

Example: Conservation Laws

Compressible Euler Equations:

$$\rho_t + \text{div}(\rho u) = 0$$

$$(\rho u^i)_t + \text{div}(\rho u^i u + p e^i) = 0$$

$$E_t + \text{div}(\rho(E+P)) = 0$$

In 1-d:

$$\begin{pmatrix} \rho \\ \rho u \\ E \end{pmatrix}_t + \begin{pmatrix} \rho u \\ \rho u^2 + p \\ (\rho(E+p)) \end{pmatrix}_x = 0$$

cons. mass
cons. mom
cons. energy

$$E = \rho e + \frac{1}{2} \rho u^2$$

$\rho \equiv$ density
 $u \equiv$ velocity

$e =$ specific internal energy $\equiv \frac{\text{energy}}{\text{mass}}$

$E \equiv \frac{\text{energy}}{\text{vol}}$

In general - an equation of state required to get 3 equations in 3 unknowns:

Eg. $e = e(v, s)$ $v = \frac{1}{\rho}$

$p = e_v(v, s)$ $s = \text{specific entropy}$

Unknowns: (ρ, u, s)

OR

$e = e(p, \rho)$

Unknowns (ρ, u, p)

General $n \times n$ system of conservation laws:

$$U_t + F(U)_x = 0$$

F - nonlinear Euler: $U = \begin{pmatrix} \rho \\ \rho u \\ E \end{pmatrix}$

$U \equiv$ "the conserved quantities"

$$\square \quad \underbrace{U}_t + \underbrace{A(U)} \cdot \underbrace{U}_x = 0 \quad A(U) \equiv \frac{\partial F^i}{\partial U^j} \quad (c) \quad (36)$$

$(n \times 1) \quad n \times n \quad n \times 1$

Under change of dependent variables:

$$V = \phi(U)$$

eg. , $\begin{pmatrix} p \\ u \\ s \end{pmatrix} = \phi \left(\begin{pmatrix} p \\ s \\ u \end{pmatrix} \right)$

Then

$$V_t^\alpha = \frac{\partial V^\alpha}{\partial U^j} U_t^j$$

$$V_x^\alpha = \frac{\partial V^\alpha}{\partial U^j} U_x^j$$

Conservation Law:

$$U_t^i + \frac{\partial F^i}{\partial U^j} U_x^j = 0$$

mult by $\frac{\partial V^\alpha}{\partial U^i}$:

$$\underbrace{\frac{\partial V^\alpha}{\partial U^i} U_t^i}_{V_t^\alpha} + \underbrace{\frac{\partial V^\alpha}{\partial U^i} \frac{\partial F^i}{\partial U^j} \frac{\partial U^\sigma}{\partial V^\tau} \frac{\partial V^\tau}{\partial U^j}}_{J^{-1} A J} U_x^j = 0$$

δ^σ_j

V_x^τ

$$V_t^\alpha + \underbrace{\frac{\partial V^\alpha}{\partial U^i} \frac{\partial F^i}{\partial U^\sigma} \frac{\partial U^\sigma}{\partial V^\tau}}_{\bar{A}^\alpha_\tau} V_x^\tau = 0$$

Conclude: under change of dependent

variable, $U_t + A(U)U_x$

transforms like:

U_t^i, U_x^i

transform like a vectors:

$\bar{V}_t^\alpha = \frac{\partial V^\alpha}{\partial U^i} U_t^i$

$A \equiv A^i_j$ transforms like a (1) - tensor:

$\bar{A}^\alpha_\beta = A^i_j \frac{\partial U^j}{\partial V^\beta} \frac{\partial V^\alpha}{\partial U^i}$

~~Conclude: For any choice of state variables~~

Conclude: The eigenvalues of A can be computed from the equation $V_t + \bar{A} \cdot V_x$ for any state variables V .

- The ^{right} eigenvectors of A transform like \bar{R} vectors: (39)

$$A R = \lambda R$$

$$\Rightarrow \bar{A} \bar{R} = \lambda \bar{R}$$

where

$$\bar{R}^\alpha = R^i \frac{\partial \bar{V}^\alpha}{\partial U^i}$$

In this example, U & V represent coordinates on the "manifold" of dependent variables

~~• The left eigenvectors of A transform like components of co-vectors:~~

$$L A = \lambda A$$

• The left eigenvectors of A transform like components of co-vectors. (39D)

$$LA = \lambda L \quad (L_i A^i_j = \lambda L_j)$$

$$\Rightarrow \bar{L} \bar{A} = \lambda \bar{L} \quad (\bar{L}_\alpha A^\alpha_\beta = \lambda L_\beta)$$

where

$$\bar{L}_\alpha = \frac{\partial U^i}{\partial V^\alpha} L_i$$

Moreover: $LA = \lambda L$ has a soln iff

$$\det |\lambda I - A| = 0$$

\Rightarrow iff λ is a root of char polyn \Rightarrow right/left eigenvalues are the same.

~~Theorem~~:

Lemma: if $LA = \lambda_1 L$, $AR = \lambda_2 R$

$\lambda_1 \neq \lambda_2$, then

$$L_\sigma R^\sigma = 0$$

P.f.

$$L_\sigma A^\sigma R^\tau = \lambda_1 L_\sigma R^\tau$$

$$= \lambda_2 L_\sigma R^\tau$$

$\Rightarrow L_\sigma R^\tau = 0$ unless $\lambda_1 = \lambda_2$ ✓

Defn: system is strictly hyperbolic if A has n real & distinct e-values at each state value \bar{U} .

oft. quoted: " L^k is orthogonal to all the right eigenvectors R_e , $k \neq e$ "

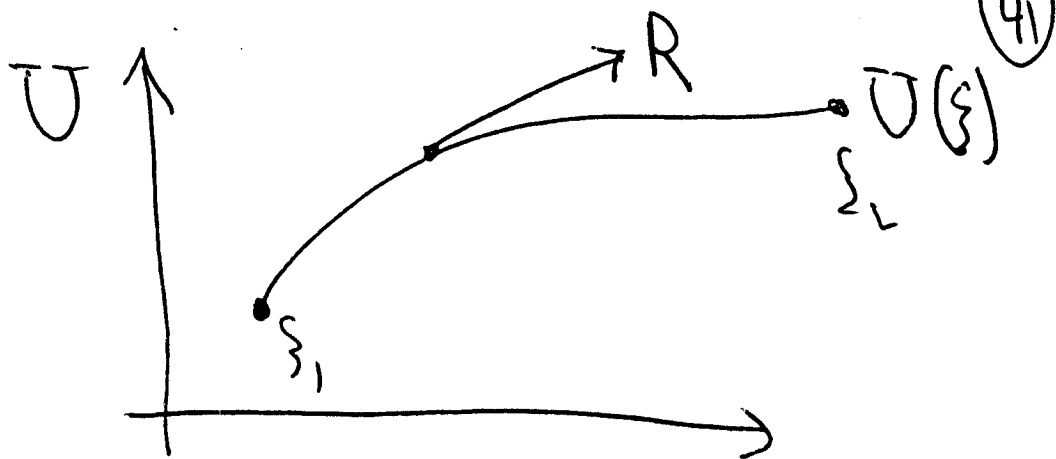
The eigenvalues & eigenvectors associated with system (c) are important as follows: (40)

Construction: Let $R^i(\bar{U})$, $\lambda(\bar{U})$ be the eigenvector - eigenvalue field associated with matrix $A(\bar{U}) \equiv A^i_j(\bar{U})$. Let $U^i(x, t)$ be any smooth function ~~that~~ ~~satisfies the condition~~ constructed as

follows: Let $\bar{U}(s)$ denote an integral curve of the vector field R^i defined on \bar{U} -space; i.e.,

$$\frac{d}{ds} U^i(s) = R^i(\bar{U}(s)).$$

Picture:



Let $\lambda(U(\xi))$ be the eigenvalue along the integral curve.

Defn: λ is genuinely nonlinear if $\nabla \lambda \cdot R_i \neq 0$
I.e., λ is monotone increasing or decreasing along $U(\xi)$.

Let $\bar{U}(x,t)$ be any function satisfying:

$$\bar{U}(x,t) \equiv U(\xi)$$

$\xi = \xi(x,t)$ a smooth fn

along lines in the xt -plane of speed $\frac{dx}{dt} = \lambda(U(\xi))$

Theorem: $\bar{U}(x,t)$ is a solution of (c). Such a soln is called a λ -simple wave.

Assume $\nabla \lambda \cdot R_i \neq 0$ along integral curve R ,
so that we can take $U \equiv U(\lambda)$ as
parameterizing R . (42)

Theorem: Let $\lambda(x, t)$ denote any smooth
solution of

$$\lambda_t + \lambda \lambda_x = 0.$$

Then

$U(\lambda(x, t))$
is a smooth solution of

$$U_t + A(U)_x = 0.$$

Cor: Since λ is a scalar, this construction
is valid in any set of variables V ,
so long as $V'(\lambda) = R$

Proof:

$$U_t = U' \lambda_t = R \lambda_t$$

$$U_x = U' \lambda_x = R \lambda_x$$

Thus

$$U_t + A(U) U_x = R \lambda_t + A R \lambda_x$$

$$= R \lambda_t + R \lambda \lambda_x = R (\lambda_t + \lambda \lambda_x) = 0$$

Comment: $U_t + F(U)_x = 0$ is in conservation form. Not all systems of form

$$\bar{V}_t + A(\bar{V}) \bar{V}_x = 0$$

can be written in conservation form —
In fact — it is the conservation form that determines the shock-waves.

Note: L_i is a left eigenvector of