$, $Xv = \mu v$ then

$$X(u \otimes v) = (\lambda + \mu)u \otimes v.$$

Characters T = complex representation of $\mathfrak{sl}_2(\mathbb{C})$

Recall: H diagonalizes, eigenvalues of H are integers.

a_i = dimension of eigenspace for H with eigenvalue i

$$\text{ch } T = \sum_i a_i q^i \leftarrow \text{Laurant polynomial in } q$$

$$\text{ch } L(n) = q^n + q^{n-2} + \dots + q^{-n}$$

Fact $\text{ch}(V \oplus W) = \text{ch}(V) + \text{ch}(W)$

$$\text{ch}(V \otimes W) = \text{ch}(V) \cdot \text{ch}(W)$$

From Lie group perspective, we can consider

$$e^{tH} = \begin{pmatrix} e^t & 0 \\ 0 & e^{-t} \end{pmatrix}$$

On an H -eigenspace with eigenvalue k , e^{tH} acts by e^{tk}

Define $q = e^t$

$$e^{tH} = \begin{pmatrix} q & 0 \\ 0 & q^{-1} \end{pmatrix} \text{ acts on } H\text{-eigenspace with eigenvalue } k \text{ by } q^k$$

Conclusion

$$\boxed{\text{ch}(V) = \text{Tr}_V \begin{pmatrix} q & 0 \\ 0 & q^{-1} \end{pmatrix}}.$$

element of $\text{SL}(2; \mathbb{C})$

$g \rightarrow \text{Tr}_V(g)$ is a function on G

No just restrict this function to $\{ \text{diagonal}, \dots \} \subset G$.

(v matrices)

$$V_2 = L(z) = \langle x^2, xy, y^2 \rangle \quad \begin{pmatrix} q & 0 \\ 0 & \bar{q}^{-1} \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} qx \\ q^{-1}y \end{pmatrix}$$

$$\mathrm{ch} V_2 = q^2 + 1 + q^{-2}$$

$$\begin{aligned} \mathrm{ch} V_2 \otimes V_2 &= (q^2 + 1 + q^{-2})(q^2 + 1 + q^{-2}) = \\ &= q^4 + 2q^2 + 3 + 2q^{-2} + q^{-4} \quad (\star) \end{aligned}$$

How to decompose $V_2 \otimes V_2$ into irreducibles?

$$(\star) = \underbrace{(q^4 + q^2 + 1 + q^{-2} + q^{-4})}_{\mathrm{ch} V_4} + \underbrace{(q^2 + 1 + q^{-2})}_{\mathrm{ch} V_2} + \underbrace{\frac{1}{2}}_{\mathrm{ch} V_0}$$

$$V_2 \otimes V_2 \cong V_4 \oplus V_2 \oplus V_0$$

Rmk $V_2 \otimes V_2 = \mathrm{Sym}^2(V_2) \oplus \Lambda^2(V_2)$

$$\begin{array}{ccc} \mathrm{SS} & \downarrow & \Downarrow \\ \begin{matrix} 3 \times 3 \text{ matrices} \\ \mathfrak{so}(3) \end{matrix} & \begin{matrix} \text{symmetric} \\ 3 \times 3 \text{ matrices} \\ \mathrm{dim} = 6 \end{matrix} & \begin{matrix} \text{skew-symmetric} \\ 3 \times 3 \text{ matrices} \\ \mathrm{dim} = 3 \end{matrix} \end{array}$$

HW 6: $V_0 \otimes V_0 = ?$

$V_1 \otimes V_1 \otimes V_1 = ?$

By decomposition above, we see that $\Lambda^2 V_2 \cong V_2$

$$\mathrm{Sym}^2(V_2) \cong V_4 \oplus V_0$$

$\mathrm{SO}(3)$ acts on the space of symmetric 3×3 matrices

preserving the matrix $\begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$ which spans a trivial representation in $\mathrm{Sym}^2(V_2)$.