

# Embedded Eigenvalues of Coupled Graph Operators

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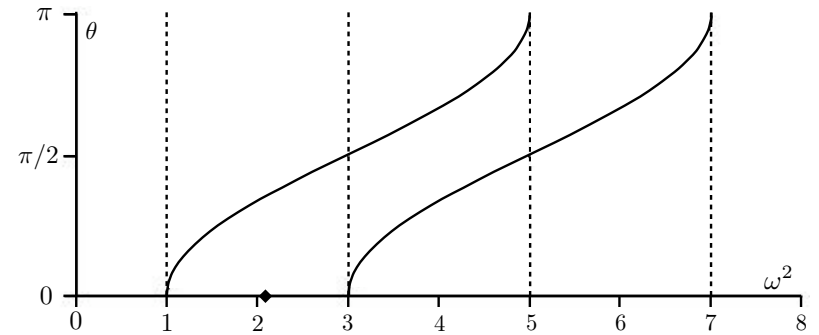
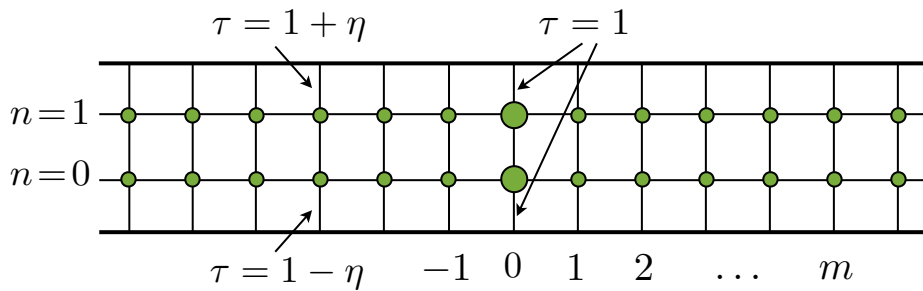


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# Spectrally embedded eigenfunctions in 1D graph operators

## 1. Coupled biased chains



Floquet solutions  $e^{i\theta m} q_n$  to  $(Au) = -\omega^2 u \rightsquigarrow$  Dispersion relation

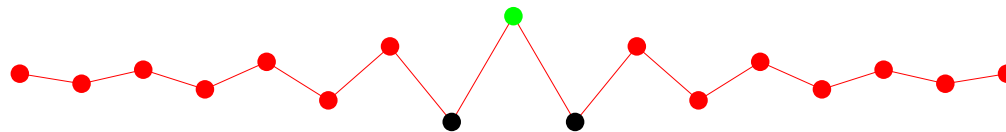
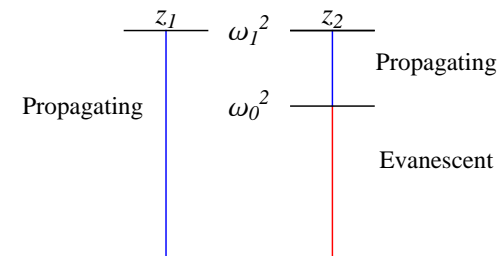
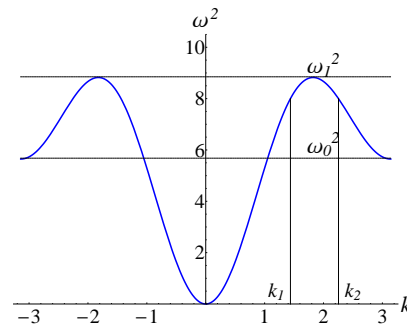
## 2. Next-nearest neighbor interactions $\rightsquigarrow$ bound state at embedded frequency

$$(Au)_n = \tau_2(u_{n+2} + u_{n-2}) + \tau_1(u_{n+1} + u_{n-1}) - 4u_n$$

Floquet solutions

$$u_n = z^n = e^{ikn}$$

of  $Au = -\omega^2 u$



## Higher-D: Embedded eigenvalues of periodic graph operators

Given a  $n$ -periodic finite-order graph operator  $A$ , its **Floquet surface** at  $\lambda \in \mathbb{C}$  is

$$\Psi_{A,\lambda} = \{ 0 \neq z \in \mathbb{C}^n : \det(\hat{A}(z) - \lambda) \neq 0 \}$$

$\hat{A}(z)$  is the multiplication-operator representation of  $A$  under the Floquet transform

$$u_n \mapsto \hat{u}_n(z) = \sum_{g \in \mathbb{Z}^n} u_{n+g} z^{-g}$$

The **spectrum** of  $A$  is

$$\sigma(A) = \{ \lambda : \Psi_{A,\lambda} \cap \mathbb{T}^n \text{ is nonempty.} \}$$

**Theorem.** (Kuchment and Vainberg)

Let  $\lambda \in \sigma(A)$ , and let the Floquet surface  $\Psi_{A,\lambda}$  be **irreducible**. If  $u$  is in  $\ell^2$ ,  $B$  is a local perturbation of  $A$ , and

$$(A + B)u = \lambda u,$$

then  $u$  **has compact support**.

P. Kuchment and B. Vainberg, [On the Structure of Eigenfunctions Corresponding to Embedded Eigenvalues of Locally Perturbed Periodic Graph Operators](#), Commun. Math. Phys. 268 (2006)

# A class of operators with reducible Fermi surface

$A$  and  $L$  : periodic self-adjoint operators on an  $n$ -periodic graph.

$$\begin{aligned} \text{bias} & : B = \cos(\theta) L \\ \text{coupling} & : \Gamma = e^{i\phi} \sin(\theta) L \end{aligned} \implies L^2 = B^2 + \Gamma\Gamma^*$$

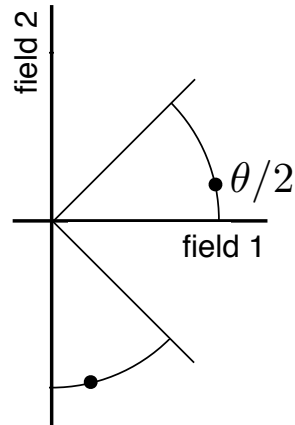
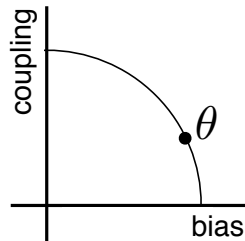
$$\mathcal{A} = \begin{bmatrix} A+B & \Gamma \\ \Gamma^* & A-B \end{bmatrix}, \quad \tilde{\mathcal{A}} = \begin{bmatrix} A+L & 0 \\ 0 & A-L \end{bmatrix}, \quad \mathcal{U} = \begin{bmatrix} e^{i\phi} \cos(\frac{\theta}{2})I & -\sin(\frac{\theta}{2})I \\ \sin(\frac{\theta}{2})I & e^{-i\phi} \cos(\frac{\theta}{2})I \end{bmatrix}$$

$$\implies \mathcal{A}\mathcal{U} = \mathcal{U}\tilde{\mathcal{A}} \quad \text{with } \mathcal{A} \text{ and } \tilde{\mathcal{A}} \text{ self-adjoint and } \mathcal{U} \text{ unitary.}$$

$$\det(\hat{\mathcal{A}}(z) - \lambda I) = \det(\hat{A}(z) + \hat{L}(z) - \lambda I) \det(\hat{A}(z) - \hat{L}(z) - \lambda I)$$

$$\implies \Psi_{\mathcal{A},\lambda} \text{ is reducible for all } \lambda.$$

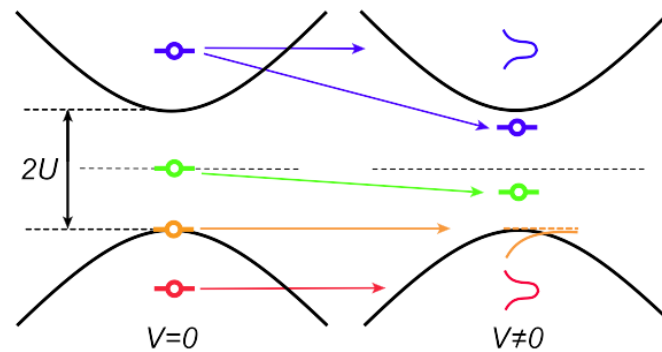
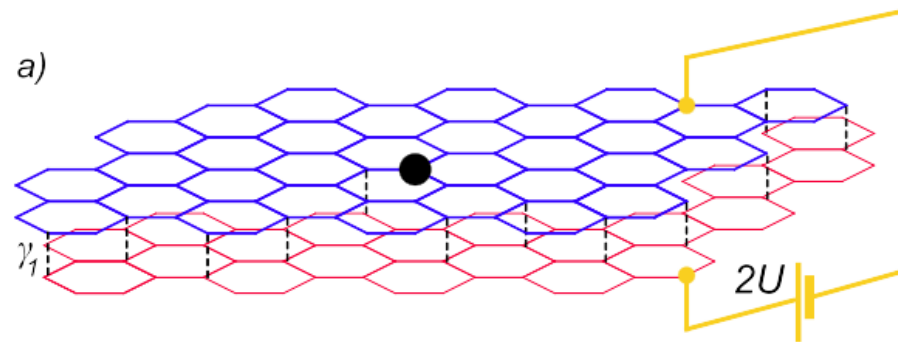
Interpretation:



field components of the columns of  $\mathcal{U}$

$\theta = \pi/2 \rightsquigarrow$  even and odd motions

## Example: Bilayer graphene



V. V. Mkhitarian and E. G. Mishchenko, [Localized states due to expulsion of resonant impurity levels from the continuum in bilayer graphene](#), Phys. Rev. Lett. 110, 086805 (2013)

## Forcing an evanescent motion within the spectrum

$$\tilde{\mathcal{A}}u = f$$

decompose  $u$  and  $f$  into decoupled components:  $u = u_+ + u_-$      $f = f_+ + f_-$

$$(A \pm L - \lambda I)u_{\pm} = f_{\pm}$$

Choose  $\lambda$  such that  $\det(\hat{A}(z) + \hat{L}(z) - \lambda I)$  has no zeros on the torus  $\mathbb{T}^n$   
and  $\det(\hat{A}(z) - \hat{L}(z) - \lambda I)$  does have zeros on  $\mathbb{T}^n$

So  $\lambda \in \sigma(\mathcal{A})$  but  $\lambda \notin \sigma(A + L - \lambda I)$ .

$$\hat{u}_+(z) = (\hat{A}(z) + \hat{L}(z) - \lambda I)^{-1} \hat{f}_+(z)$$

$$\hat{u}_-(z) = 0$$

## Evanescent spectrally embedded eigenfunctions (not compactly supported)

Localized perturbation of  $A$ :  $A + V$

To construct an embedded eigenvalue for the perturbed operator

$$\mathcal{A} + \mathcal{V} = \begin{bmatrix} A + V + B & \Gamma \\ \Gamma^* & A + V - B \end{bmatrix},$$

construct a non-embedded one for  $A + V + L - \lambda I$   
that is in the spectrum of  $A + V - L - \lambda I$ .

Want  $(\hat{A}(z) + \hat{L}(z) - \lambda I)\hat{u}_+(z) = -\hat{V}(z)\hat{u}_+(z) =: \hat{f}_+(z)$

**Question:** Given  $f_+$  and the solution  $u_+$ ,  
can one find a multiplication operator  $(Vu)_n = v_n u_{+n}$  such that  $v_n u_n = f_{+n}$  ?

One needs  $u_{+n} \neq 0$  for all  $n$  such that  $f_{+n} \neq 0$ .

This is not always possible—but it is if the forcing is localized at one point:  $f_{+n} = \delta_{0n}$ .