

Introduction to General Relativity

Math 280, UC-Davis - Temple

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- Grading - Hand in FIP Homework problems
- Office Hrs: T-Th after class / 3148 MSB/Appt
- Text: Introduction to General Relativity
by Adler, Bazin & Schiffer

Out of Print - copy of text together with
brief article on history on file in Math Off.

No Textbook Required / Notes should be self-contained

- Articles & Notes can be found on my
webpage: <http://www.math.ucdavis.edu>

References:

- (1) R. Wald, General Relativity (Standard)
- (2) S. Weinberg, Gravitation & Cosmology:
Principles and applications of GR (Great physics)
- (3) Misner, Thorne & Wheeler, Gravitation
("Telephone book" - full of geometrical & physical insights)
- (4) Hawking & Ellis, Large Scale Structure of Space time (short on intuition / ch 2 gives brief intro to diff. geom / ch 6-10 covers singularity theorems)
- (5) S. Chandrasekhar, The Mathematical Theory of Black Holes (Good advanced book)
- (6) B. Dubrovin & A. Fomenko Modern Geometry Methods & Applications Vol I (Great Exposition)

- (7) Ø Grøn & S. Hervik Einstein's General (3A)
Theory of Relativity (Great book/advanced)
- (8) L. Hughston & K. Tod An Introduction to
General Relativity (Excellent introduction)
- (9) Sachs & Wu, General Relativity for
Mathematicians (Avoid!)
- (10) Peebles, Cosmology (standard)
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Famous Classical Texts:

- (1) Theory of Relativity, Eddington '20's
- (2) Relativity, Thermodynamics and Cosmology
R. Tolman, Oxford Press, 1934
- (3) Whittaker ≈ 1950 Credited Spec. Rel. to Lorentz & Poincaré.

DIFFERENTIAL Geometry

- (1) Differential Geometry, Vol I, II SPIVAK
- (2) Foundations of Differentiable Manifolds & Lie Groups F. W. Warner

History

- (1) Subtle is the Lord: the science and life of Albert Einstein, A. Paes

Special Relativity

- (1) Spacetime Physics Taylor and Wheeler

Introduction

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■ The Einstein Equation (1916):

$$(1) \quad G = K T$$

- First and simplest ^{- covariant} field equation for gravitational field
- Replaced Newton's Law of Gravitation (1687)

$$(2) \quad \vec{F} = -G_0 \frac{M_1 M_2}{r^2} \vec{r}$$

$$\begin{array}{ccc} & G = K T & \\ \nearrow & & \nwarrow \\ \text{Einstein} & & \text{Stress Energy Tensor} \\ \text{Curvature Tensor} & & \text{(Einstein)} \end{array}$$

Here :

$$K = \frac{8\pi G_0}{c^4}$$

is forced by the assumption that (1) and (2) agree in the limit of low velocities & weak gravitational fields: This is the

Correspondence Principle

History

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❑ Fundamental Question: Why do the planets move around the sun?

Newton: Ans: planets go around the sun because the sun "pulls" on the planets with an inverse square "force" (1687)

$$\vec{F} = M \vec{a}$$

$$\vec{F} = -G_0 \frac{M_s M_p}{r^2} \hat{r}$$

Einstein: Ans: planets go around the sun because the enormous mass of the sun (99% of mass) causes spacetime to be curved and planets move in "st lines" thru a "curved space" (1916)

$$G = K T$$

↗
Einstein
Curvature
Tensor

↖ stress Energy Tensor

"mass-energy & mass-energy
flux causes curvature"

• In 1687 (Principia) Newton showed that these laws explained the 3 Kepler laws:

- ① Planets move in ellipses around sun with sun at one focus
- ② equal areas in equal time
- ③ $\frac{L^3}{T^2} = \text{const}$ indept of planet.

Triumph 1846 John Adams & Urbain
Leverrier predict position of Neptune
from deviations in orbit of Uranus from
Newtonian prediction.

⌘ Problems with Newton Theory "Conceptually"

- (1) "force" as "pull" across empty space?
- (2) ∞ speed of propagation
- (3) Mysterious "Equivalence Principle"

Grav. Mass = Inertial Mass

Ex: let $F = M_E \vec{a}$ = force exerted by sun on earth

$$\vec{F} = \underset{\nearrow}{M_E} \vec{a} = -G_0 \frac{M_S \cancel{M_E}}{r^2} \overset{\nwarrow}{\vec{r}}$$

inertial mass
that resists acceleration

gravitational mass
appearing in 3rd law

Conclude: acceleration is independent of mass
of earth: i.e., a feather would
follow same trajectory as earth
if given same initial conditions:

⇒ Planetary motions appear to be more
a property of the space than of the
object moving thru the space (eg, due to "force"
on the object)

5 (d)

$$(4) \quad \vec{F} = -G_0 \frac{M_s M_e}{r^2} \vec{r} = M_e \vec{a}$$

is a coordinate dependent expression.

For a system of particles with positions

$$\underline{x}_1, \dots, \underline{x}_N \quad \underline{x}_i = (x_i^1, x_i^2, x_i^3) = \underline{x}_i(t)$$

$$m_n \frac{d^2 \underline{x}_n}{dt^2} = -G \sum_k \frac{m_n m_k (\underline{x}_n - \underline{x}_k)}{|\underline{x}_n - \underline{x}_k|^3}$$

Equations invariant under coord changes

$$\underline{x}' = R \underline{x} + \underline{v} t + \underline{d}$$

$$t' = t + \tau$$

$\Leftarrow$ Galileo Group

R arbitrary rotation (any orthogonal matrix
 $RR^T = Id$)

But only these:

$\Rightarrow$ Galilean Invariance.

Theory is not coordinate indept: eg: what determines these frames, and why do the stars appear to (coincidentally) be fixed in these frames?

Principle of Mach: Somehow the masses in the universe determine the inertial frames

☐ Einstein's Theory resolves these problems:

- Assumption: all properties of gravitational field are determined by a $sg(-1, 1, 1, 1)$ metric g on spacetime: in any coord system, x , g is a matrix field

$$g_{ij}(x) \equiv 4 \times 4 \text{ symmetric, nonsingular, bilinear form of } sg(-1, 1, 1, 1)$$

- free fall paths: geodesics of g
 - "aging time" for traveller along path $x(s)$
 $\equiv$ arclength of path as determined by g
 - Non-rotating coordinate frames determined by $\parallel$ -translation given by connection assoc. with g .
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Q: What is the equation that the gravitational field g_{ij} should satisfy?

$$\text{Einstein: } G = kT$$

$\nearrow$
 Einstein
 Curvature
 Tensor

$\nwarrow$
 Stress
 Energy
 Tensor

- Fundamental Result: 1859 Riemann defined the curvature ~~in the~~ associated with a Riemannian metric g at each point.
 $\Rightarrow$ Riemann Curvature Tensor.
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- In a coordinate system $x = (x^0, x^1, x^2, x^3)$ on spacetime, $x^0 = ct$ (5)

$$G_{ij} = K T_{ij} \quad ij = 0, \dots, 3$$

$$G_{ij} = R_{ij} - \frac{1}{2} R g_{ij}$$

G_{ij} and T_{ij} are 4×4 matrix (or tensor) fields defined on spacetime

$R_{ij} \equiv$ Ricci curvature tensor

$R \equiv$ scalar curvature

$g_{ij} \equiv$ spacetime metric

inverse of g_{ij}

$$R_{ij} = R^{\sigma}{}_{i\sigma j}, \quad R = R^{\sigma}{}_{\sigma} = g^{\sigma\tau} R_{\tau\sigma}$$

Einstein Summation Convention: sum repeated up-down indices from 0 to 3

$R^i{}_{jkl} \equiv$ Riemann Curvature Tensor (1851)

$$R^i{}_{jkl} = \underbrace{\Gamma^i_{jl,k} - \Gamma^i_{jk,l}}_{\text{"curl"}} + \underbrace{\Gamma^i_{\sigma k} \Gamma^{\sigma}_{jl} - \Gamma^i_{\sigma l} \Gamma^{\sigma}_{jk}}_{\text{"bracket"}}$$

$$,k \equiv \frac{\partial}{\partial x^k}$$

$$\Gamma_{ik}^i = \frac{1}{2} g^{i\sigma} \{ -g_{ik,\sigma} + g_{\sigma i,k} + g_{k\sigma,i} \} \quad (6)$$

Conclude: G_{ij} is a complicated function of the metric entries g_{ij} and their derivatives up to order 2.

• In empty space, $T_{ij} = 0 \Rightarrow$

Einstein: $G_{ij}[g] = 0 \quad \begin{cases} \text{2nd order PDE} \\ \text{in } g_{ij} \end{cases}$

Newton: $\Delta \Phi = 0 \quad \begin{cases} \text{2nd order PDE} \\ \text{in grav. potential} \end{cases}$

$\Phi \equiv$ gravitational potential:

$$\nabla \Phi = - \frac{\text{force}}{\text{mass}} \quad \text{at } P \in \mathbb{R}^3$$

- For a perfect fluid,

$$T_{ij} = (\rho c^2 + p) u_i u_j + p g_{ij}$$

$\rho \equiv \frac{\text{mass-energy}}{\text{3-vol}}$ in "frame of particle"

$$u^i \equiv \frac{dx^i}{d\tau} \equiv \text{"4-velocity of fluid particle"}$$

$p \equiv \text{pressure}$

Einstein: $G_{ij} = K T_{ij}$ $\left\{ \begin{array}{l} \text{Source involves} \\ \text{mass and momentum} \\ \text{densities and their fluxes} \end{array} \right.$

$$G_{00} = 1$$

Newton: $\Delta \Phi = 4\pi \rho$ $\left\{ \begin{array}{l} \rho \equiv \text{mass density} \\ \text{at fixed time.} \end{array} \right.$

- In limit of low velocities, the (0,0)-equn

$$G_{00} = K T_{00} \rightarrow \Delta \Phi = 4\pi \rho \quad \checkmark$$

- Problem with Newton - RHS $\rho(x,t)$ must evolve according to some equation which expresses cons. of mass & momentum - but Φ is determined instantaneously, indept of derivatives of ρ .

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• Indeed: In the classical case,

cons. of mass: $\boxed{\rho_t + \operatorname{div} \rho u = 0}$ 1 eqn

$u = \frac{dx}{dt} = 3\text{-velocity of fluid particle.}$

Newton's law: ($\Omega(t)$ moving with fluid)

$$\frac{d}{dt} \int_{\Omega(t)} \rho u^i d^3x = - \int_{\partial\Omega} p n^i d^2S - \int_{\Omega(t)} \nabla \Phi \rho d^3x$$

$\uparrow$
 outer normal

(grav. force on ρd^3x is $\nabla \Phi \rho d^3x$)

$$\Leftrightarrow \int_{\Omega(t)} (\rho u^i)_t + \operatorname{div}(\rho u^i) d^3x = \int_{\Omega(t)} -\frac{\partial}{\partial x^i} p - \nabla \Phi \rho d^3x$$

$$\Leftrightarrow \boxed{(\rho u^i)_t + \operatorname{div}(\rho u^i + p e^i) = -\rho \nabla \Phi^i} \quad 3 \text{ eqn's}$$

Newton

$\Rightarrow$ 4 eqn's in 4 unknowns ρ, u^i, Φ .

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- For Einstein, the conservation equations are built into $G_{ij} = \kappa T_{ij}$ automatically because G_{ij} is constructed to satisfy

$$\text{div } G = 0$$

 $\Rightarrow$

$$\text{div } T = 0$$

$\Leftrightarrow$ the relativistic Euler equations which express cons. of mass & energy.

$$T^{ij} = (\rho c^2 + p) u^i u^j + p g^{ij}$$

$$G = G[g^{ij}]$$

$g^{ij} \equiv$ gravitational metric tensor

Possible Topics :

I. Differential Geometry

- M^4 Manifold of events
- Tangent bundle / Cotangent bundle / Tensors / Differential Forms / Lie Derivatives
- Integration / maps / pullback / orientation
- Metric / $\parallel$ -translation / Connection / Covariant Derivatives / Curvature / Hodge - *

II. Einstein's Theory

- Special Relativity
 - Stress Energy Tensor / Perfect Fluid / E & M
 - $G = kT$ / Newtonian Limit / Red Shift / Fermi Transport / Coriolis force
 - Special Solutions / Schwarzschild Soln / Birkhoff Thm / Max' Symm Spaces / Robertson-Walker Metric / Interior Schwarzs. / De Sitter spaces / Oppenheimer Snyder
 - Stability
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- Initial value problem / well-posedness / equation count - gauge freedom / Bianchi identities
- Shock-Waves / Special Relativity / Te-Smo Result