

Fast Interface Models with Dimensional Reduction

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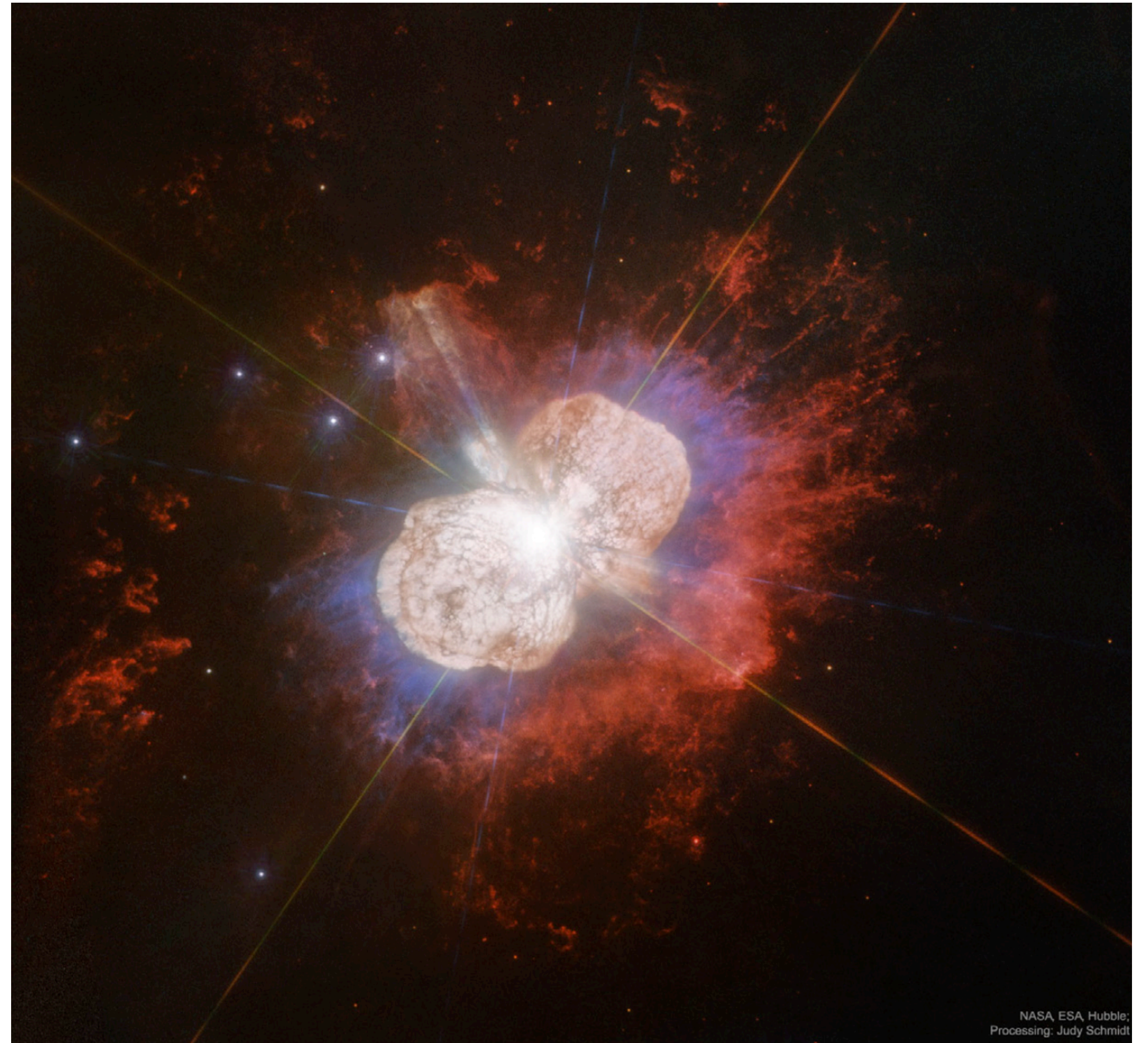
Joint work with Steve Shkoller

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Methodology

Multiscale Flow via Interface Models

- Action at many scales
- Direct numerical simulation is computationally expensive
- **Idea:** decompose flow into simple constituent parts
 - e.g. transport, diffusion, shock waves, contacts, acoustic waves, flame fronts, etc.
- Model using interface methods

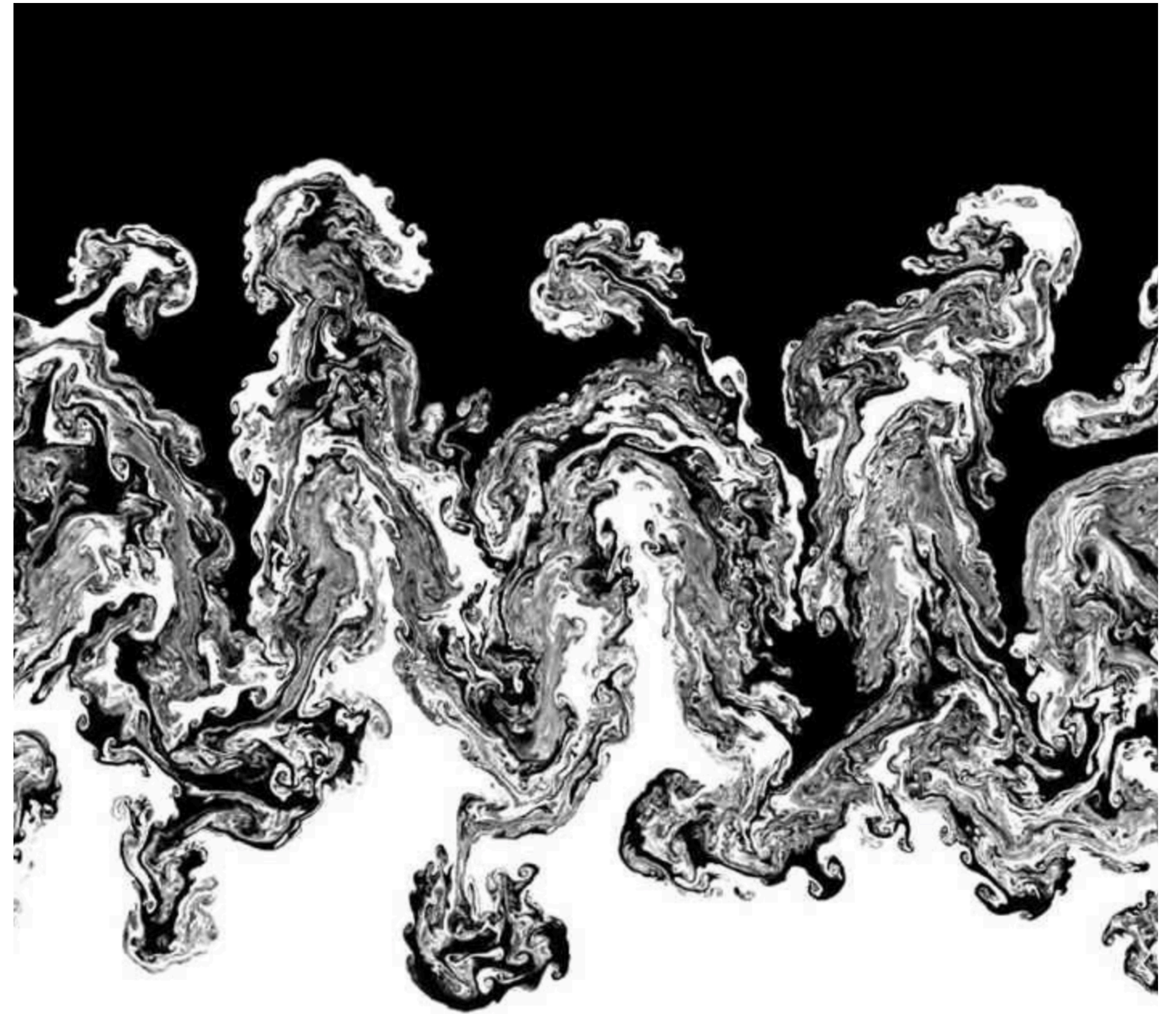


NASA, ESA, Hubble;
Processing: Judy Schmidt

Interface Methods

Strengths and Weaknesses

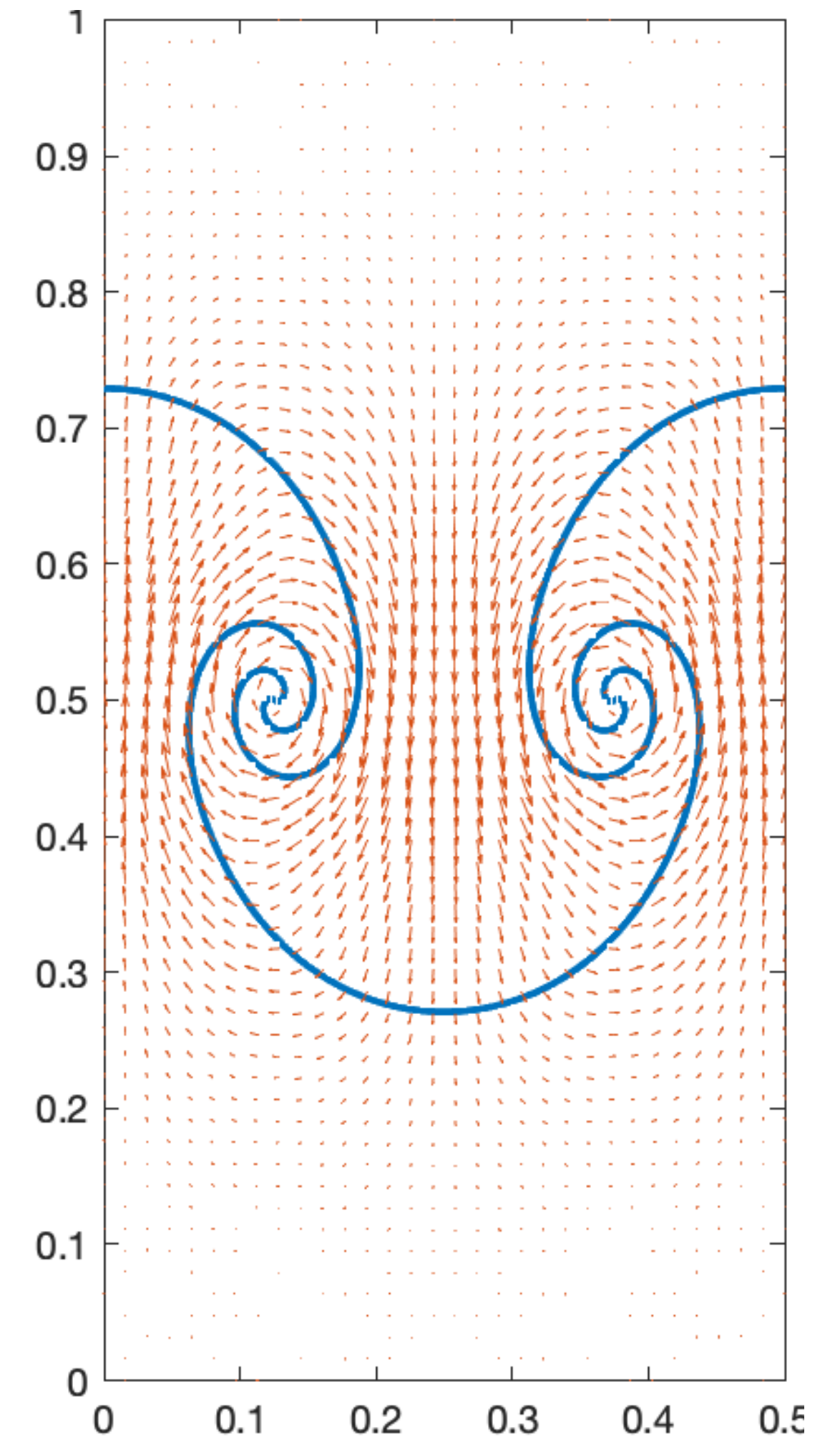
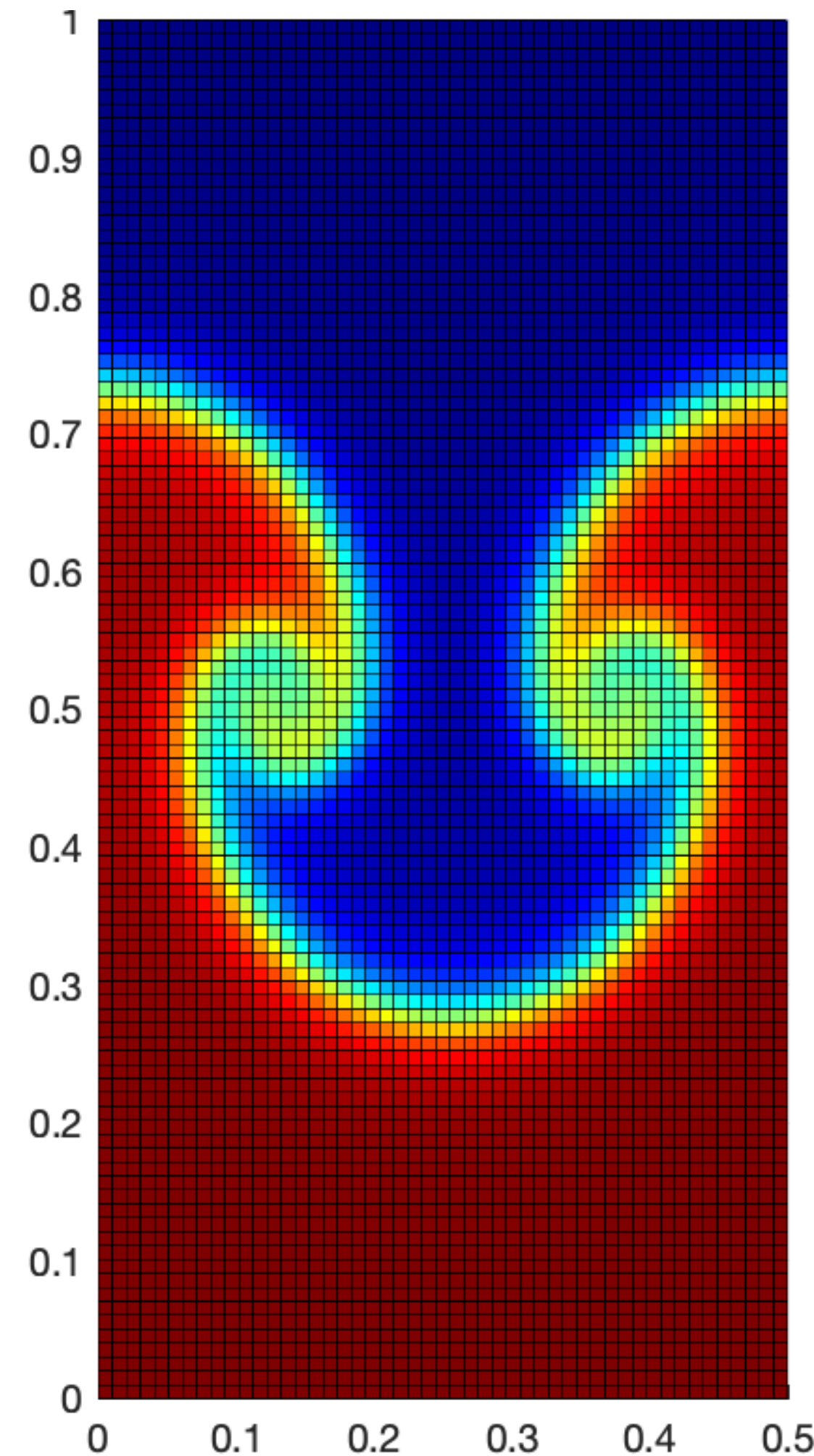
- Fluid systems often dominated by propagating waves and surfaces of discontinuity.
- Interface methods offer geometric flexibility, computational efficiency, and concise representation.
- Less useful at scales where flow is dominated by mixing or turbulence.



Incompressible Flow

Contact Discontinuities

- Nonlocal effects of incompressible flow recovered using boundary integral methods.
- We developed the **z-model** for contact discontinuities in 2d & 3d.
- Matches analytical and experimental predictions for Rayleigh-Taylor Instability (RTI).

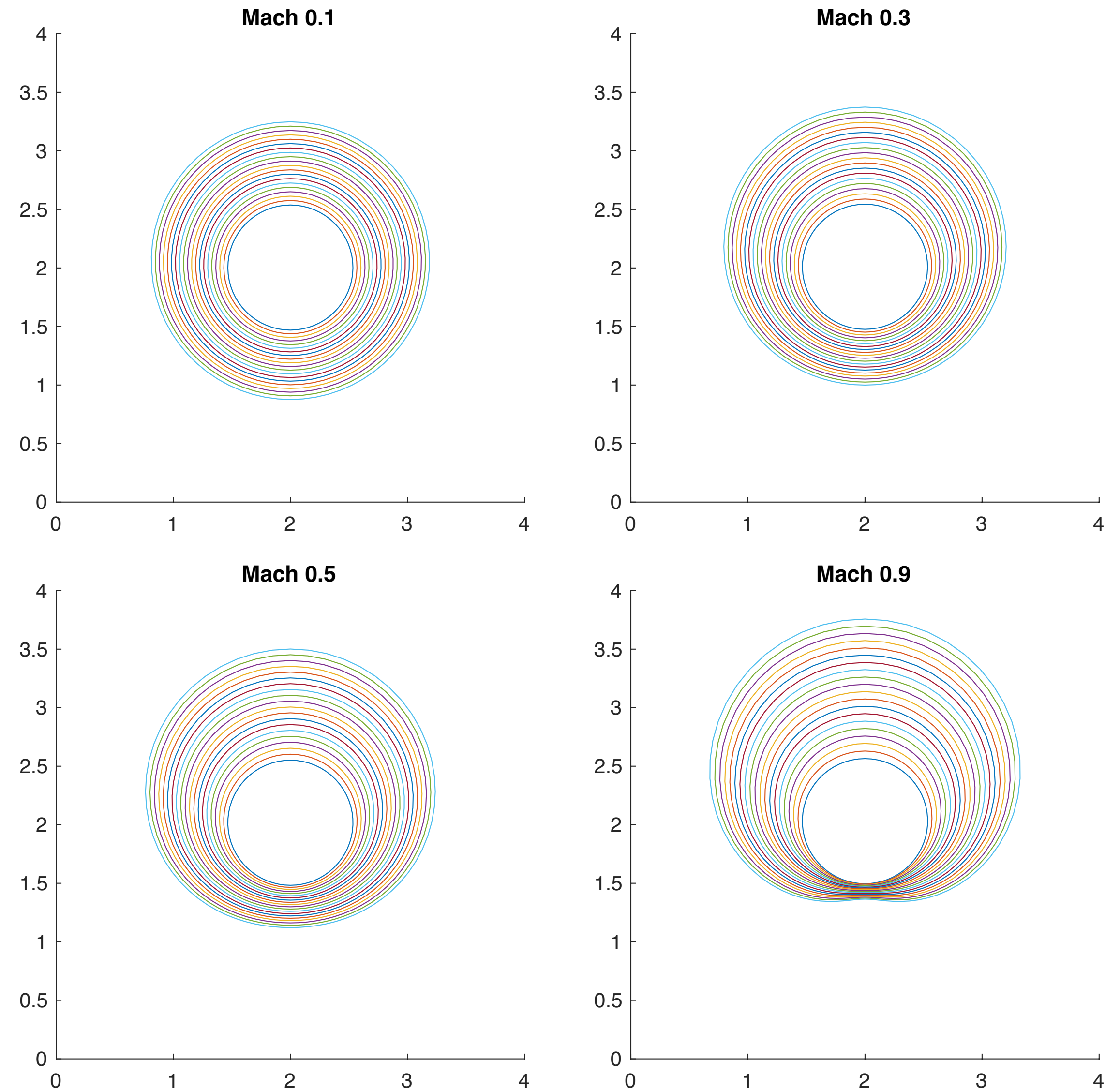


Rayleigh-Taylor instability (z-model with reconstructed density/velocity)

Compressible Flow

Nonlinear Sound Waves

- Novel geometric coordinates which align with propagating wave or surface of discontinuity.
- The **wza-model** for nonlinear sound waves in 2d.
- Extension to shock waves/vorticity waves/contacts in progress



Sound wave emitted from moving object at various speeds (abc-model)

Relevant Publications

1. Computational Methods for the Euler Equations of Gas Dynamics

1. J. REISNER, J. SERENCSEA, AND S. SHKOLLER, **A space-time smooth artificial viscosity method for nonlinear conservation laws**, *Journal of Computational Physics*, 235, 912–933, (2013)
2. R. RAMANI, J. REISNER, AND S. SHKOLLER, **A space-time smooth artificial viscosity method with wavelet noise indicator and shock collision scheme, Part 1: the 1-D case**, *Journal of Computational Physics*, 387, 81–116, (2019)
3. R. RAMANI, J. REISNER, AND S. SHKOLLER, **A space-time smooth artificial viscosity method with wavelet noise indicator and shock collision scheme, Part 2: the 2-D case**, *Journal of Computational Physics*, 387, 45–80, (2019)
4. R. RAMANI and S. SHKOLLER, **A multiscale model for Rayleigh-Taylor and Richtmyer-Meshkov instabilities**, *Journal of Computational Physics*, 405, 109177, (2020)

2. Fast interface Models for Contact Discontinuities in Incompressible Flow

1. A. CHENG, R. GRANERO-BELINCHON, S. SHKOLLER, and J. WILKENING, **Rigorous asymptotic models of water waves**, *Water Waves*, 1, 71--130, (2019)
2. J. CANFIELD, N. DENISSEN, M. FRANCOIS, R. GORE, R. RAUENZAHN, J. REISNER, S. SHKOLLER, **A comparison of interface growth models applied to Rayleigh-Taylor and Richtmyer-Meshkov instabilities**, *Journal of Fluids Engineering*, 142(12), 121108, (2020)
3. G. PANDYA and S. SHKOLLER, **Interface models for three-dimensional Rayleigh-Taylor instability**, *Journal of Fluid Mechanics*, 959, A10, (2023)

3. Fast Interface models for Sound Waves, Shock Waves, Vorticity Waves, and Contacts in Compressible Flow

1. S. SHKOLLER and V. VICOL, **The geometry of maximal development for the Euler equations**, (2023), *preprint*.
2. G. PANDYA and S. SHKOLLER, **Interface model for nonlinear sound wave propagation based on new geometric Riemann variables**, *In Preparation*.

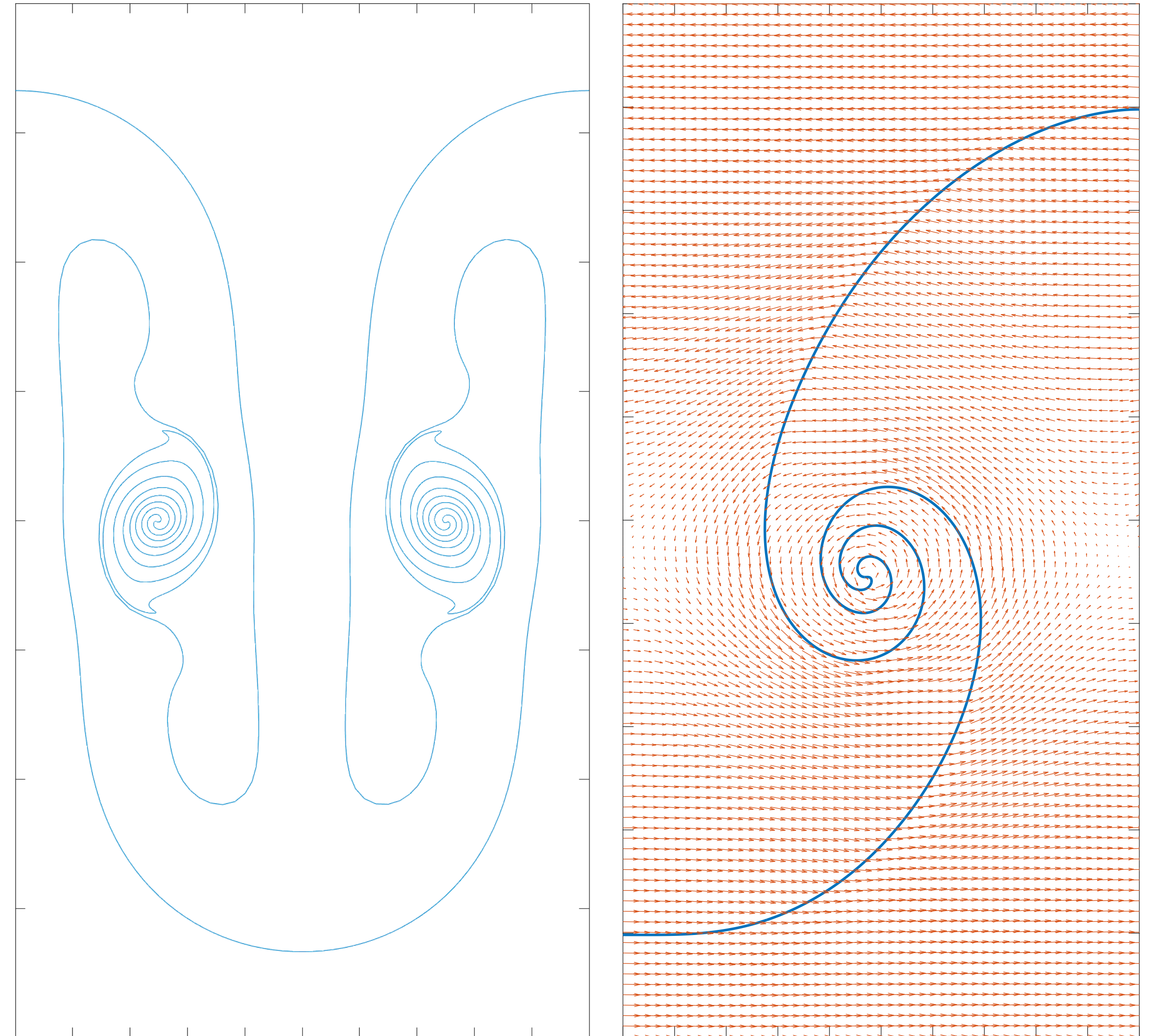
The z-Model

(P. and Shkoller, JFM '23)

- Interface model for the interface position and amplitude of vorticity.
- Models the deposition of vorticity on density discontinuities:

$$\frac{D\omega}{Dt} = \frac{\nabla \rho \times \nabla p}{\rho^2}.$$

- Velocity reconstructed using nonlocal Biot-Savart integral.



z-model interface, Rayleigh-Taylor instability

with reconstructed velocity

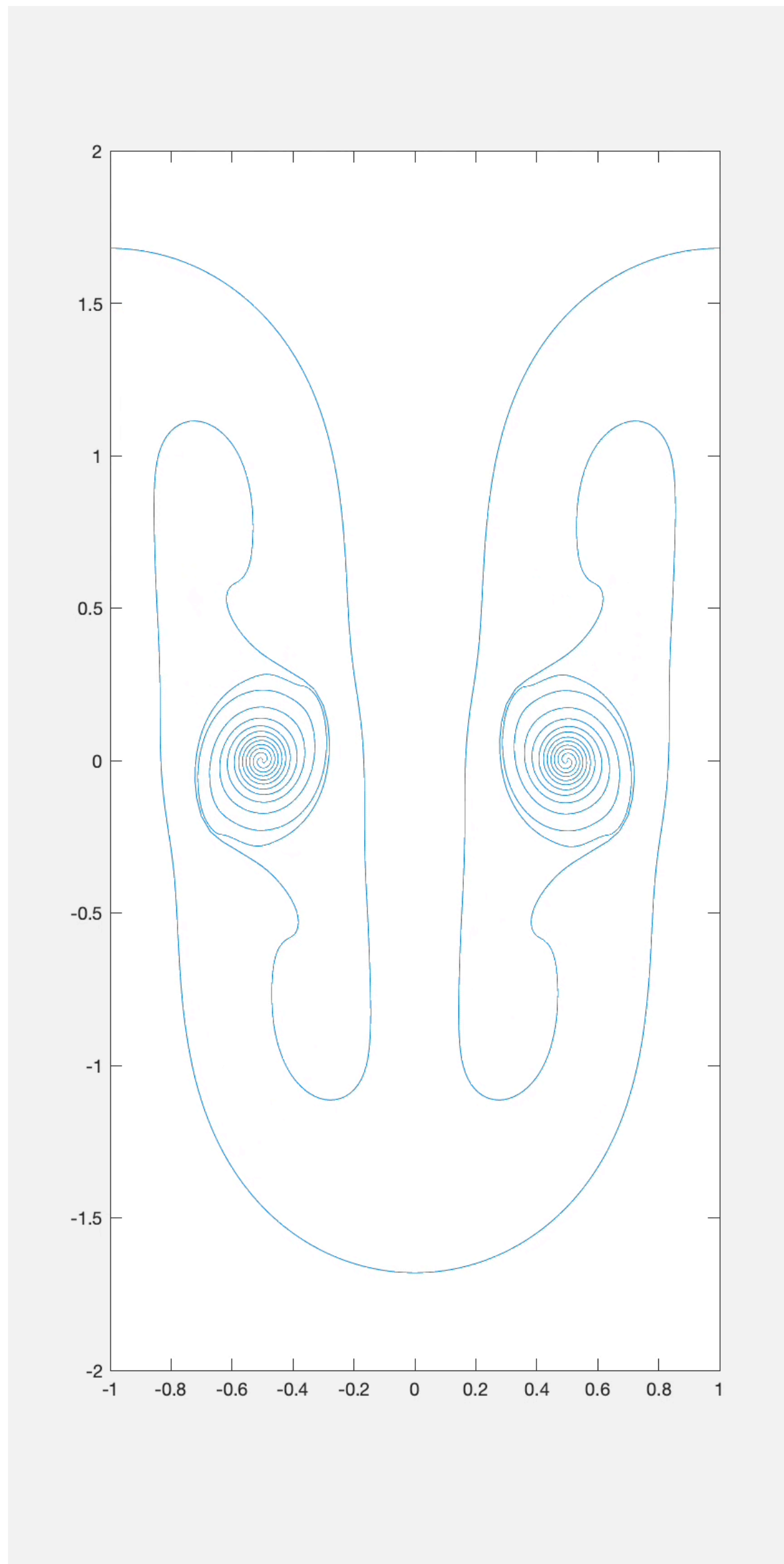
Velocity Reconstruction

- An arbitrary velocity field u can be reconstructed from its **vorticity** $\omega = \text{curl } u$ and **compression** $\chi = \text{div } u$:

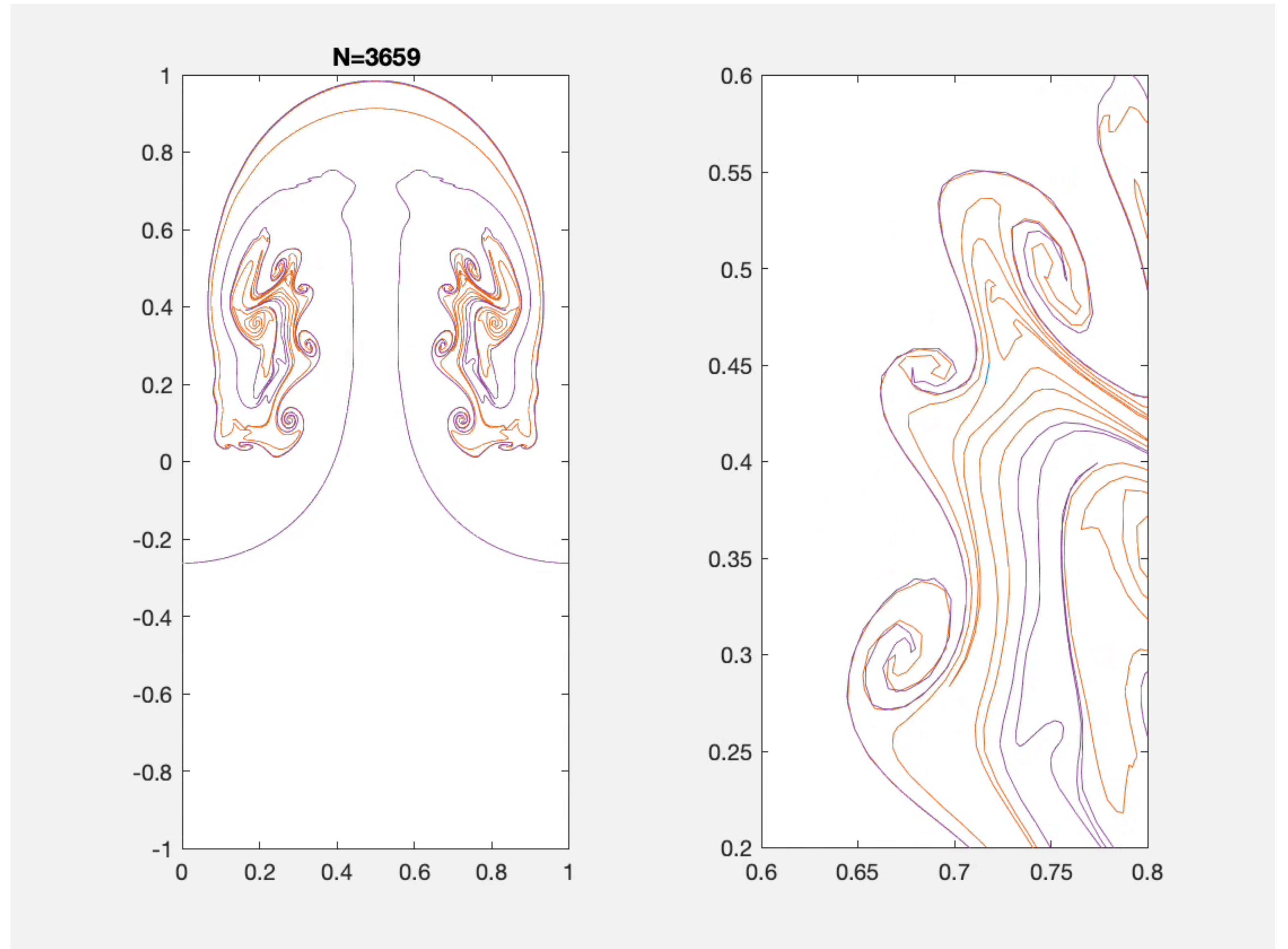
$$u(x) = \frac{1}{4\pi} \iiint_{\mathbb{R}^3} \frac{\omega(x') \times (x - x')}{|x - x'|^3} + \frac{\chi(x')(x - x')}{|x - x'|^3} dx'$$

- Given a vortex sheet at $\Gamma(t)$ in incompressible flow with amplitude of vorticity $\bar{\omega}(x, t)$, $x \in \Gamma(t)$, the velocity is given by the *Biot-Savart law*,

$$u(x, t) = \frac{1}{4\pi} \iint_{\Gamma} \frac{\bar{\omega}(x', t) \times (x - x')}{|x - x'|^3} dS(x')$$

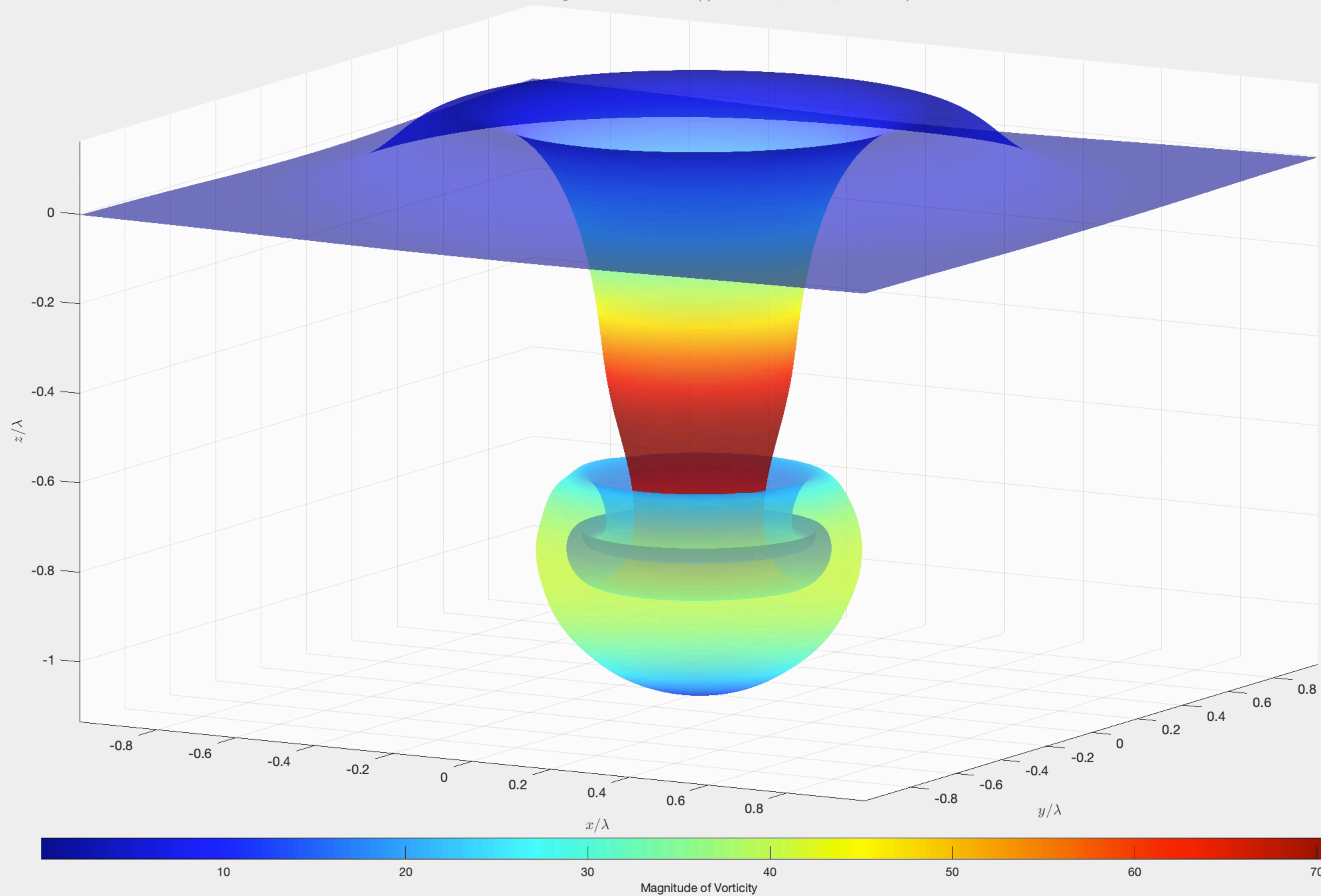


Single-mode Rayleigh-Taylor instability

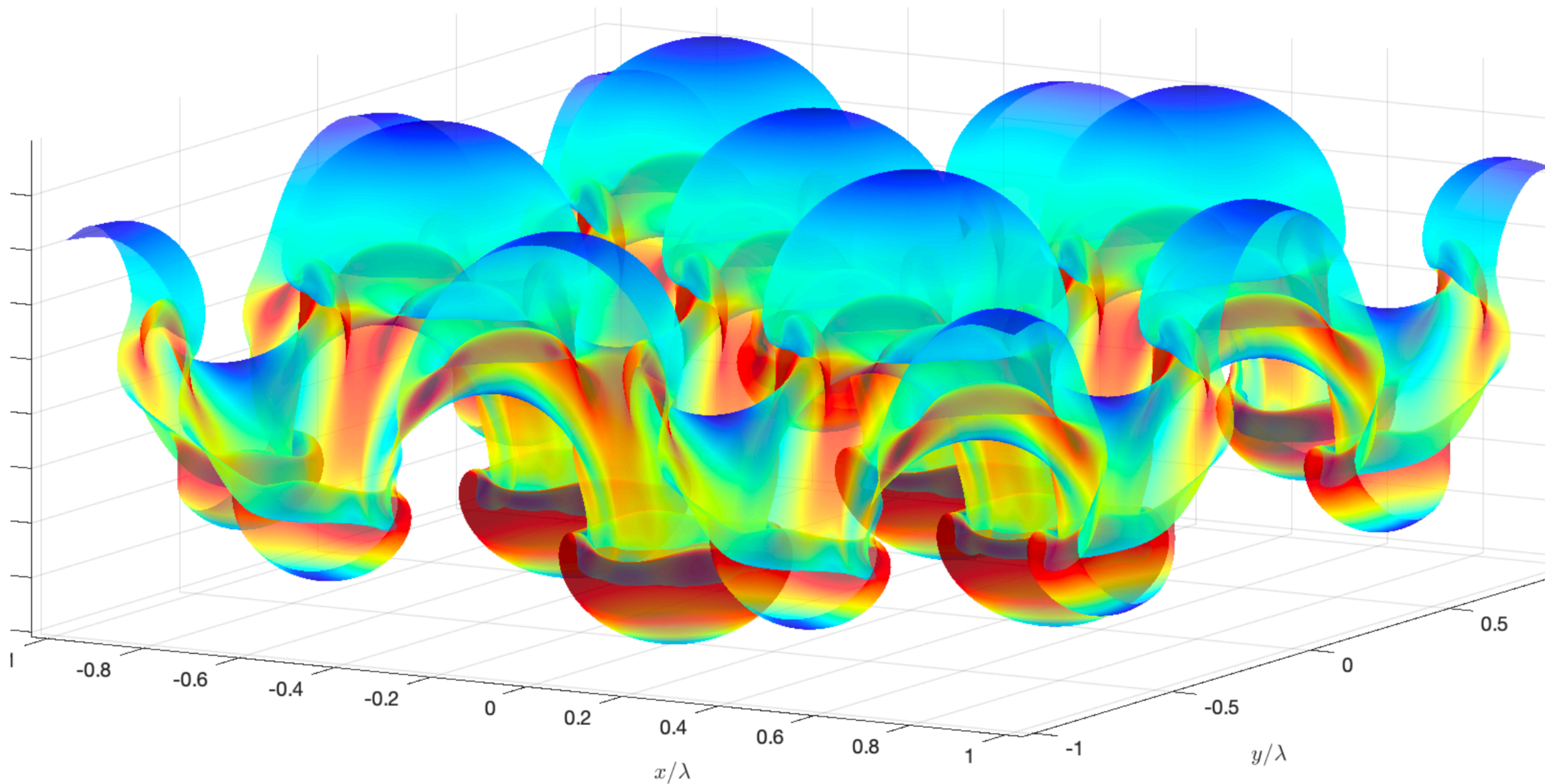


Rising bubble interacting with unstable density interface

256×256 Higher-Order Interface ($t/\tau = 2.368, A = 0.5, \epsilon = 0.0039$)



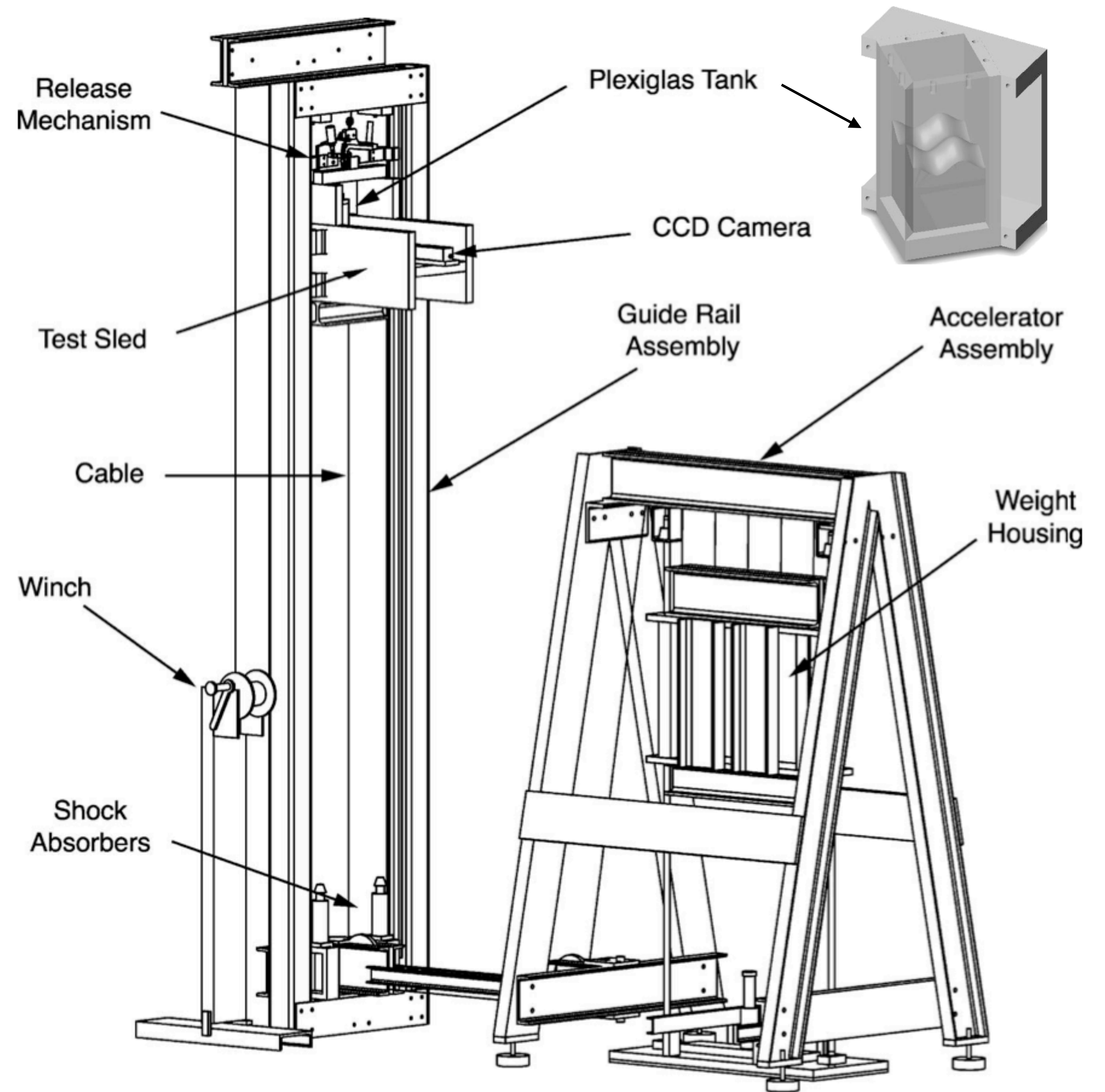
3D Single-Mode RTI



Comparison With Experiments

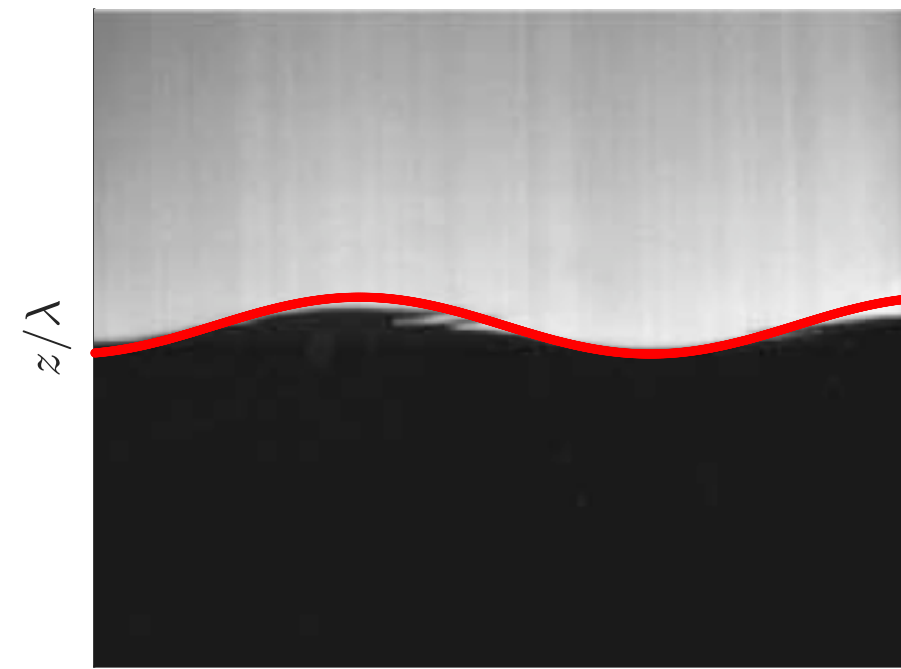
Single-Mode RT Instability

- A standing wave is excited in a tank containing two fluids, the heavier fluid on the bottom.
- The tank is then accelerated downward, so that the acceleration points from the heavier fluid into the lighter fluid.
- The resulting Rayleigh-Taylor instability is photographed along the tank's diagonal using PLIF (planar laser induced fluorescence).



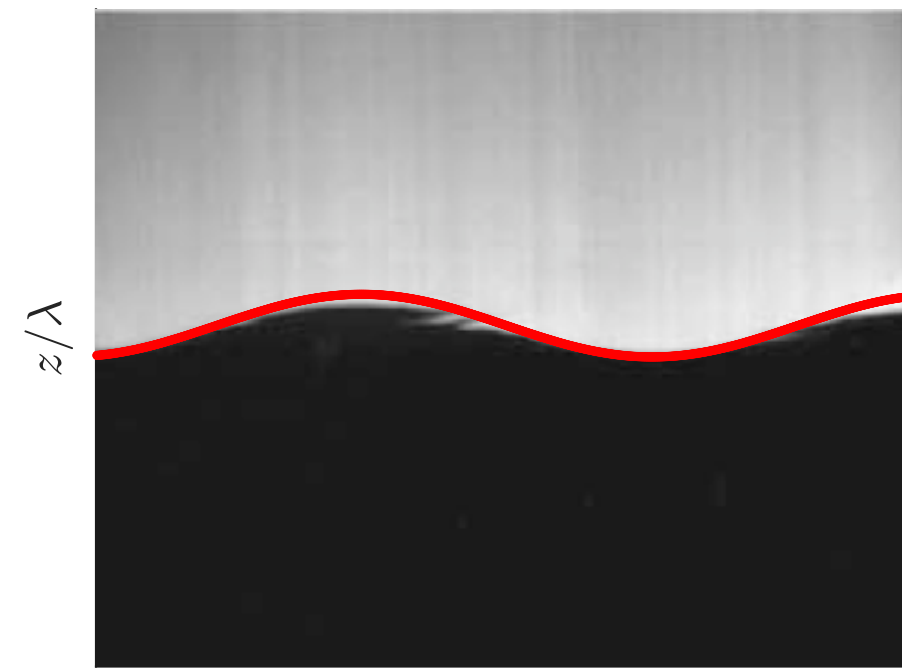
Experimental setup of Waddell/Jacobs and Wilkinson/Jacobs

Diagonal cross-section, $t/\tau = 0$



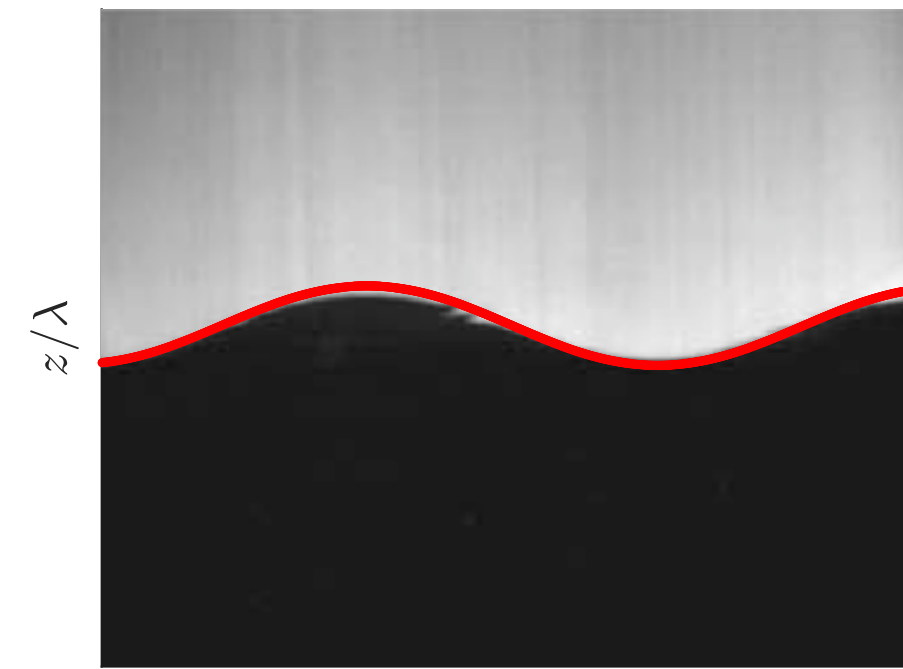
(diagonal distance)/ λ

Diagonal cross-section, $t/\tau = 0.21$



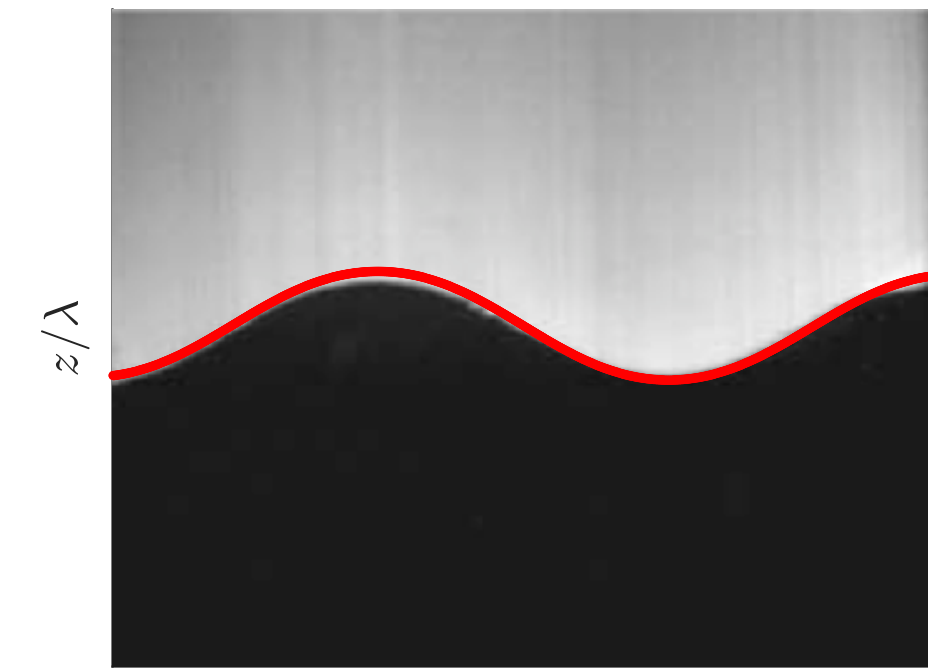
(diagonal distance)/ λ

Diagonal cross-section, $t/\tau = 0.42$



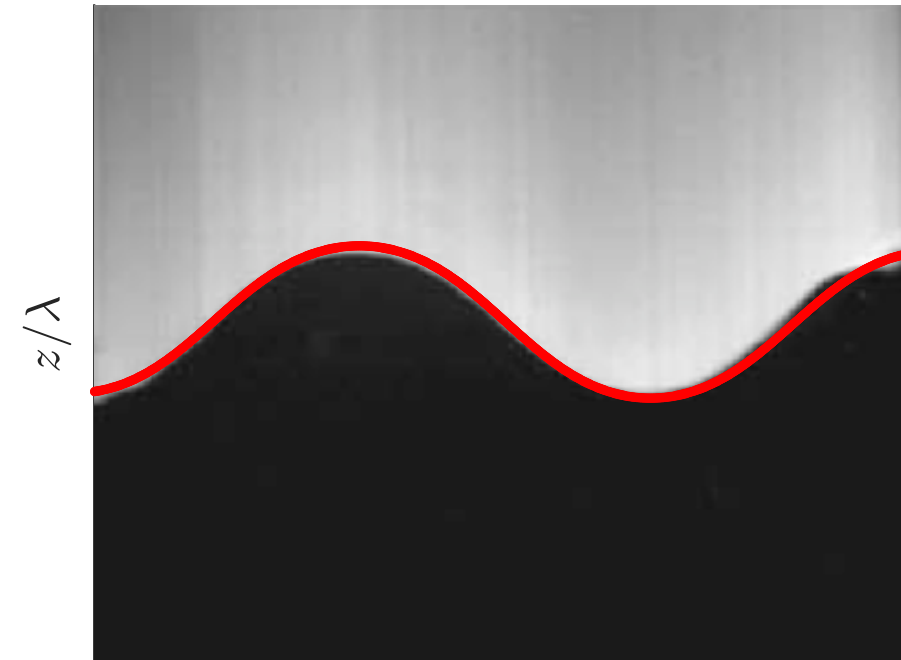
(diagonal distance)/ λ

Diagonal cross-section, $t/\tau = 0.63$



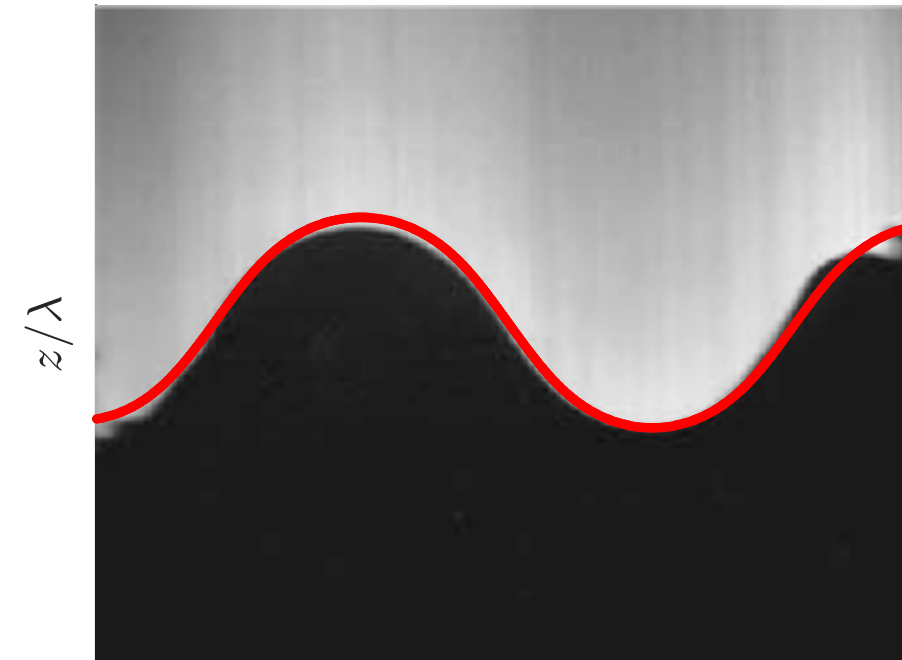
(diagonal distance)/ λ

Diagonal cross-section, $t/\tau = 0.83$



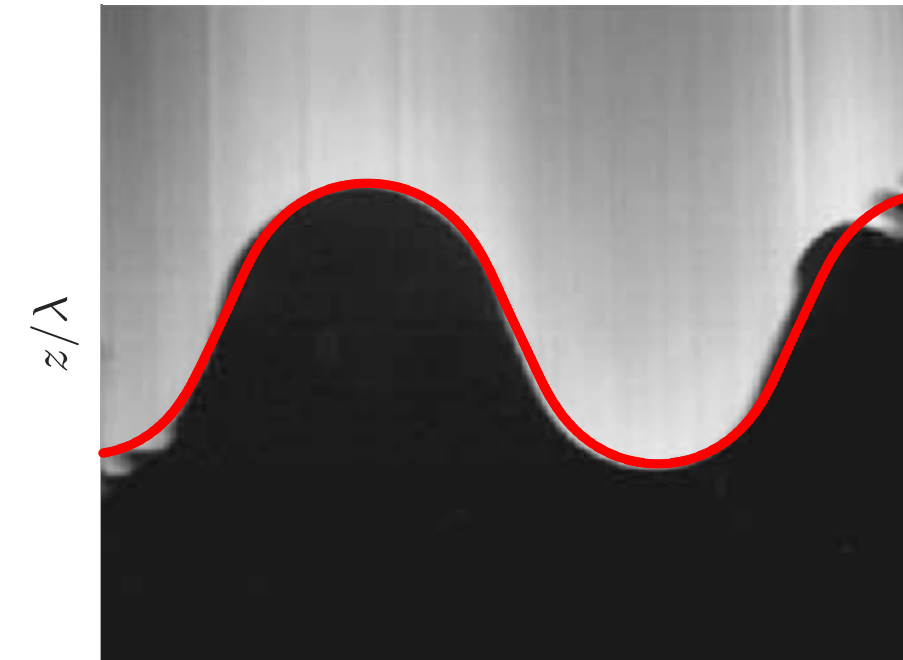
(diagonal distance)/ λ

Diagonal cross-section, $t/\tau = 1.04$



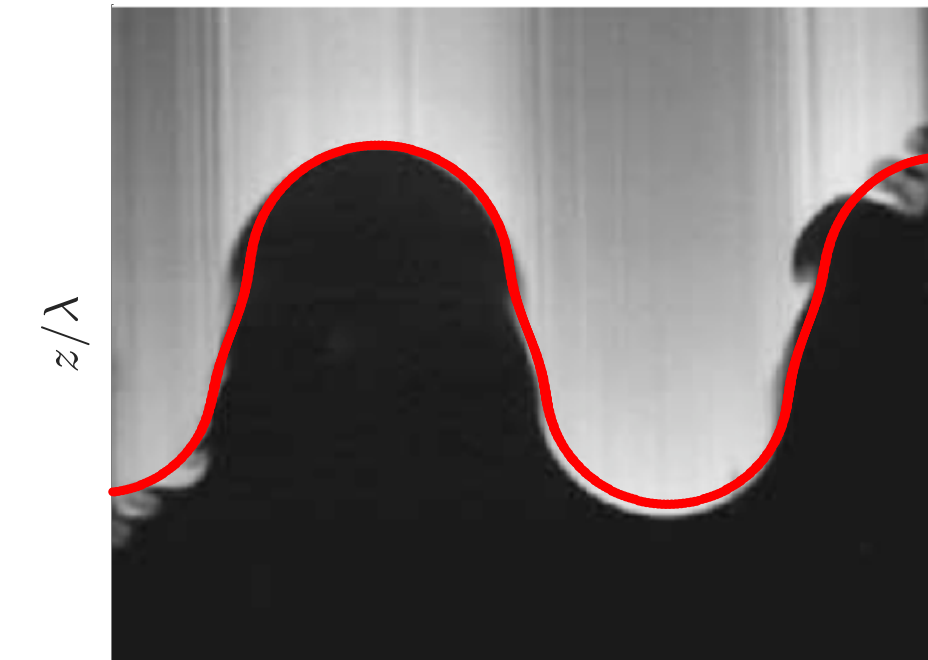
(diagonal distance)/ λ

Diagonal cross-section, $t/\tau = 1.25$



(diagonal distance)/ λ

Diagonal cross-section, $t/\tau = 1.45$



(diagonal distance)/ λ

Diagonal cross-section, $t/\tau = 1.66$



(diagonal distance)/ λ

Diagonal cross-section, $t/\tau = 2$



(diagonal distance)/ λ

Comparison of z -Model to
Experimental Data of Wilkinson/
Jacobs, Showing Diagonal
Cross-Section

Comparison With Experiments

Rocket Rig Experiment

- Theory predicts that the RT mixing layer grows quadratically with time:

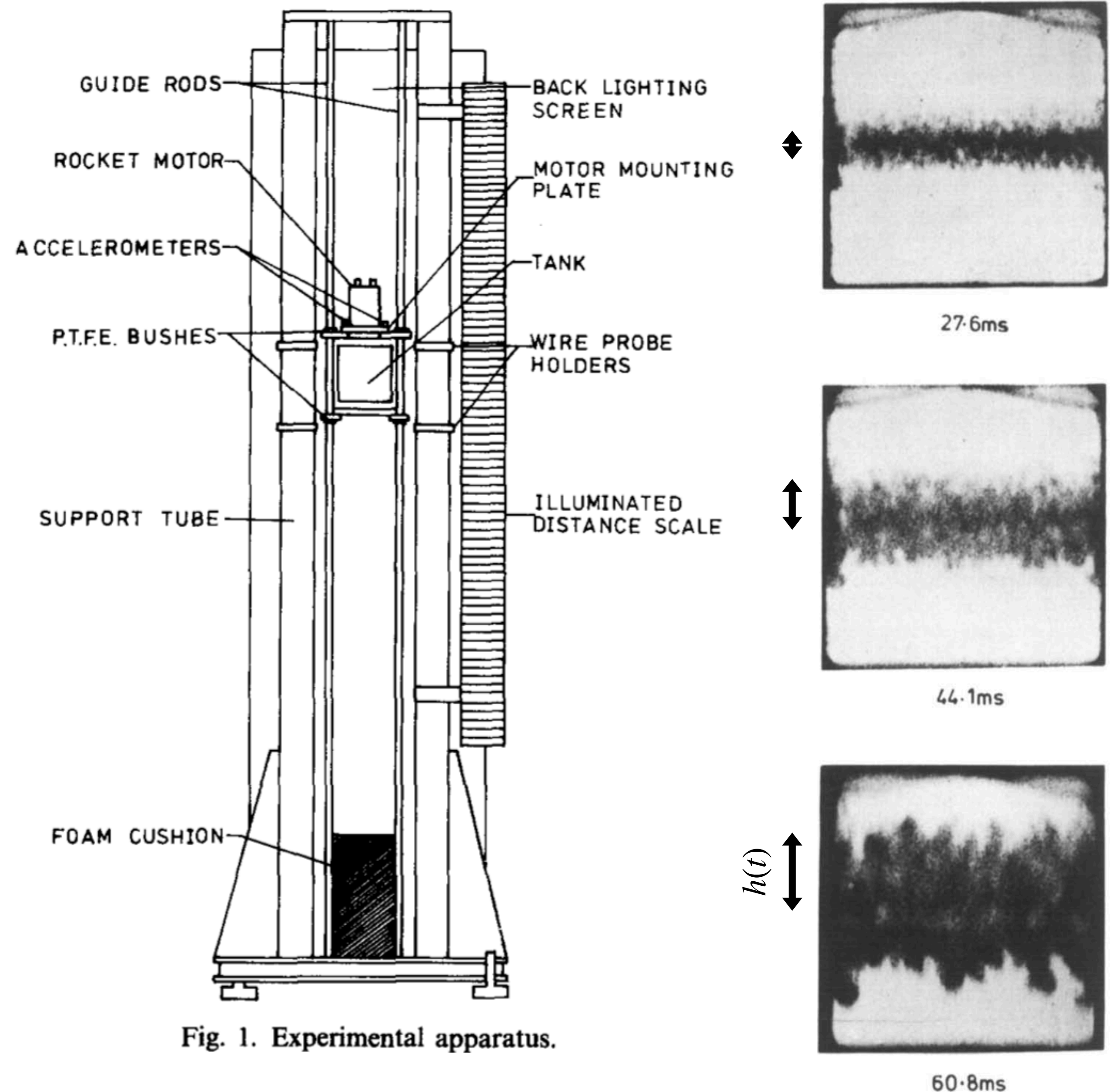
$$h - h_0 = \alpha A g t^2$$

- Experiments of Read and Youngs found constant of proportionality

$$\alpha \approx 0.06 \text{ or } 0.07$$

- Model of Cabot/Cook & Ristorcelli/
Clark measures α dynamically:

$$\alpha = \dot{h}^2 / (4Agh)$$



Comparison With Experiments

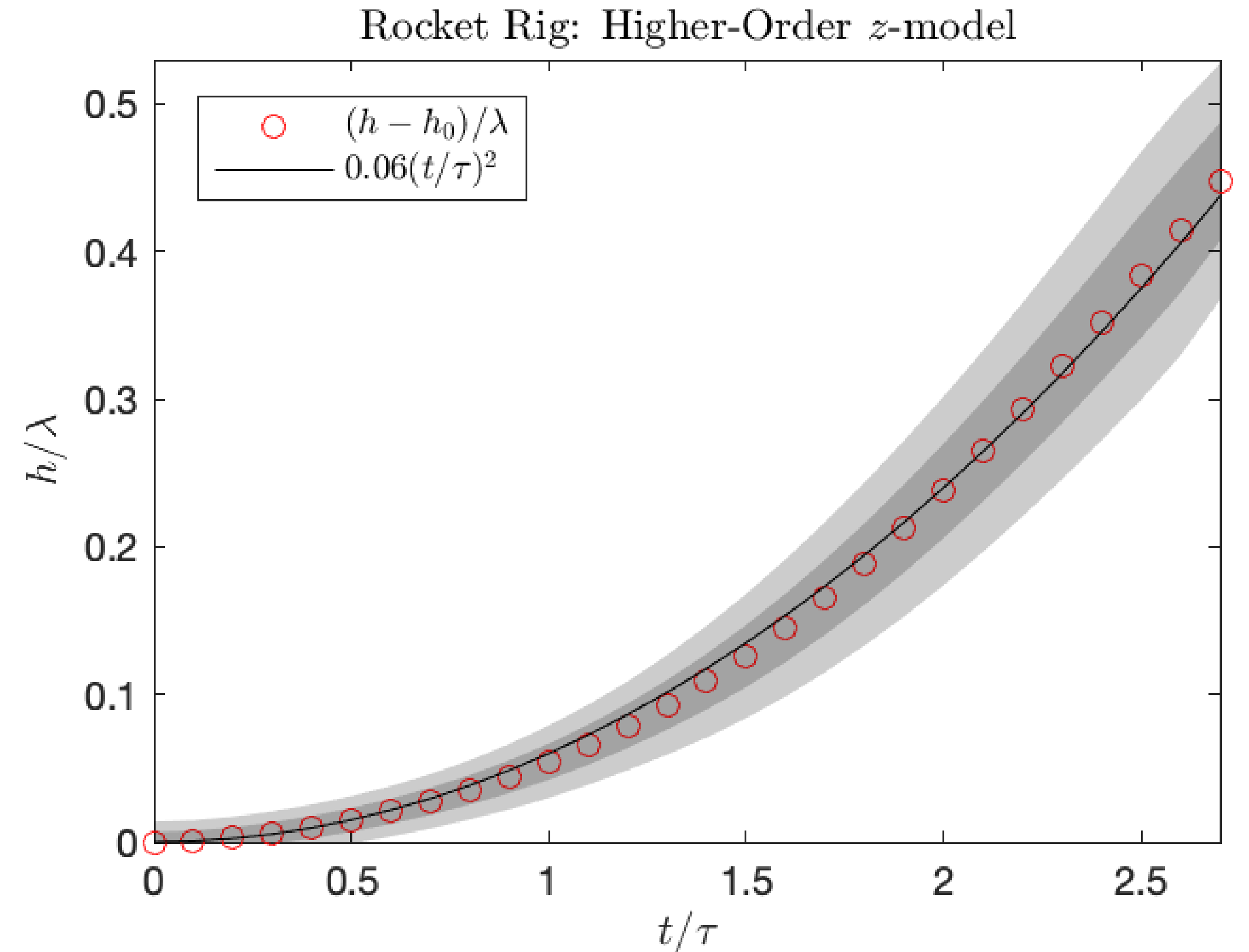
Rocket Rig

We begin with randomized initial data

$$z_3 = \text{Re} \sum_{k_1, k_2 = -N}^N c_{k_1, k_2} e^{2\pi i(k_1 s_1 + k_2 s_2)}$$

where the coefficients c_{k_1, k_2} are sampled from a normal distribution.

We averaged the result of 100 runs, using parameters that match Read's NaI/Pentane experiment.



Width of mixing layer, averaged over 100 runs. The mean is shown in red, and the grey bands show one and two standard deviations. Here λ is a characteristic length scale and $\tau = (\lambda/Ag)^{1/2}$ is the characteristic time scale. The simulations and the experiment were both run for 60ms, or approximately 2.7τ .

Comparison With Experiments

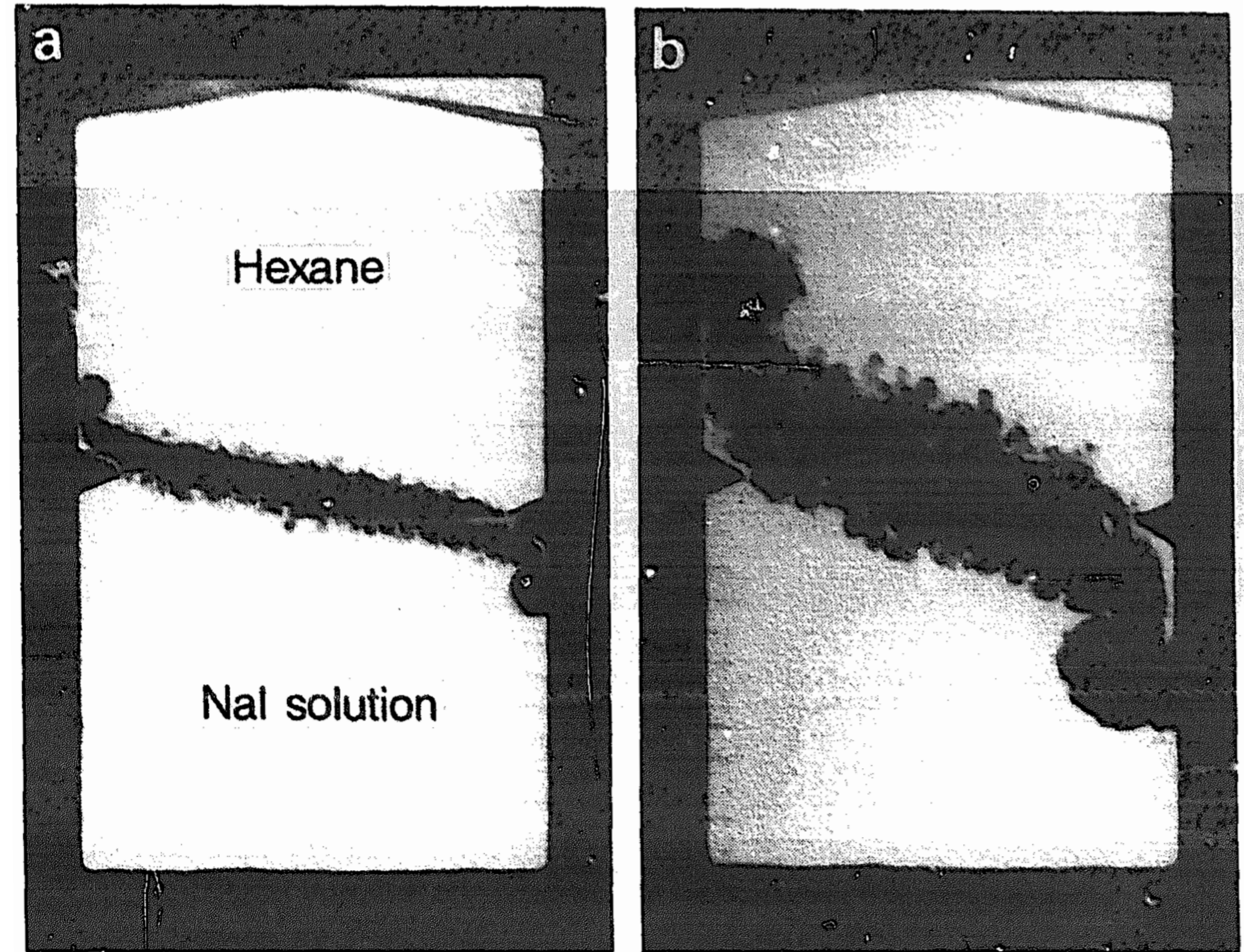
Tilted Rocket Rig

- This experiment is similar to the rocket rig, but the tank is tilted at a slope of 1/10.
- Our initial data has the form

$$z_3 = \frac{s_1}{10} + \text{Re} \sum_{k_1, k_2 = -N}^N c_{k_1, k_2} e^{2\pi i(k_1 s_1 + k_2 s_2)}$$

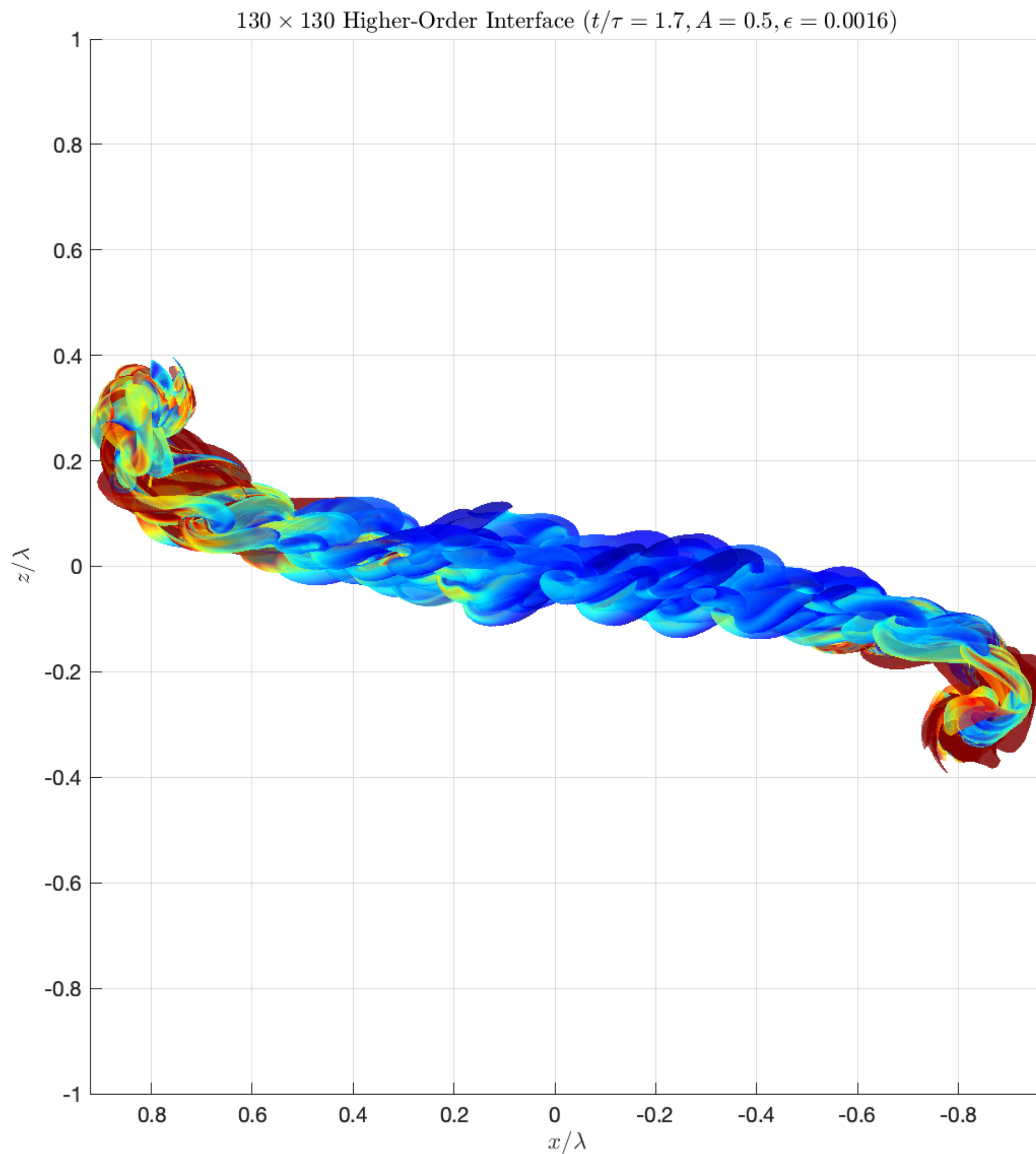
where the coefficients are again drawn from a normal distribution.

- The interface eventually flips over, so that the heavy fluid is on the bottom.

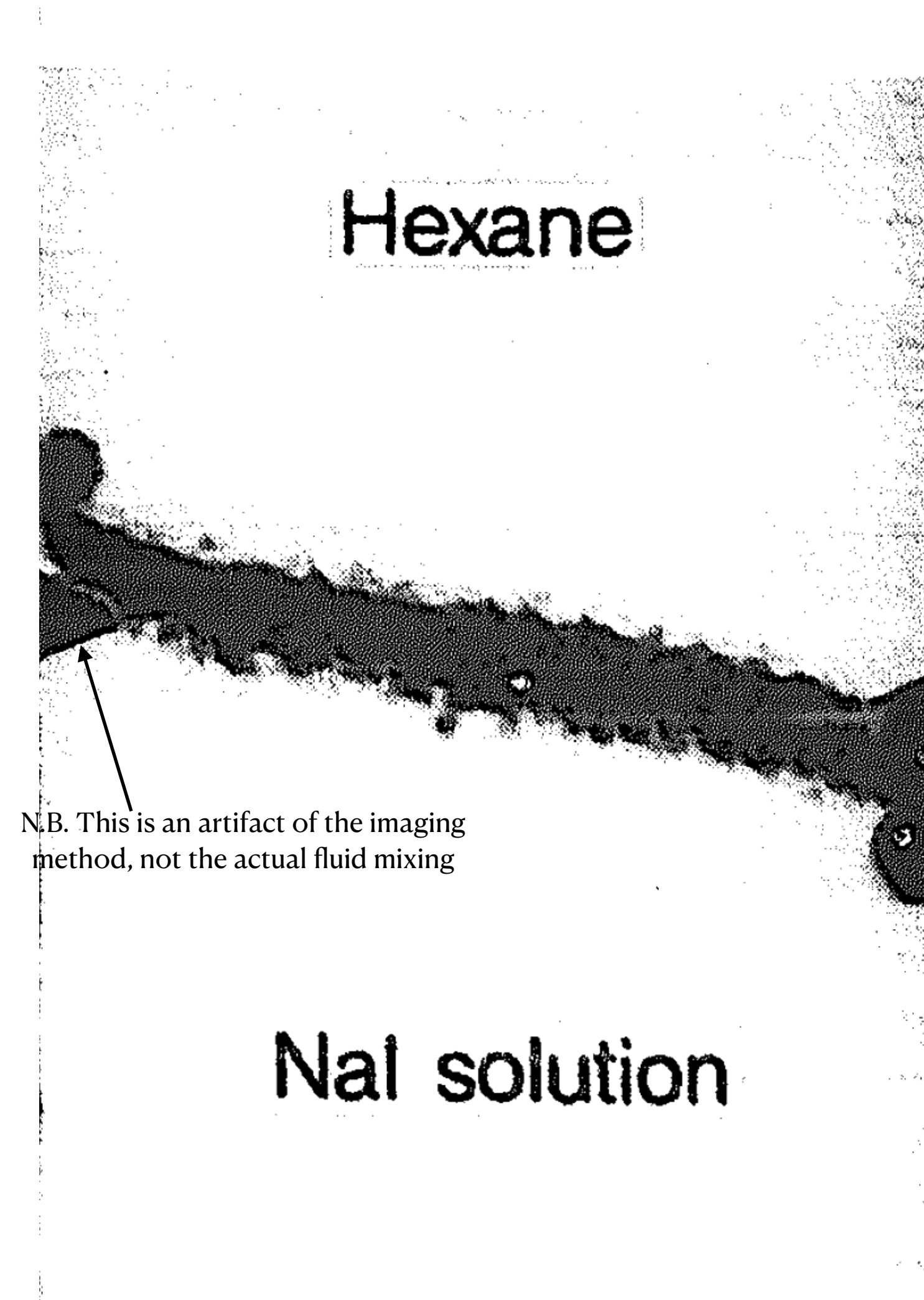


Experimental photographs by Youngs

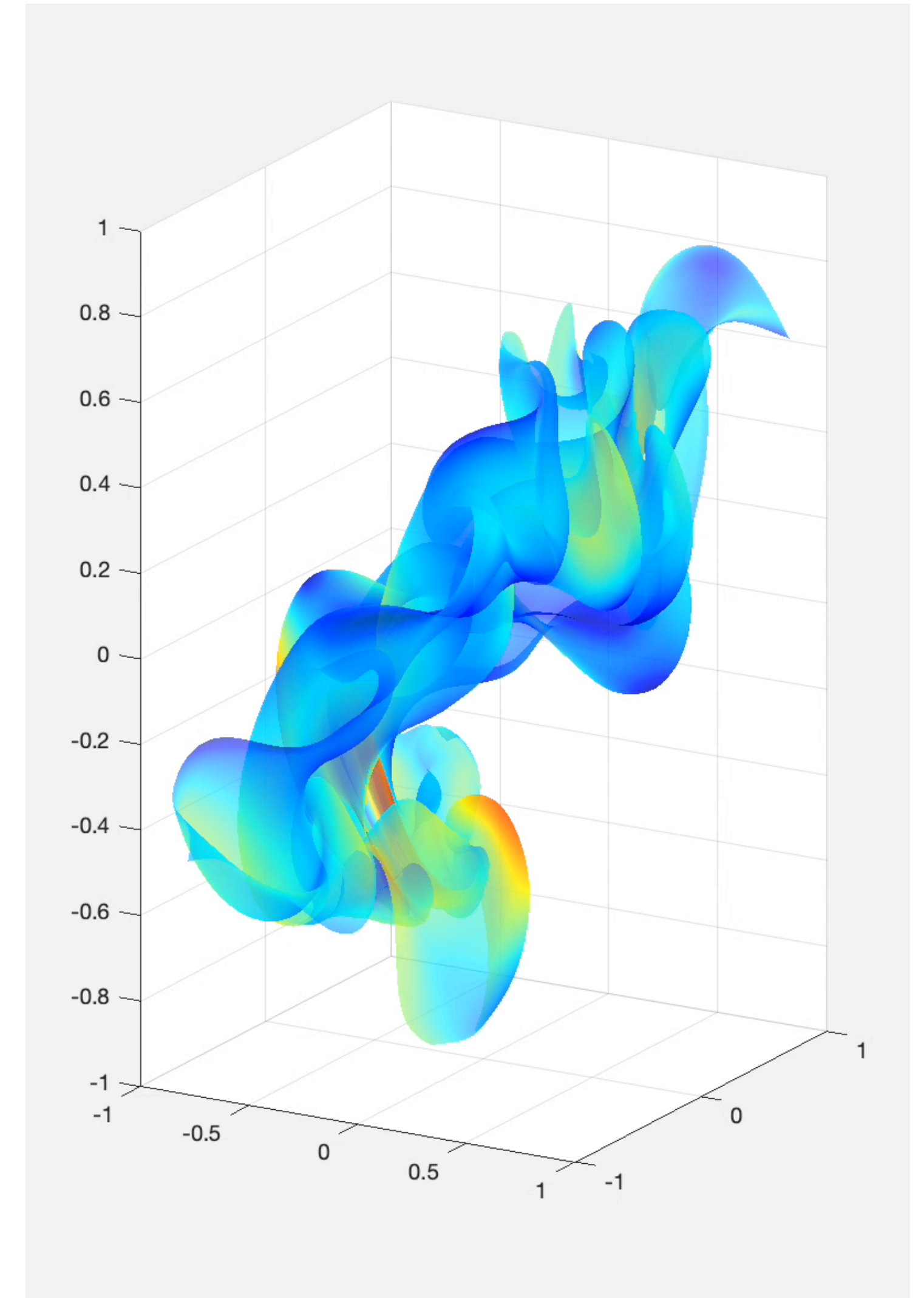
Comparison With Experiments



Side view of z -model at $t = 1.7\tau$



An experiment of Youngs, showing the tilted rig after 39ms, or approximately 1.7τ .

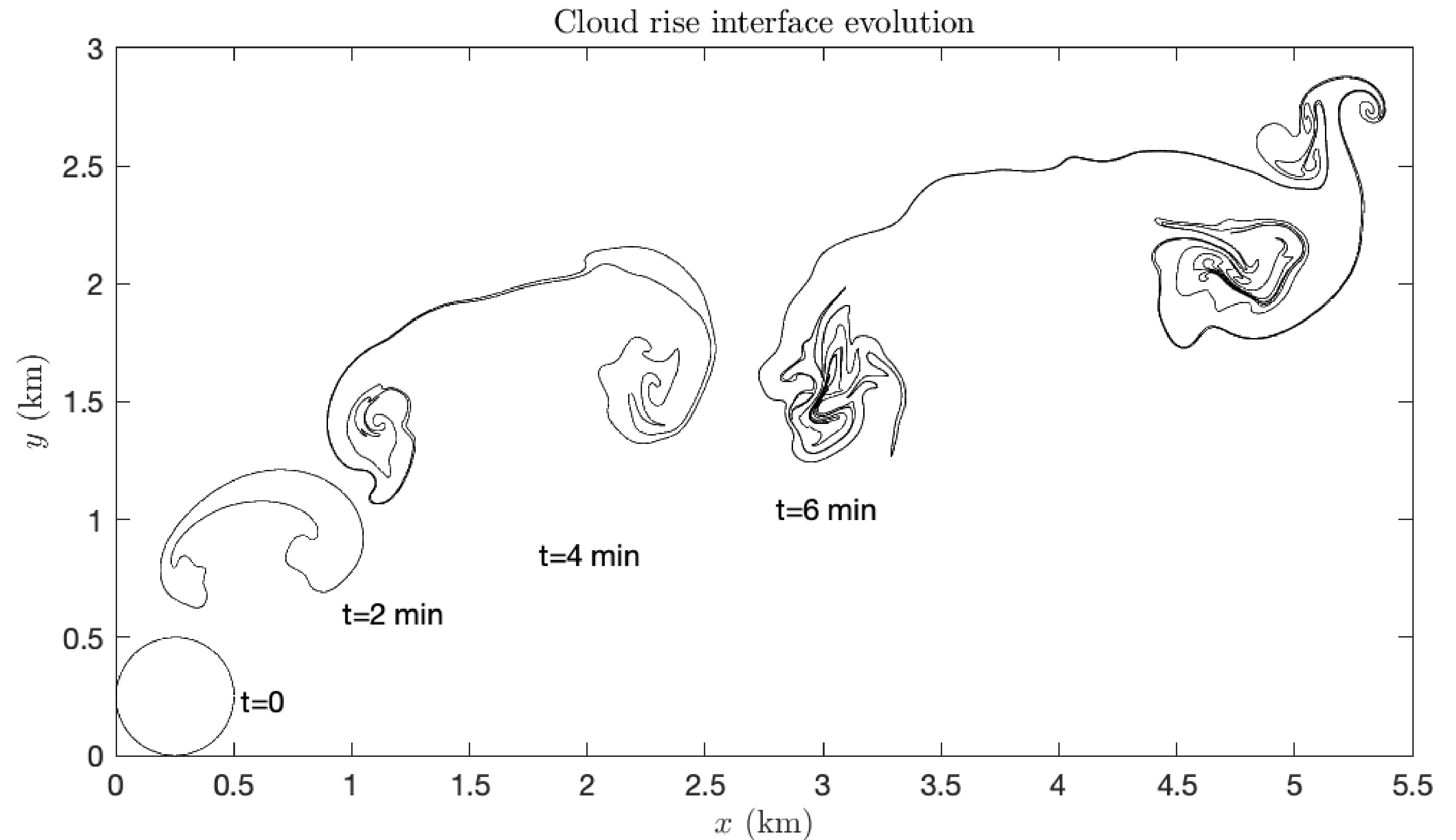


The tilted rig "flipping over" in our simulation.

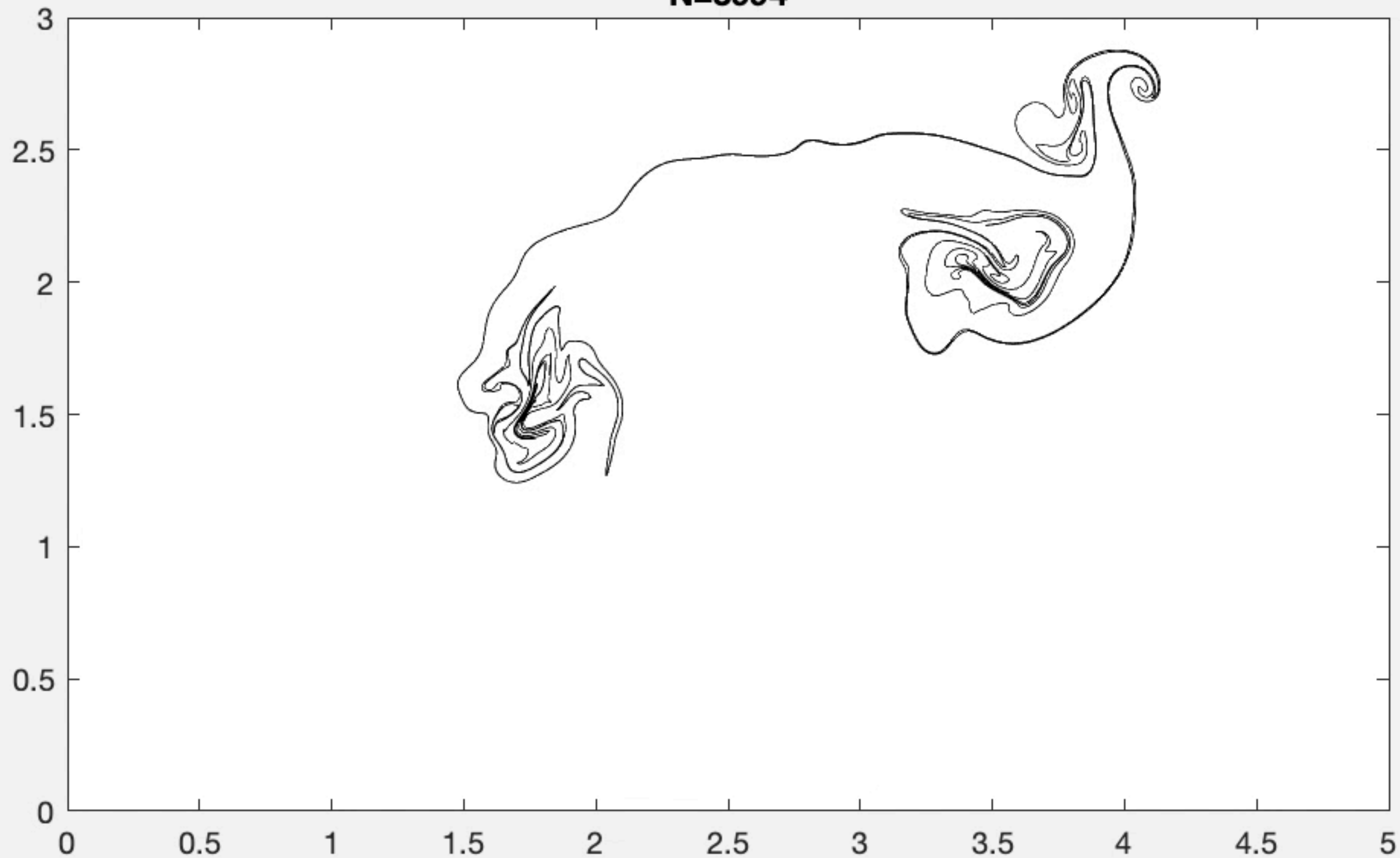
Cloud Rise

z-model + density stratification + shear flow

- A sphere of warm air rising into density stratification w/ horizontal wind shear.
- Baroclinic generation of vorticity dominates the flow, so we can model this using the z-model



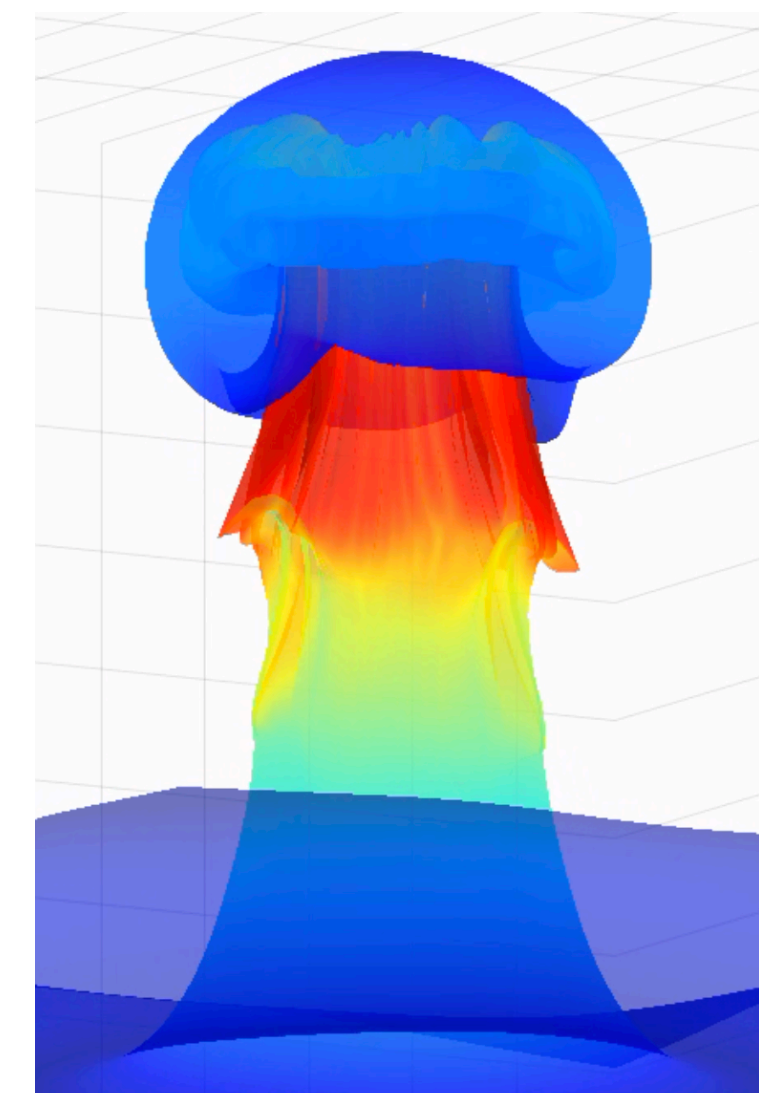
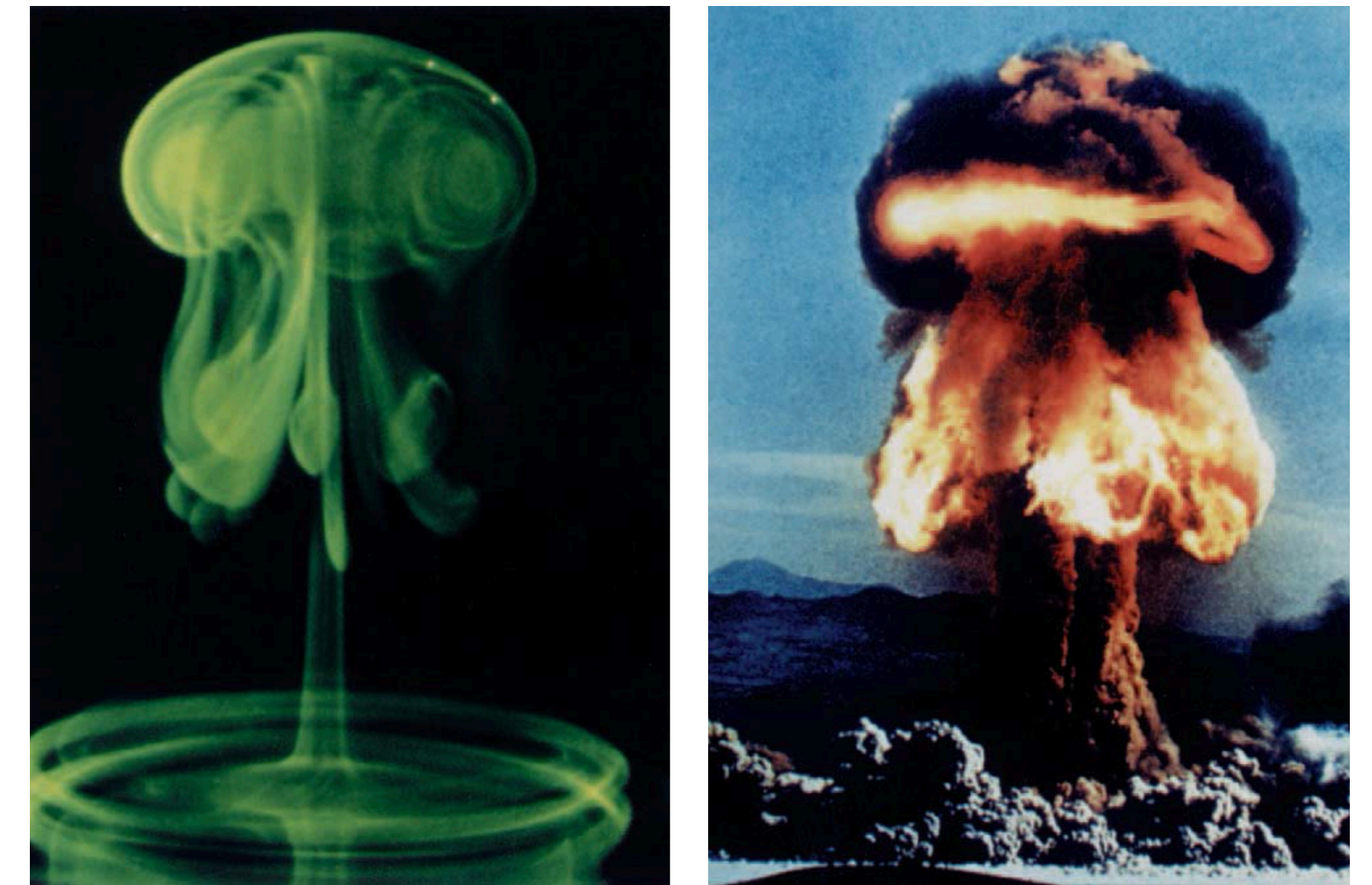
N=3994



Diffuse
Interface



Sharp
Interface



Rayleigh-Taylor Instability

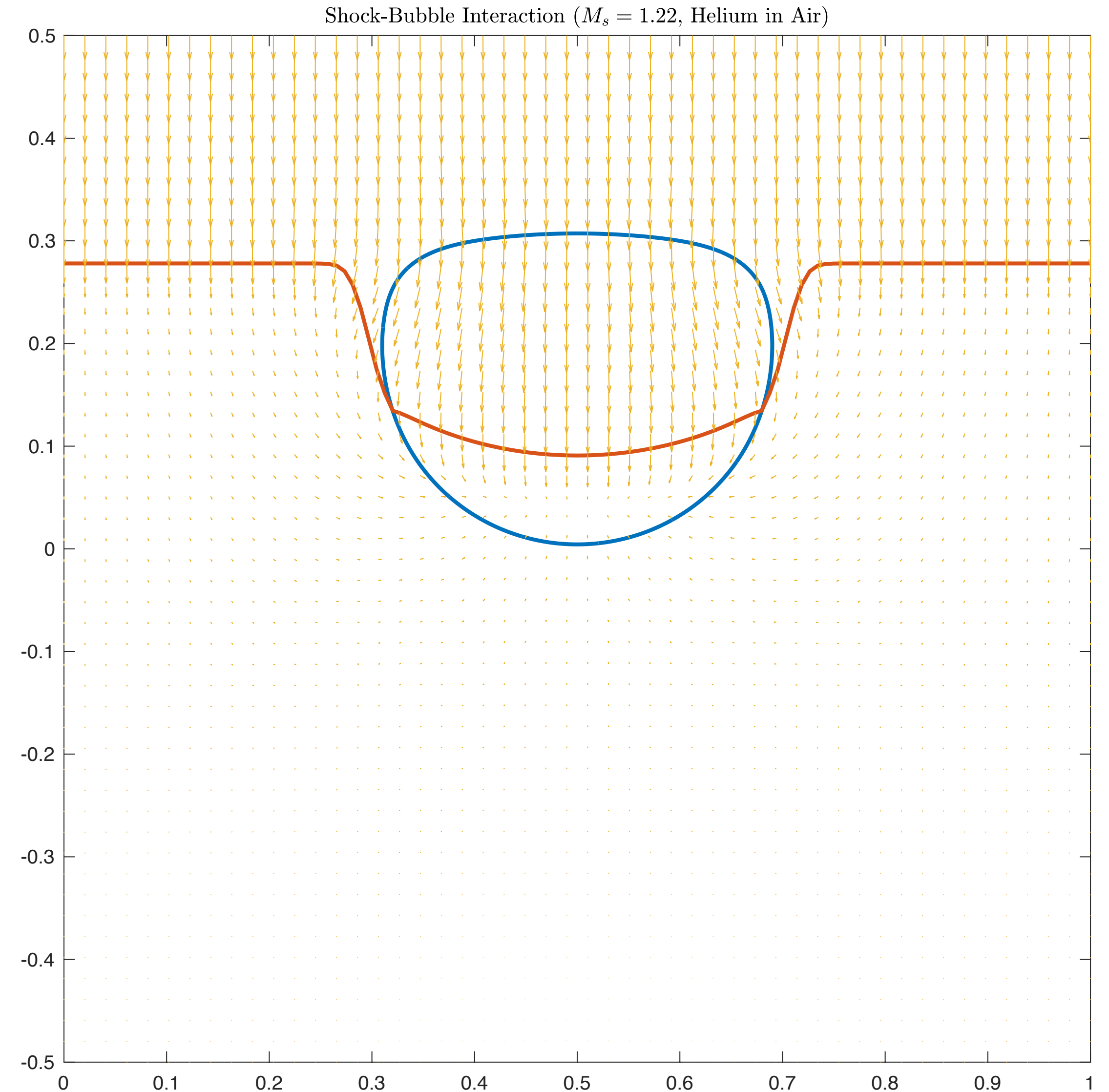
- 1) water droplet impact
- 2) plume of superheated gas from atomic bomb
- 3) z-model simulation

Rayleigh-Taylor instability in clouds vs. z-model

Shock-Contact interaction

Z-model + GSD

- Coupling contact model to compressible flow field for RMI (Ramani and Shkoller '20)
- Coupling contact discontinuities to geometrical shock dynamics (GSD), with full velocity reconstruction.
- Comparison with experiments of Haas/Sturtevant on shock-bubble interaction



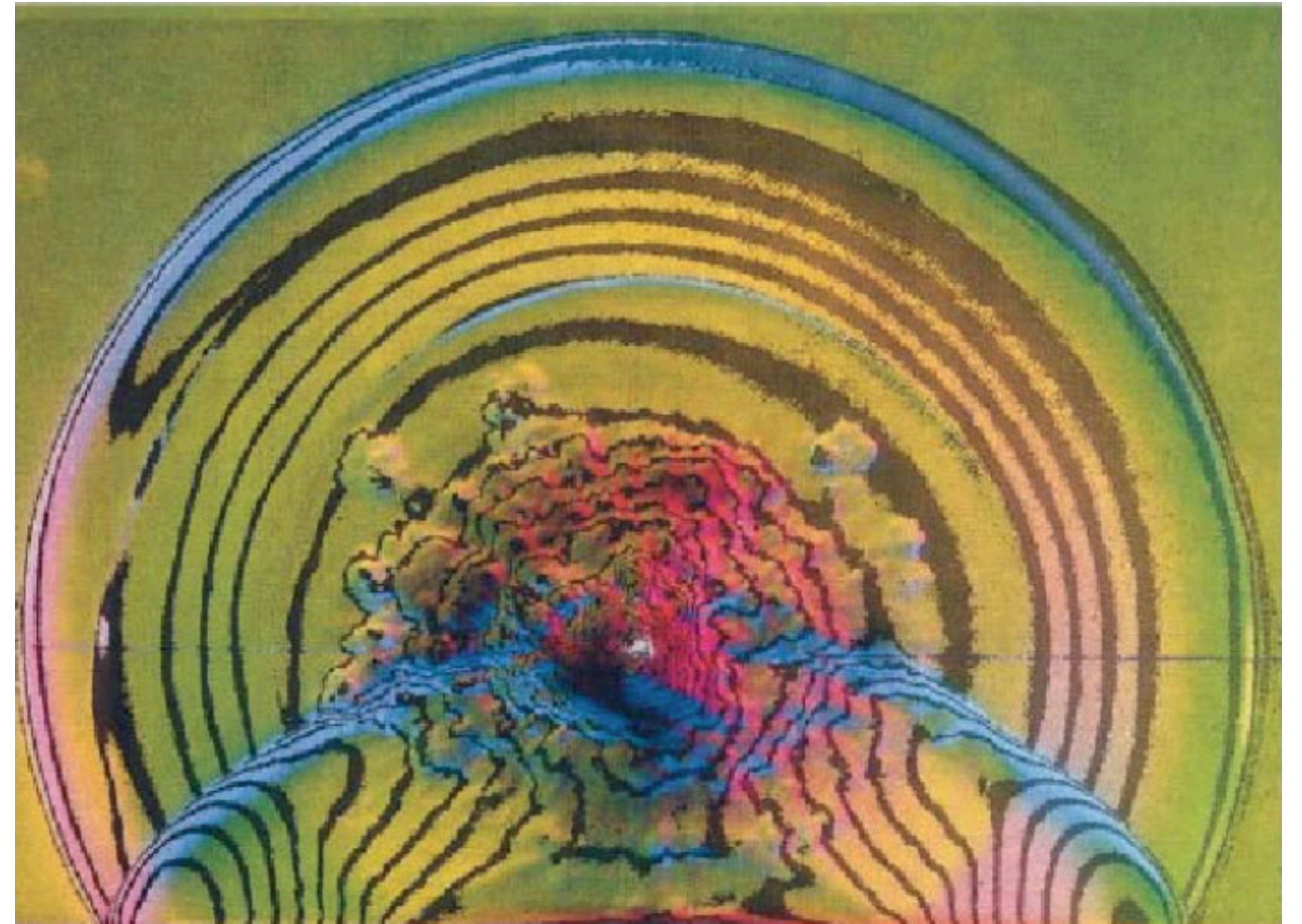
Shocks

Geometrical Shock Dynamics (Whitham, Schwendeman)

- Interface model for the shock position and shock Mach number

$$M = \frac{\text{shock speed}}{\text{sound speed}}.$$

- Models shock wave propagating into quiescent gas (e.g. blast wave from explosion, shock wave from supersonic airplane).



Schlieren/Interferometer image of blast wave (Kleine/Takayama)

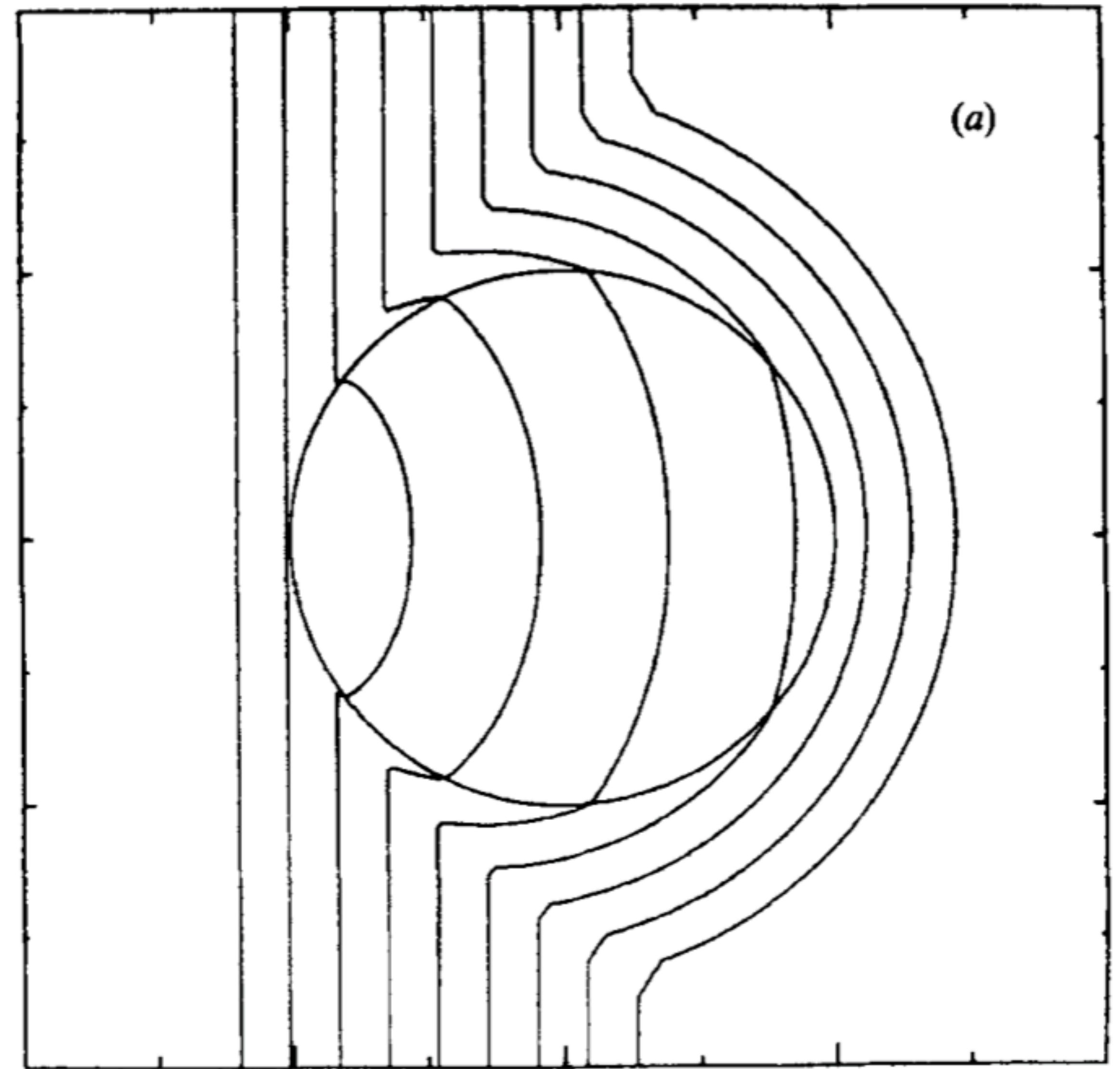
Shocks

Nonuniform Media

- When the sound speed changes, the Mach number changes too:

$$\Delta M = \frac{M \cos \theta}{\beta(M, \gamma)} \frac{\Delta c}{c}.$$

- In particular, where the shock encounters a density discontinuity, the shock develops a corner.



Shock propagating through sphere of lighter gas (Schwendeman)

Shocks and Contacts

Coupling the z-model and the kinematic shock model

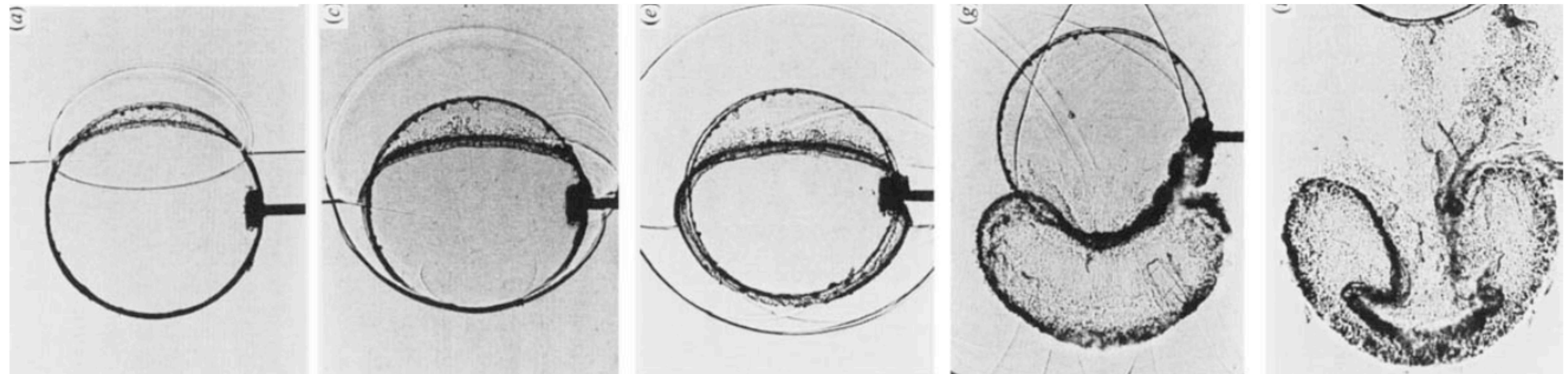
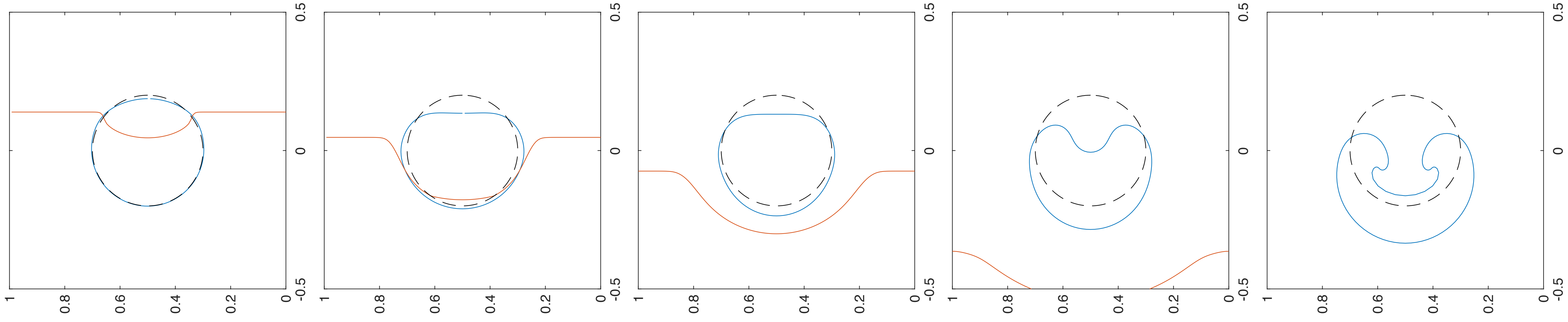
- The contact is affected by the shock via baroclinic vorticity production:

$$\Delta\omega = \left(\frac{1}{\rho^+} - \frac{1}{\rho^-} \right) \Delta p \sin \theta.$$

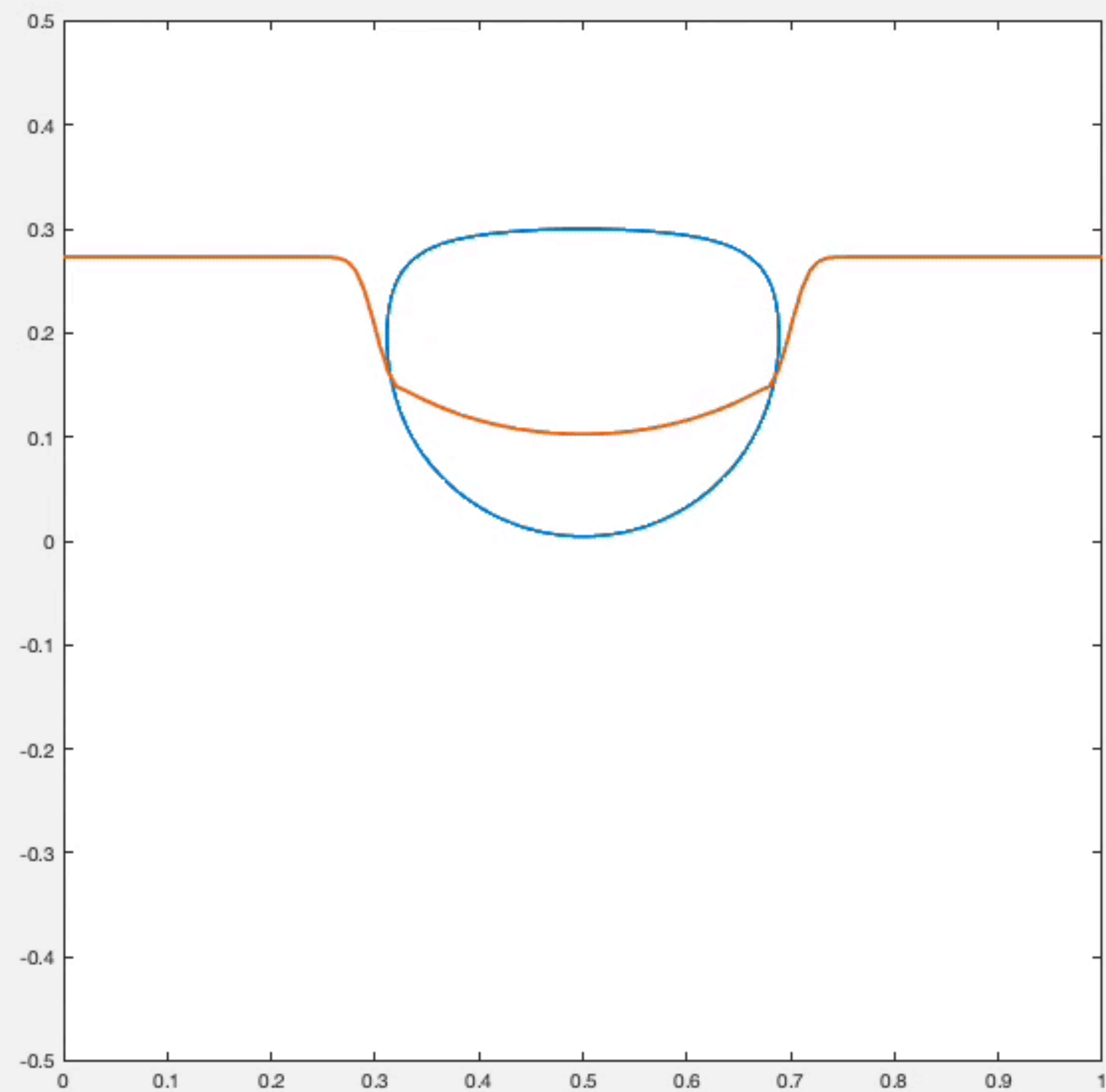
- The shock is affected by the contact via the change in sound speed:

$$\Delta M = \frac{M \cos \theta}{\beta(M, \gamma)} \frac{\Delta c}{c} \qquad c^2 = \frac{\gamma p}{\rho}.$$

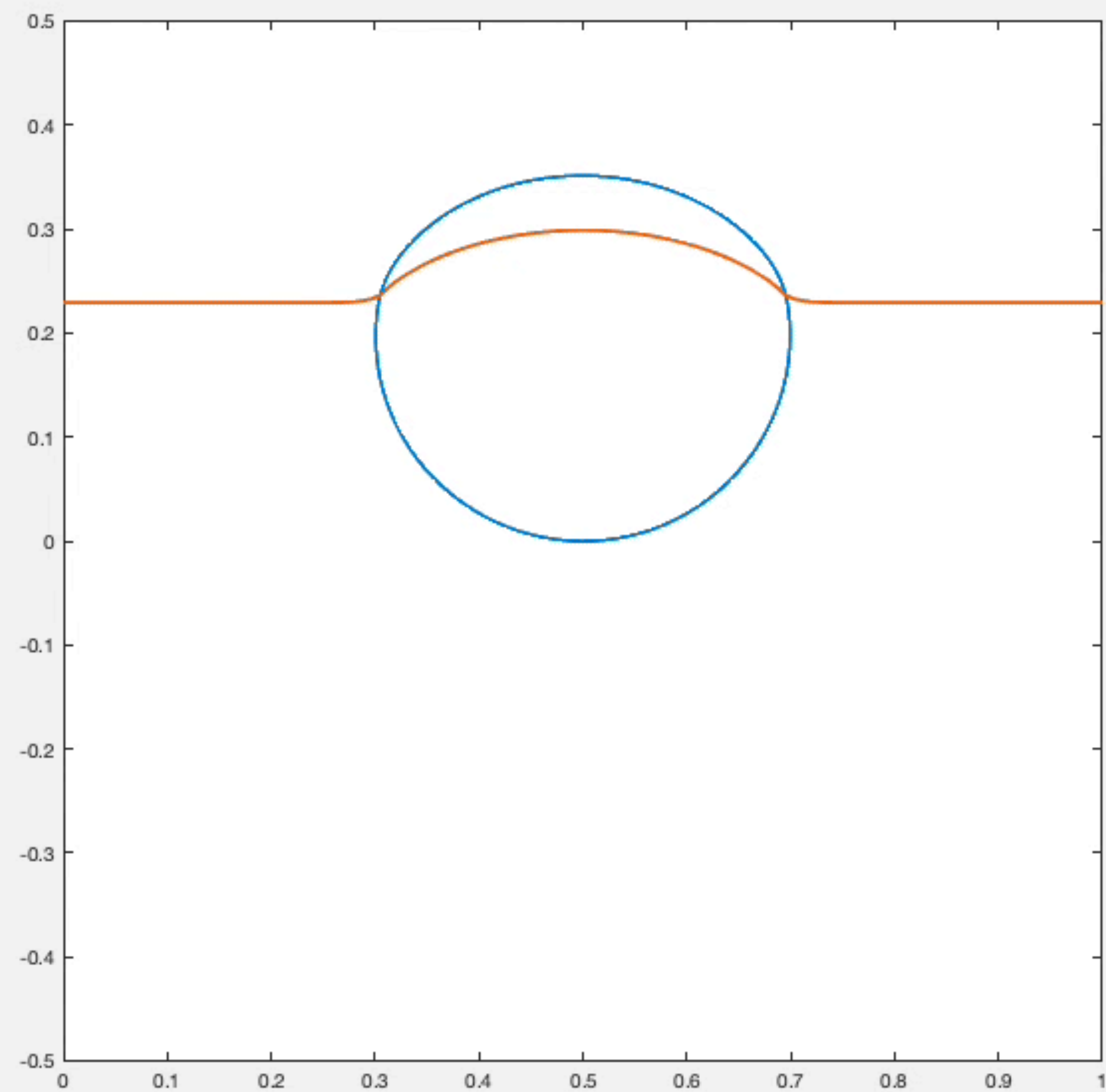
- The full velocity is reconstructed from the amplitudes of vorticity and compression, using our boundary integral formulation.



Mach 1.22 shock propagating through helium bubble in air;
our simulation vs. experiments of Haas/Sturtevant.



Air-Helium

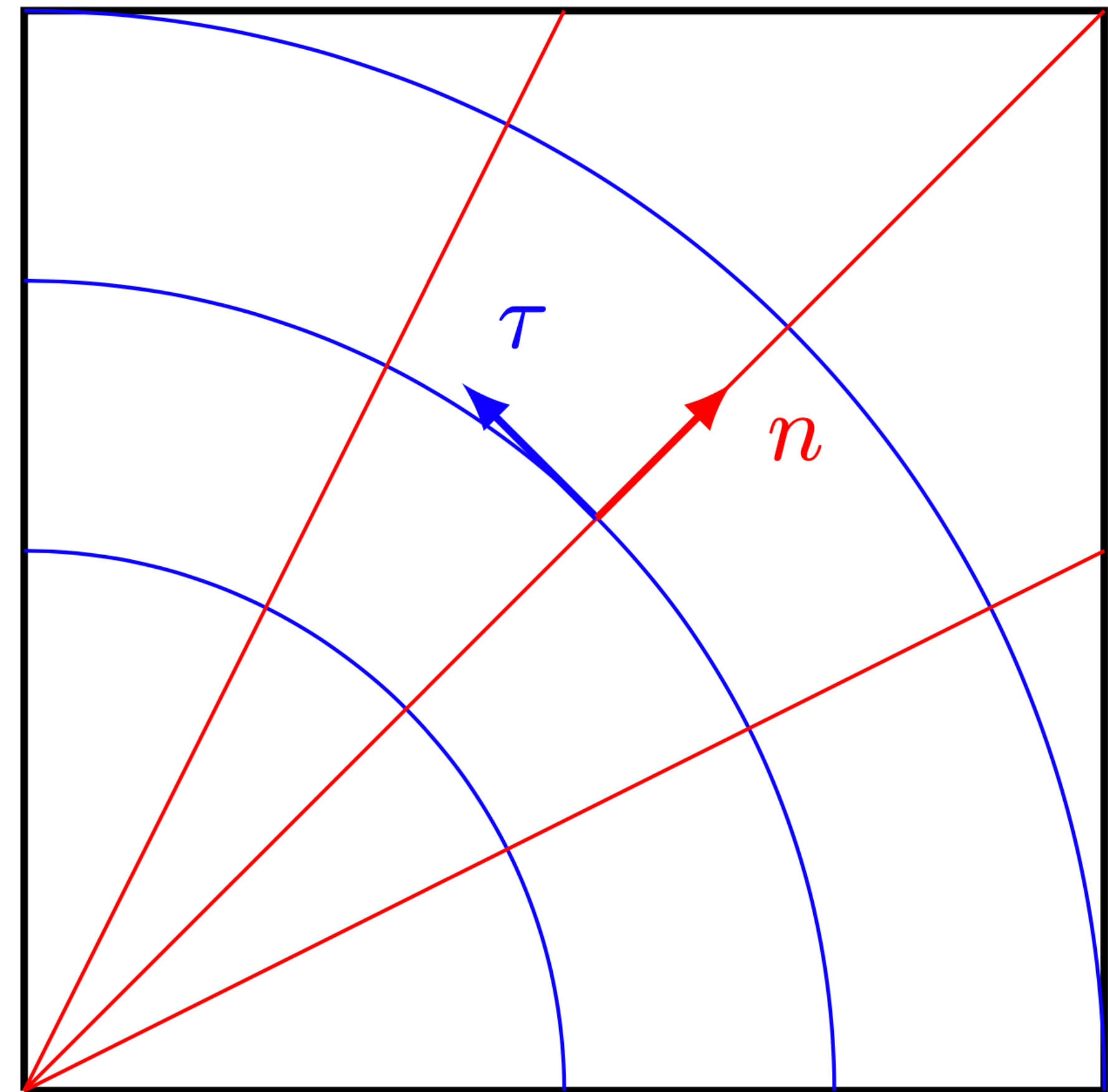


Air-Freon

Nonlinear Sound Waves

The wza-model

- Compressible Euler equations have three distinct characteristic velocities.
- Introduce new geometric coordinate system that moves with fast acoustic characteristic
- Simultaneously solve for both physical variables and geometric unknowns



Characteristic Surfaces

Fast Acoustic Characteristic

- Geometric coordinates

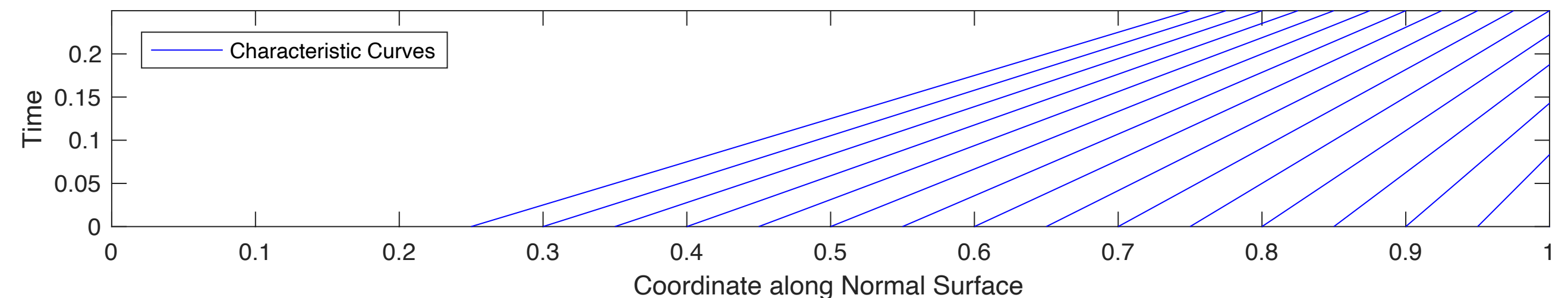
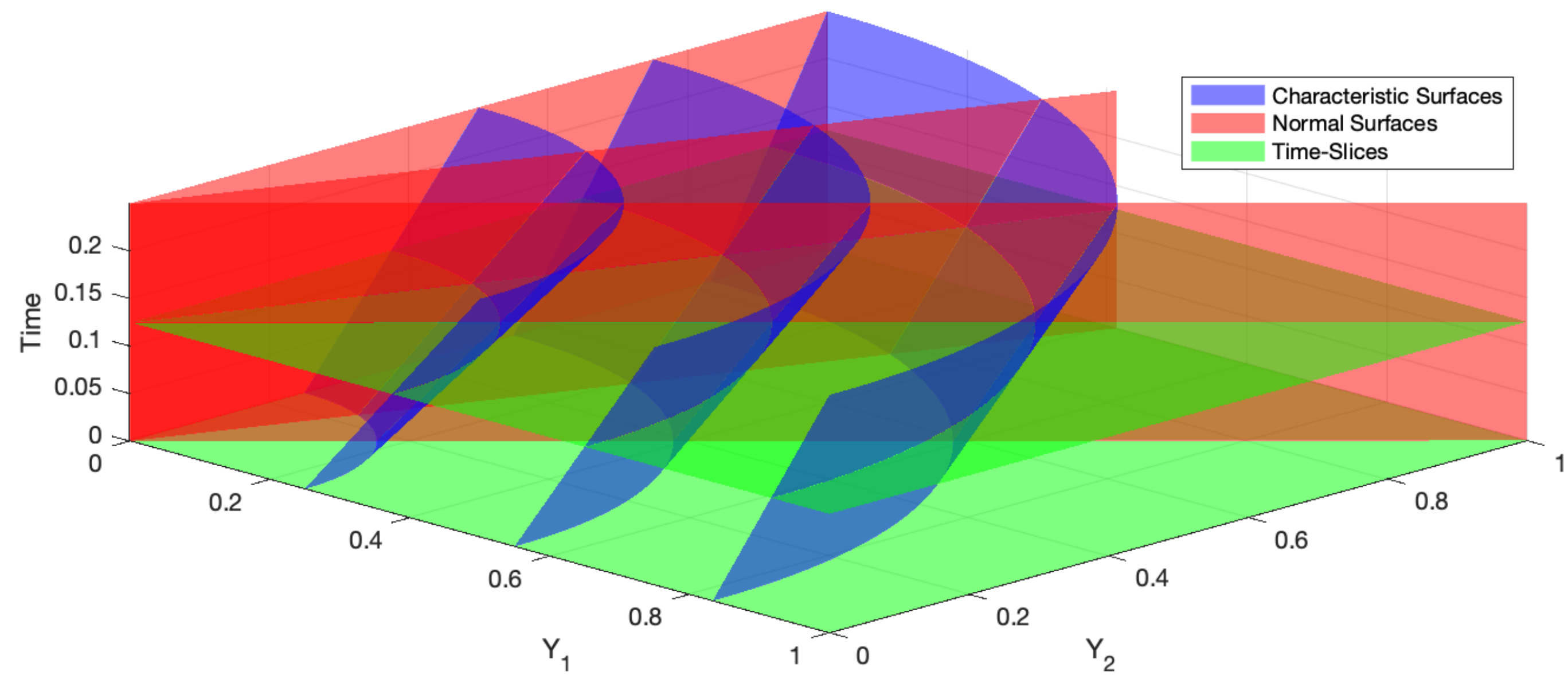
$$y_1 = \eta_1(x_1, x_2, t), \quad y_2 = \eta_2(x_1, x_2, t)$$

- Wavefronts given by $x_1 = \text{const.}$
- Tangent and normal directions

$$\tau = \frac{(\partial_2 \eta_1, \partial_2 \eta_2)}{|(\partial_2 \eta_1, \partial_2 \eta_2)|}, \quad n = \frac{(\partial_2 \eta_2, -\partial_2 \eta_1)}{|(\partial_2 \eta_1, \partial_2 \eta_2)|}$$

- Time-evolution given by

$$\partial_t \eta = u \circ \eta + (c \circ \eta)n$$



Geometric Riemann Variables

The $wz\alpha$ -model

- Riemann invariants given by $w = u \cdot n + c/\alpha$, $z = u \cdot n - c/\alpha$, $a = u \cdot \tau$.
- In 1d, w is constant along fast acoustic characteristic; in 2d, the evolution of w only depends on derivatives tangent to wavefront.

$$\left\{ \begin{array}{l} \partial_t \eta = a\tau + \left(\frac{1+\alpha}{2}w + \frac{1-\alpha}{2}z \right) n \\ \partial_t w = -a\partial_\tau \left(\frac{1+\alpha}{2}w + \frac{1-\alpha}{2}z \right) - \frac{\alpha}{2}(w-z)\partial_\tau a + \left(\frac{\alpha}{4}(w^2 - z^2) - a^2 \right) \kappa + a\partial_n n \cdot \tau \\ \partial_t z = -a\partial_\tau \left(\frac{1+\alpha}{2}w + \frac{1-\alpha}{2}z \right) + \frac{\alpha}{2}(w-z)\partial_\tau a - \left(\frac{\alpha}{4}(w^2 - z^2) + a^2 \right) \kappa + \alpha(w-z)\partial_n z - a\partial_n n \cdot \tau \\ \partial_t a = -\frac{\alpha}{4}(w-z)\partial_\tau(w-z) + \frac{1}{2}(w+z)\left(\frac{1+\alpha}{2}w + \frac{1-\alpha}{2}z \right) + \frac{1}{2}a(w+z)\kappa + \frac{\alpha}{2}(w-z)\partial_n a \end{array} \right.$$

Evolution of wavefront

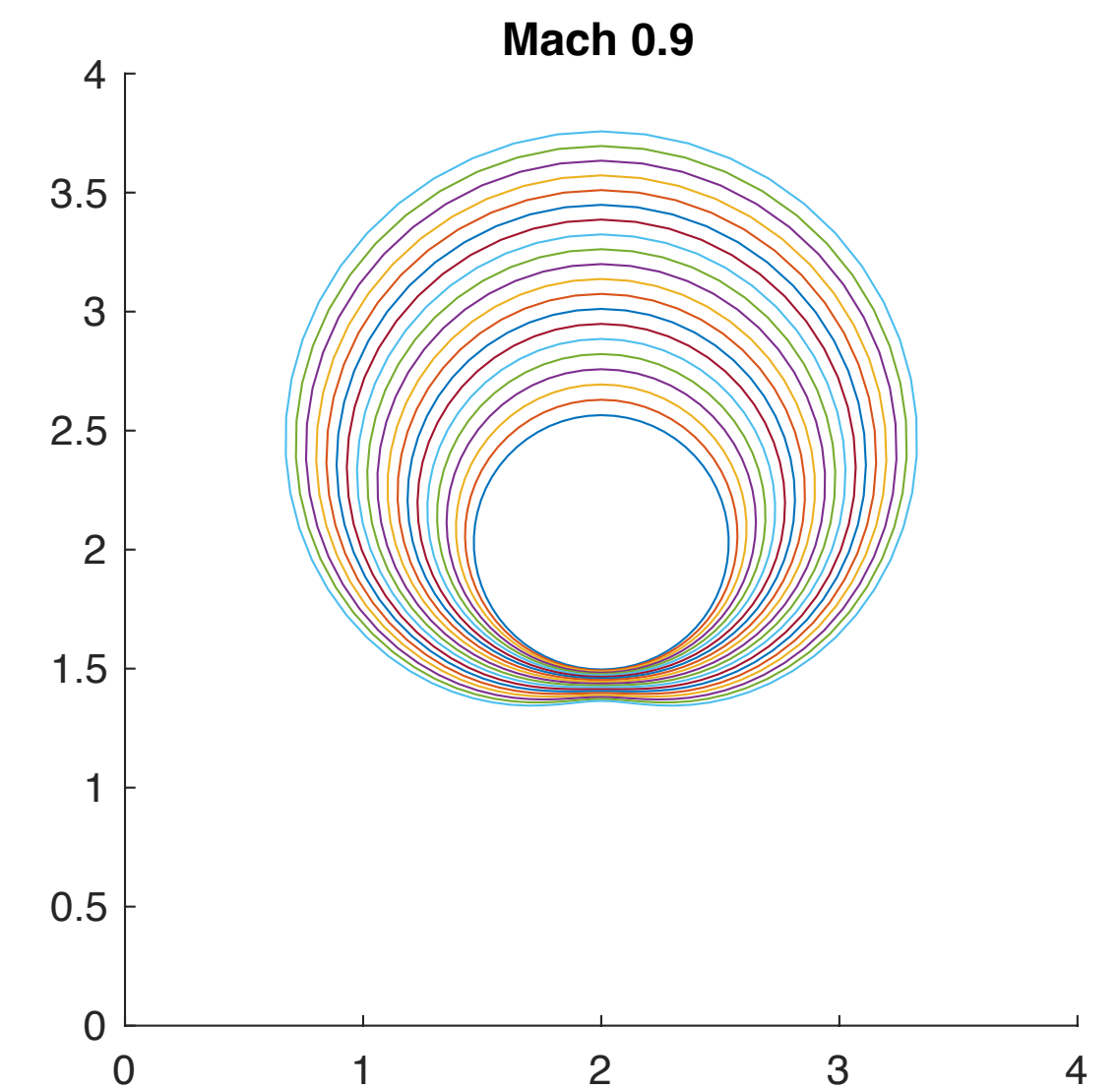
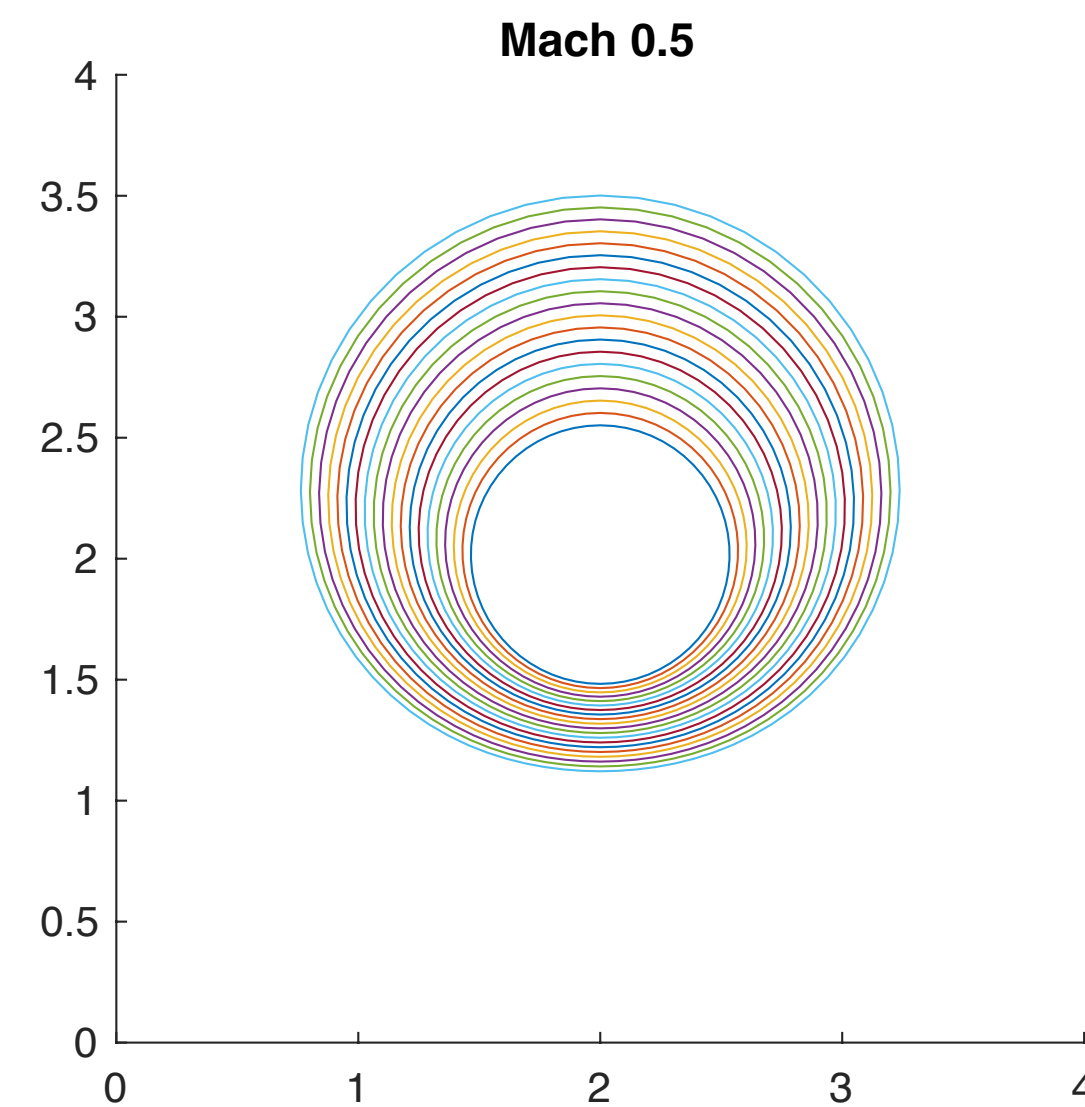
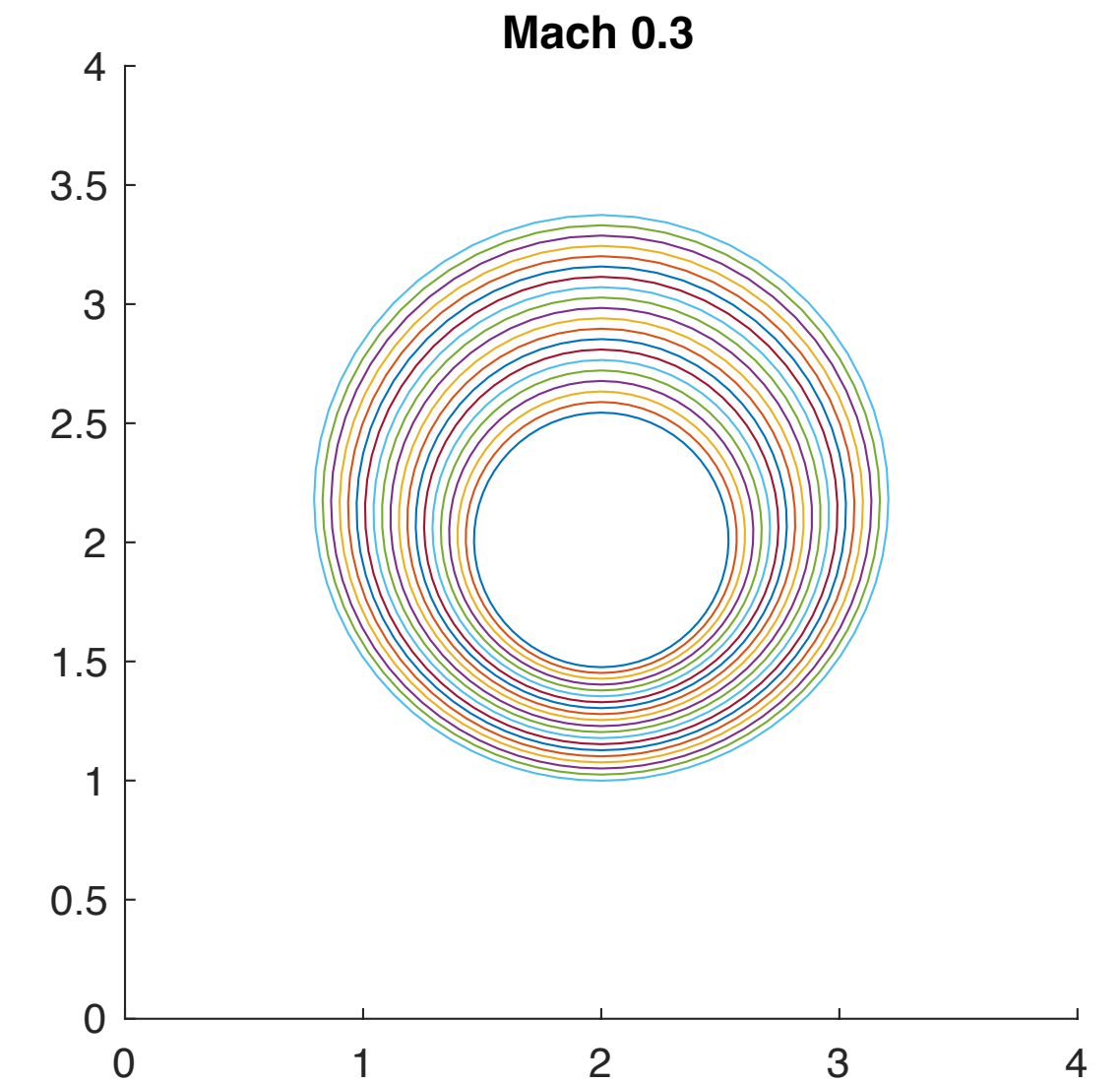
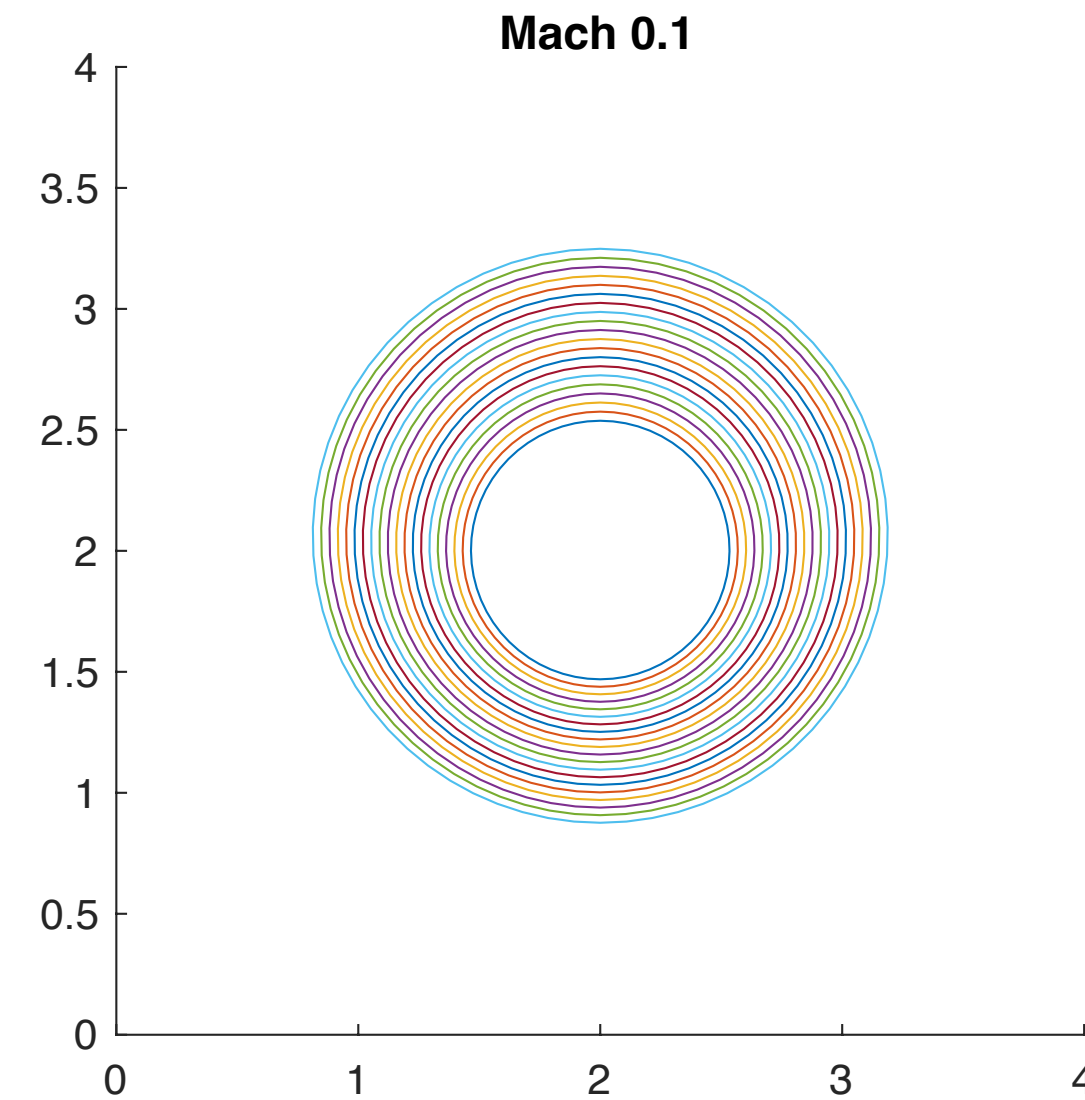
Evolution of geometric Riemann variables along wavefront

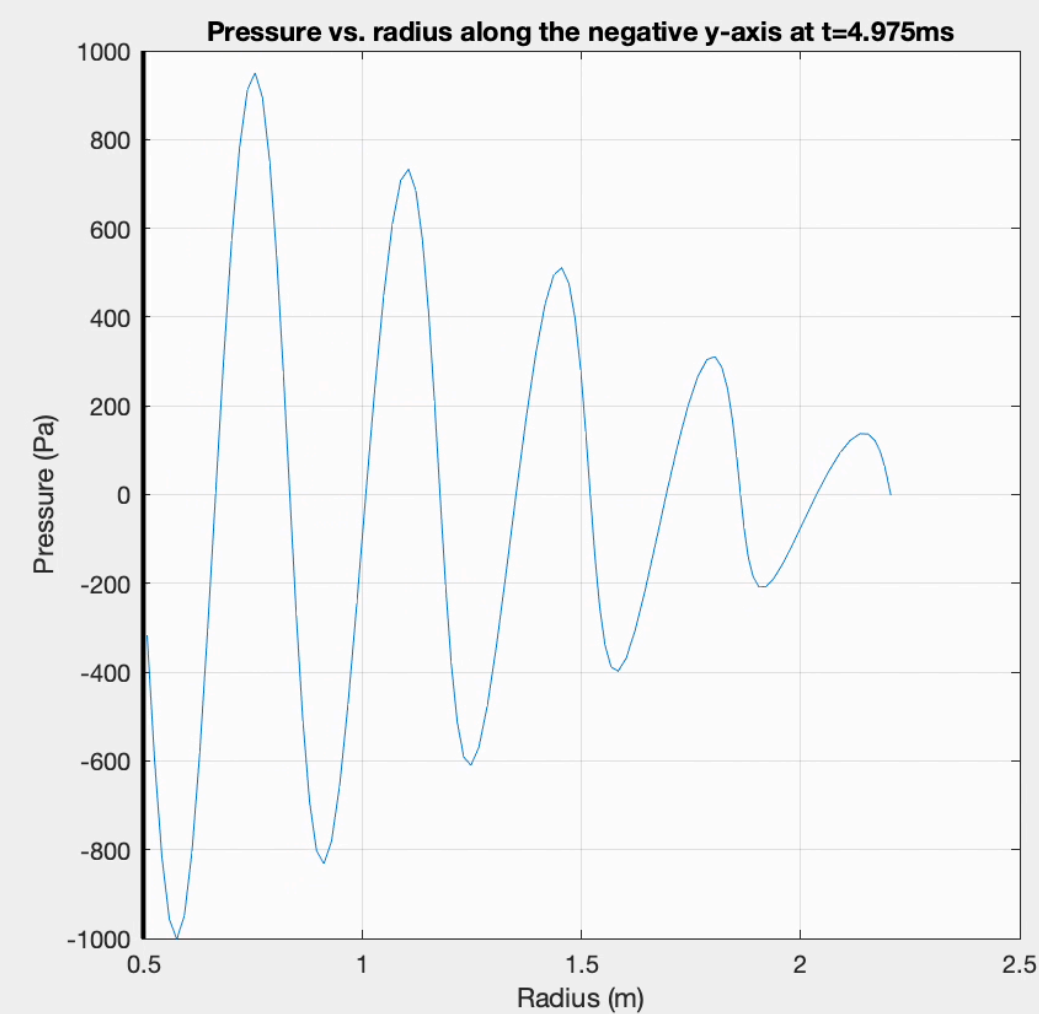
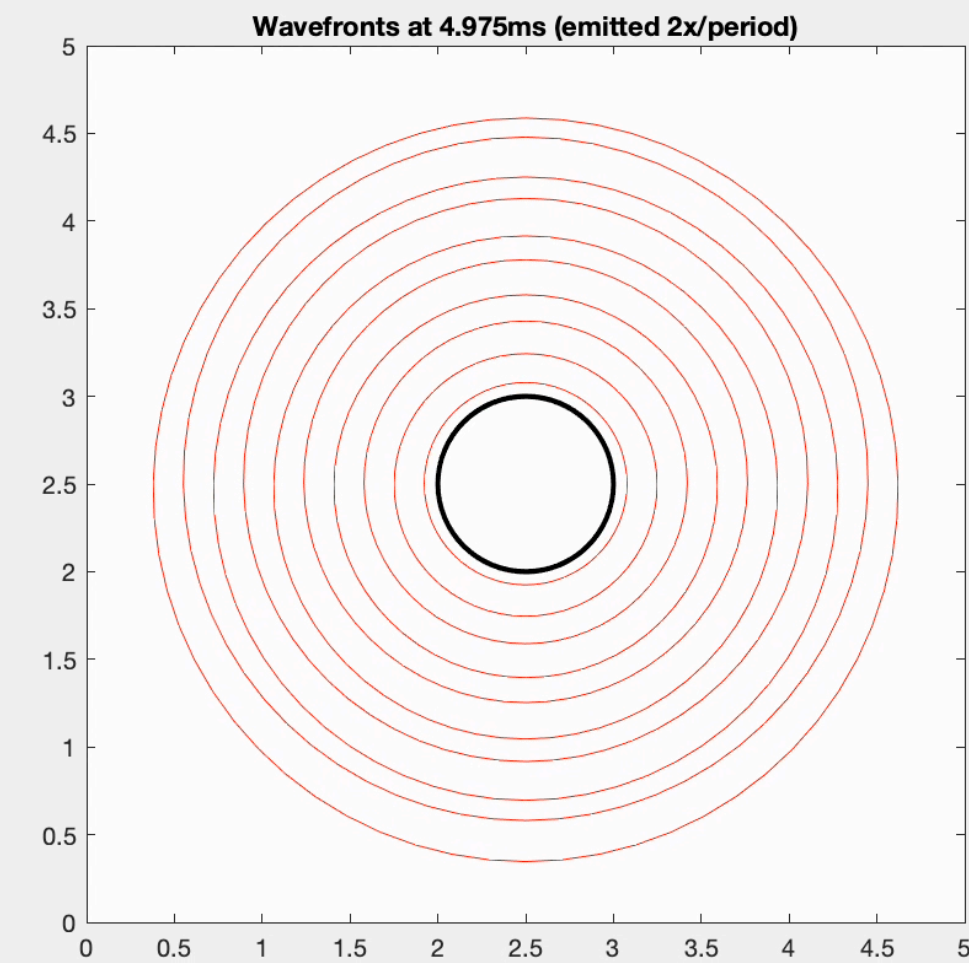
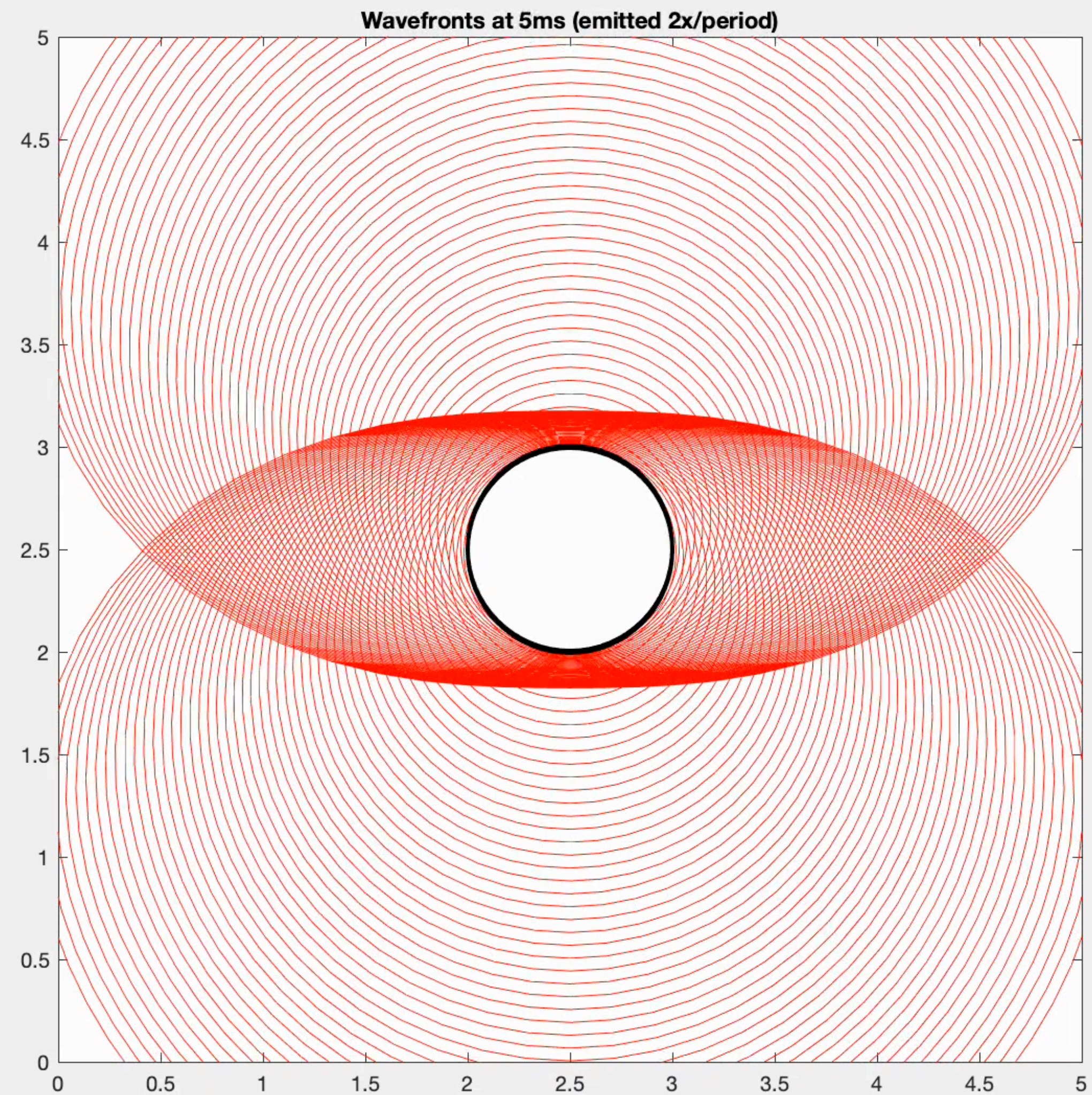
Normal derivatives of subdominant Riemann variables are negligible

Sound Waves

Interface Model

- Following the fast characteristic, we obtain an interface model for acoustic waves.
- Approaches geometrical optics in the low Mach number limit
- Tracks Riemann invariants along fast acoustic characteristics
- Captures nonlinear effects (steepening of acoustic waves into shock waves)





Continuous “emission” of wavefronts allows for reconstruction of sound wave profiles

Future Work

- Models in development:
 - Interface model for compressible vorticity waves and contact discontinuities.
 - Interface model for shock waves using our framework.
 - Steepening of nonlinear sound waves into shock waves.
 - Incompressible contact model including viscous effects
- Coupling of interface models:
 - Shock-contact interaction
 - Shock-sound wave interaction

A black and white artistic photograph. In the background, a person's face is shown in profile, looking towards the right. The face is softly lit, with the eyes closed or looking down. In the foreground, a liquid, possibly water or a light-colored oil, is captured in a dynamic splash or ripple, creating a series of peaks and valleys. The liquid is positioned horizontally across the middle of the frame, partially obscuring the lower part of the face. The overall mood is serene yet dynamic, with a focus on texture and light.

Thank You!