

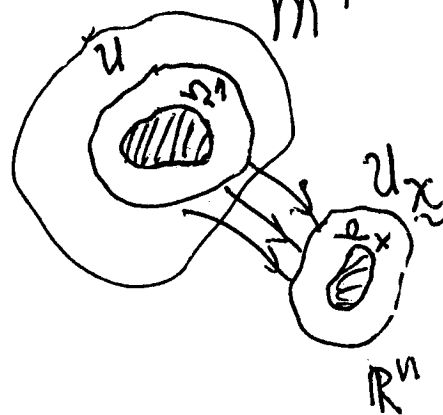
Differential Forms / Integration (Local meaning)

- General Manifold M^n , no metric (Warner)
- $\Omega^n \subset M^n$ open, n -dimensional

$$f : M^n \rightarrow \mathbb{R}$$

Assume: Ω^n covered by single coord chart x

$$\begin{aligned} x : U_x &\rightarrow \mathbb{R}^n, \quad x(U) = U_x \\ \Omega^n &\subset U_x, \quad x(\Omega^n) = \Omega_x \end{aligned}$$



- Problem: How to give

$$\int_{\Omega_x} f$$

a coordinate independent meaning.

- Note: Under coordinate transformation $x \rightarrow y$, integral transforms by

$$\int_{\Omega_x} f dx^1 \cdots dx^n = \int_{\Omega_y} f \left| \frac{\partial x}{\partial y} \right| dy^1 \cdots dy^n, \quad (*)$$

where

$$\left| \frac{\partial x}{\partial y} \right| = \det \left(\frac{\partial x^i}{\partial y^j} \right).$$

E.g., if we interchange x^i, x^j , then integral changes by factor (-1) . $\Rightarrow$ (*) has a "double valued" meaning - $\pm$ for opposite oriented "loop travel".

- Idea: define $dx^1 \cdots dx^n = dx^1 \wedge \cdots \wedge dx^n$ in such a way that under a change of coordinates,

$$dx^1 \wedge \cdots \wedge dx^n = \left| \frac{\partial x}{\partial y} \right| dy^1 \wedge \cdots \wedge dy^n. \quad (*)$$

Then, if we let

$$\alpha = dx^1 \wedge \cdots \wedge dx^n,$$

then

$$\int_{\Omega^n} f \alpha$$

is a coord. indep't expression, (modulo a $\pm$ sign to be discussed later)

- To achieve (*), it suffices to make $dx^1 \cdots dx^n$ a completely anti-symmetric n -tensor:

Defn: $T_{i_1 \cdots i_k} dx^{i_1} \otimes \cdots \otimes dx^{i_k}$ is anti-symmetric $\binom{0}{k}$ -tensor, $1 \leq k \leq n$, if

$$T_{i_{\pi(1)} \cdots i_{\pi(k)}} = (-1)^\pi T_{i_1 \cdots i_k}$$

where π is a permutation of indices $i_1 \cdots i_n$

$(-1)^\pi = -1$ if π is odd - it takes odd # of interchanges to take $i_1 \cdots i_k \rightarrow \pi(i_1) \cdots \pi(i_k)$

$(-1)^\pi = 1$ if π is even.

Immediate: $T_{i_1 \cdots i_k} = 0$ if $i_p = i_q$, so it suffices to sum over seq's of distinct indices $i_1 \cdots i_k$ in (A)

Lemma: Anti-symmetry is a coord. indep property of a $\binom{0}{k}$ tensor. (HW FIP)

- Explicit examples of $\binom{0}{k}$ -tensors:

To construct one, we need only define an anti-symmetric ^{linear} operator on k vectors.

$k=n$

Given: $X_1, \dots, X_n$ in TM^n

In x-words: $X_{\underline{x}} = X_{\underline{x}}^i \frac{\partial}{\partial x^i}$

$\Rightarrow$ $n \times n$ coordinate matrix $X_{\underline{x}}^i$

$$X_{\underline{x}}^i = \begin{bmatrix} | & | & & | \\ X_1^i & X_2^i & \dots & X_n^i \\ | & | & & | \end{bmatrix}_{\underline{x}}^i$$

Defn: $dx^1 \wedge \dots \wedge dx^n (X_1, \dots, X_n) = \det(X_{\underline{x}}^i)$

Clearly: linear, anti-symmetric

Problem: how does it transform under coordinate transformation $\underline{x} \rightarrow \underline{y}$?

• Soln: $\det(X^i_x) \equiv \sum_{\pi} (-1)^{\pi} X_1^{\pi(1)} \cdots X_n^{\pi(n)} \quad (5)$

"sum over all permutations π of $(1, \dots, n)$."

$$= \sum_{\pi} (-1)^{\pi} dx^{\pi(1)} \otimes \cdots \otimes dx^{\pi(n)} (x_1, \dots, x_n)$$

$$= T_{i_1 \dots i_n} dx^{i_1} \otimes \cdots \otimes dx^{i_n} (x_1, \dots, x_n)$$

where $T_{i_1 \dots i_n} = (-1)^{\pi} \quad \pi: (1, \dots, n) \rightarrow (i_1, \dots, i_n)$

Now let $\tilde{x} \rightarrow \tilde{y}$ be a coord transformation:

$$T = T_{i_1 \dots i_n} dx^{i_1} \otimes \cdots \otimes dx^{i_n}$$

$$\Rightarrow T = \overline{T}_{\alpha_1 \dots \alpha_n} dy^{\alpha_1} \otimes \cdots \otimes dy^{\alpha_n}$$

where

$$\overline{T}_{\alpha_1 \dots \alpha_n} = T_{i_1 \dots i_n} \frac{\partial x^{i_1}}{\partial y^{\alpha_1}} \cdots \frac{\partial x^{i_n}}{\partial y^{\alpha_n}}$$

(6)

But:

$$T_{i_1 \dots i_n} \frac{\partial x^{i_1}}{\partial y^{\alpha_1}} \dots \frac{\partial x^{i_n}}{\partial y^{\alpha_n}} = \sum_{\pi} (-1)^{\pi} \frac{\partial x^{\pi(1)}}{\partial y^{\alpha_1}} \dots \frac{\partial x^{\pi(n)}}{\partial y^{\alpha_n}}$$

$$= (-1)^{\bar{\pi}} \det \frac{\partial x}{\partial y} \quad \bar{\pi} : (1, \dots, n) \rightarrow (\alpha_1, \dots, \alpha_n)$$

$\therefore$ in y -coordinates,

$$T = T_{\alpha_1 \dots \alpha_n} dy^{\alpha_1} \otimes \dots \otimes dy^{\alpha_n} = \sum_{\bar{\pi}} (-1)^{\bar{\pi}} \det \frac{\partial x}{\partial y} dy^{\alpha_1} \otimes \dots \otimes dy^{\alpha_n}$$

$$= \det \frac{\partial x}{\partial y} \sum_{\bar{\pi}} (-1)^{\bar{\pi}} dy^{\alpha_1} \otimes \dots \otimes dy^{\alpha_n}$$

$$\equiv \det \frac{\partial x}{\partial y} dy^1 \wedge \dots \wedge dy^n.$$

Conclude:

$$dx^1 \wedge \dots \wedge dx^n = \det \frac{\partial x}{\partial y} dy^1 \wedge \dots \wedge dy^n$$

has the desired transformation properties.

- Invariant defn of integral: we call (7)
 $\alpha = a(x) dx^1 \wedge \dots \wedge dx^n$, (a a scalar)
 a volume form, or n -form. Then (depending on coords)

$$\int_{\Omega} f \alpha = \int_{\Omega_x} f(x) a(x) dx^1 \wedge \dots \wedge dx^n$$

by which we mean

$$= \int_{\Omega_x} f(x) a(x) dx^1 \wedge \dots \wedge dx^n.$$

Then in y -coords,

$$\alpha = a(x(y)) \det \frac{\partial x}{\partial y} dy^1 \wedge \dots \wedge dy^n$$

$$\Rightarrow \int_{\Omega} f \alpha = \int_{\Omega_y} f(y) a(x(y)) \det \frac{\partial x}{\partial y} dy^1 \wedge \dots \wedge dy^n$$

$$= \int_{\Omega_y} f(y) a(x(y)) \det \frac{\partial x}{\partial y} dy^1 \wedge \dots \wedge dy^n$$

gives the invariant defn of $\int_{\Omega} f \alpha$.

(8A)

• Another Way: Define

$$\int_{\Omega} f \alpha \equiv \int_{\Omega_x} f(x) \underbrace{\alpha\left(\frac{\partial}{\partial x^1}, \dots, \frac{\partial}{\partial x^n}\right)}_{a(x) \cdot 1} dx^1 \cdots dx^n$$

Then in y -coordinates

$$\int_{\Omega} f \alpha \equiv \int_{\Omega_y} f(x) \underbrace{\alpha\left(\frac{\partial}{\partial y^1}, \dots, \frac{\partial}{\partial y^n}\right)}_{a(x) \det \frac{\partial x}{\partial y} \cdot 1} dy^1 \cdots dy^n$$

In this picture, " $\alpha(X_1, \dots, X_n)$ eats the n -vectors $X_1, \dots, X_n$ and evaluates the α -measure of the n -volume of the n -opiped spanned by $X_1, \dots, X_n$."

Really: $\alpha_p(X_1, \dots, X_n) \equiv \frac{\alpha\text{-Volume of } X_1, \dots, X_n}{\text{coord volume}}$ in coord system with $\frac{\partial}{\partial x^i} = X_i$ at P .

- ORIENTATION: Let σ be an n -form defined on M . Then in coordinates,

$$\sigma = a(\underline{x}) dx^1 \wedge \dots \wedge dx^n.$$

Defn: an orientation on M is a non-zero n -form defined on M .

I.e., given any coord system $\underline{x}$, we say $\underline{x}$ is pos oriented wrt σ if $a(\underline{x}) > 0$. Since in y -words,

$$\sigma = a(\underline{x}(\underline{y})) \det \frac{\partial \underline{x}}{\partial \underline{y}} dy^1 \wedge \dots \wedge dy^n$$

$$\sigma = \bar{a}(\underline{y}) dy^1 \wedge \dots \wedge dy^n, \quad \bar{a} = a \det \frac{\partial \underline{x}}{\partial \underline{y}},$$

it follows that $\underline{y}$ is pos oriented wrt σ iff $\det \frac{\partial \underline{x}}{\partial \underline{y}} > 0$, as we want.

(8c)

Defn: a manifold M is orientable if $\exists$ an orientation for M .

Homework: show that the sphere is orientable but the torus is not.

Recall: a metric on M is a ^{symmetric} nonsingular $(0,2)$ -tensor ~~def~~ which defines a quadratic form on each $T_p M$: i.e.

$$g \equiv g_{ij} dx^i \otimes dx^j$$

$$(g^{ij}) = (g_{ij})^{-1}$$

$$X = a^i \frac{\partial}{\partial x^i}, Y = b^j \frac{\partial}{\partial x^j}$$

$$\langle X_p, Y_p \rangle = g(X, Y)_p = g_{ij} a^i b^j$$

• Canonical map $T_p M \rightarrow T_p^* M$ (raising/lowering)

$$T_p M \ni X \leftrightarrow \omega \in T_p^* M$$

$$a^i \frac{\partial}{\partial x^i} \leftrightarrow a_i dx^i$$

$$a_i = g_{i\sigma} a^\sigma$$

$$g^{i\sigma} a_\sigma = a^i$$

$$\Rightarrow \langle X, Y \rangle = \omega(Y)$$

$$\forall Y \in T_p M$$

• This extends to a mapping of contravariant indices to covariant ones:

$$T_{ij} dx^i \otimes dx^j \leftrightarrow T^{ij} \frac{\partial}{\partial x^i} \otimes \frac{\partial}{\partial x^j} \quad \parallel \quad T^{ij} = T_{\sigma\tau} g^{\sigma i} g^{\tau j}$$

etc

- The canonical mapping extends to raising & lowering of any indices on a tensor:

$$T^i_j \frac{\partial}{\partial x^i} \otimes dx^j \leftrightarrow T_{ij} dx^i \otimes dx^j$$

$$T_{ij} = T^\sigma_j g_{\sigma i}$$

- Thm: The raising / lowering of two covariant / contravariant indices preserves symmetry / antisymmetry wrt the two indices.

Eg. $T_{ij}^{---} = T^{\sigma\tau}{}^{---} g_{i\sigma} g_{j\tau}$

then $T_{ij}^{---} = T_{ji}^{---}$ iff $T^{\sigma\tau}{}^{---} = T^{\tau\sigma}{}^{---}$

and $T_{ij}^{---} = -T_{ji}^{---}$ iff $T^{\sigma\tau}{}^{---} = -T^{\tau\sigma}{}^{---}$

FIP

(90)

Defn: $\Lambda^k \equiv$ set of all completely antisymmetric
 $\binom{0}{k}$ -tensors $\equiv$ differential forms
of order k

$$\equiv \{ T_{i_1 \dots i_k} dx^{i_1} \otimes \dots \otimes dx^{i_k} :$$

$$T_{i_1 \dots i \dots j \dots i_k} = -T_{i_1 \dots j \dots i \dots i_k} \}$$

$\Lambda_n \equiv$ set of all completely antisymmetric -
 $\binom{k}{0}$ -tensors $\equiv$ co-differential forms
of order k

$$\equiv \{ T^{i_1 \dots i_k} \frac{\partial}{\partial x^{i_1}} \otimes \dots \otimes \frac{\partial}{\partial x^{i_k}} :$$

$$T^{i_1 \dots i \dots j \dots i_k} = -T^{i_1 \dots j \dots i \dots i_k} \}$$

(9d)

• Conclude: The metric g induces a canonical mapping from $\Lambda^k \rightarrow \Lambda_k$ by

$$T_{i_1 \dots i_k} dx^{i_1} \otimes \dots \otimes dx^{i_k} \mapsto T^{i_1 \dots i_k} \frac{\partial}{\partial x^{i_1}} \otimes \dots \otimes \frac{\partial}{\partial x^{i_k}}$$

$$T_{i_1 \dots i_k} = T^{\sigma_1 \dots \sigma_k} g_{\sigma_1 i_1} \dots g_{\sigma_k i_k}.$$

This in turn induces an inner product on Λ^k and Λ_k = e.g., $S, T \in \Lambda_k$,

$$S = S^{i_1 \dots i_k} \frac{\partial}{\partial x^{i_1}} \otimes \dots \otimes \frac{\partial}{\partial x^{i_k}}$$

$$T = T^{i_1 \dots i_k} \frac{\partial}{\partial x^{i_1}} \otimes \dots \otimes \frac{\partial}{\partial x^{i_k}}$$

$$\langle S, T \rangle = S_{i_1 \dots i_k} T^{i_1 \dots i_k}$$

• It turns out - this determines k -Vol induced by g (as we shall see)

start

9e

□ $\exists$ a natural volume form associated with a metric g defined on M :

- $\forall p \in M \exists$ coord system defined in a nbhd of p in which

$$g_{ij} = d_{ij} = \text{diag}\{\sigma_1, \dots, \sigma_n\}$$

$$\sigma_1 = \dots = \sigma_r = -1, \sigma_{r+1} = \dots = \sigma_n = 1$$

$$s \equiv \text{sgn}(g) \equiv \sum_i \sigma_i, \quad s = -r + (n-r) = n-2r$$

Assume g has $\text{sgn}(g) = s$ at every p

eg. Pos. definite case: $\text{sgn}(g) = n$, $d_{ij} = \delta_{ij}$

Lorentzian case $n=4$, $s=2$, $d_{ij} = \eta_{ij}$

- We derive a volume form that is naturally associated with g :

Fix $p \in M$, and choose a coord. system $\underline{x}$ that is orthonormal at p , i.e.,

$$g_{ij}(p) = d_{ij} \equiv \text{diag} \{ \sigma_1, \dots, \sigma_n \}.$$

Note first that

$$\det(d_{ij}) = (-1)^r,$$

and let

$$g = g(\underline{x}) \equiv \det g_{ij}(\underline{x}).$$

Recall that $\det T_j^i$ of a $\binom{1}{1}$ -tensor is a coordinate independent number at each p , but not so for $\binom{0}{2}$ -tensors; thus $g = g(\underline{x})$ is a coordinate dependent number at each pt p .

In $\underline{x}$ -coord's,

$$g = \det g_{ij} = \det d_{ij} = (-1)^r \text{ at } p.$$

- Since $\underline{x}$ is o.n. at p , it is natural to define the volume form so that $(\frac{\partial}{\partial x^1}, \dots, \frac{\partial}{\partial x^n})_p$ span the unit volume, i.e. (11)

$$\eta = \pm dx^1 \wedge \dots \wedge dx^n.$$

(We can choose $\pm$ here to set a positively oriented frame). To see that this is consistent, let $\underline{y}$ be any other coord frame that is o.n. at p , i.e.

$$\bar{g}_{\alpha\beta}(\underline{x}_0) = \delta_{\alpha\beta} = \text{diag}\{\sigma_1, \dots, \sigma_n\}.$$

Then

$$\bar{g}_{\alpha\beta} = g_{ij} \frac{\partial x^i}{\partial y^\alpha} \frac{\partial x^j}{\partial y^\beta}$$

$$\det \bar{g}_{\alpha\beta} = \det g_{ij} \left(\det \frac{\partial x^i}{\partial y^\alpha} \right)^2$$

$$\Rightarrow \det \frac{\partial x^i}{\partial y^\alpha} = \pm 1 \quad \left(\begin{array}{l} \pm \text{ depending on} \\ \text{whether } \underline{x} \text{ is pos} \\ \text{oriented wrt } \underline{y} \end{array} \right)$$

Thus in another ON frame $\underline{y}$ that is positively oriented wrt $\underline{x}$ we have (12)

$$\eta = + \det\left(\frac{\partial y^k}{\partial x^i}\right) dy^1 \dots dy^n$$
$$= + dy^1 \dots dy^n.$$

Conclude: The volume form η agrees with the Euclidean coordinate volume in each positively oriented ON-frame, at P .

• Theorem: The volume form η which has value

$$\eta = dx^1 \dots dx^n$$

in each positively oriented frame $\underline{x}$ at P , $g_{ij}(\underline{x}(P)) = \delta_{ij}$, takes the value

$$\eta = \sqrt{|g|} dy^1 \dots dy^n$$

in an arbitrary frame positively oriented with $\underline{x}$.

Proof: Let $\underline{y}$ be an arbitrary coordinate system pos. oriented wrt $\underline{x}$. Then in ~~the~~ coordinates,

$$\eta = dx^1 \dots dx^n = \left(\det \frac{\partial \underline{x}}{\partial \underline{y}} \right) dy^1 \dots dy^n.$$

But also

$$\det d_{ij} = \det g_{ij} = \det \left(\bar{g}_{AB} \frac{\partial y^A}{\partial x^i} \frac{\partial y^B}{\partial x^j} \right)$$

$$\Rightarrow \det(d_{ij}) = \det \bar{g}_{AB} \left| \det \frac{\partial \underline{y}}{\partial \underline{x}} \right|^2$$

In particular, $\left| \det \frac{\partial \underline{y}}{\partial \underline{x}} \right|^2 > 0 \Rightarrow$

$$\text{sign} \{ \det d_{ij} \} = \text{sign} \{ \det \bar{g}_{AB} \} = (-1)^r.$$

Also

$$\boxed{\det \frac{\partial \underline{y}}{\partial \underline{x}} = \sqrt{\frac{(-1)^r}{\det \bar{g}_{AB}}} \Leftrightarrow \det \frac{\partial \underline{x}}{\partial \underline{y}} = \sqrt{(-1)^r \det \bar{g}_{AB}}}$$

thus

$$\eta = \sqrt{(-1)^r \det \tilde{g}_{AB}} \, dy^1 \wedge \dots \wedge dy^n$$

$$\boxed{\eta = \sqrt{|g|} \, dy^1 \wedge \dots \wedge dy^n}$$

gives a formula for the volume form
in any coord frame y that is positively
oriented wrt $\underline{x}$

Notation : $\varepsilon \equiv dx^1 \cdots dx^n$ is a coordinate dependent $\binom{0}{n}$ -tensor called the Levi-Civita tensor. It evaluates coordinate n -volume. Thus

$$\eta = \pm \sqrt{g} \varepsilon$$

where we choose $+$ if $\underline{x}$ is positively oriented with the canonical coordinate system (at each p) in which

$$g_{ij} = d_{ij} = \text{diag}\{\sigma_1, \dots, \sigma_n\}$$

thus - in ON frames, the metric-volume form agrees with the Euclidean coord. volume form ε .

▣ k -dimensional volume induced by g :

→ Gen theory w/o metric:

Defn: Let $\Lambda^k(p)$, $\Lambda_n(p)$ denote the space of completely antisymmetric $\binom{0}{k}$ -tensors, $\binom{k}{0}$ -tensors, respectively, defined at $p \in M$. Let Λ^k , Λ_n denote the space of completely antisymmetric $\binom{0}{k}$ -, $\binom{k}{0}$ -tensor fields defined on M , resp.

Lemma ①: $\Lambda^k(p)$ & $\Lambda_n(p)$ are vector spaces at each $p \in M$.

We now construct a natural basis for $\Lambda^k(p)$, $\Lambda_n(p)$.

• Note: The construction of the spaces Λ^k , Λ_n applies to an arbitrary manifold w/o a metric g .

• Let $\underline{x}$ be a coord system defined in a nbhd of $p \in M$. Define

$$dx^{i_1} \wedge \dots \wedge dx^{i_n} \in \wedge^k$$

(*)

at each $p \in M$ by:

$$dx^{i_1} \wedge \dots \wedge dx^{i_n} (X_1, \dots, X_n)$$

$$= \det_{i_1 \dots i_n} \begin{bmatrix} \frac{1}{X_1} & \dots & \frac{1}{X_n} \\ \vdots & & \vdots \end{bmatrix}$$

where $\begin{bmatrix} \frac{1}{X_1} & \dots & \frac{1}{X_n} \\ \vdots & & \vdots \end{bmatrix}$ denotes the $n \times n$ matrix

obtained by putting the $\underline{x}$ -coord of X_i in the i th col., and $\det_{i_1 \dots i_n}$ denotes the

det taken in the $i_1 \dots i_n$ rows, ~~re-ordered by~~
in the order $i_1 \dots i_n$.

• Note (*) is a coord. dependent expression

- Observe: $dx^{i_1} \wedge \dots \wedge dx^{i_k}$ defines a multilinear, antisymmetric function on k -vectors $X_1, \dots, X_k$, and thus defines an element of $\wedge^k$. Moreover, by antisymmetry,

$$dx^{i_1} \wedge \dots \wedge dx^{i_k}(X_1, \dots, X_k) = 0$$

if $i_\ell = i_m$ or $X_\ell = X_m$. FIP

- We obtain the tensor components of $dx^{i_1} \wedge \dots \wedge dx^{i_k}$ in terms of the basis $\{dx^{i_1} \otimes \dots \otimes dx^{i_k}\}$ for $\mathcal{Y}_k^0 \supset \wedge^k$ at each $p \in M$:

$$\begin{aligned} dx^{i_1} \wedge \dots \wedge dx^{i_k} \left[\frac{\partial}{\partial x^{j_1}}, \dots, \frac{\partial}{\partial x^{j_k}} \right] &= \sum_{\pi} (-1)^\pi X_1^{i_{\pi(1)}} \dots X_k^{i_{\pi(k)}} \\ &= \sum_{\pi} (-1)^\pi dx^{i_{\pi(1)}} \otimes \dots \otimes dx^{i_{\pi(k)}} (X_1, \dots, X_k) \\ &= T_{j_1, \dots, j_k}^{i_1, \dots, i_k} dx^{j_1} \otimes \dots \otimes dx^{j_k} \end{aligned}$$

where

(19)

$$T_{j_1 \dots j_n} = \begin{cases} 1 & \text{if } (j_1 \dots j_n) = (i_1 \dots i_n) \\ (-1)^\pi & \text{if } (j_1 \dots j_n) = (i_{\pi(1)} \dots i_{\pi(n)}) \\ 0 & \text{otherwise.} \end{cases}$$

Note: By defn of $dx^{i_1} \wedge \dots \wedge dx^{i_n}$ as taking a $k \times k$ det., we have

$$dx^{i_1} \wedge \dots \wedge dx^{i_k} \wedge dx^{j_1} \wedge \dots \wedge dx^{j_k} = - dx^{i_1} \wedge \dots \wedge dx^{j_1} \wedge \dots \wedge dx^{i_k}$$

Theorem: $\{dx^{i_1} \wedge \dots \wedge dx^{i_n}, 0 \leq i_1 < i_2 < \dots < i_n \leq n\}$
forms a basis for $\mathcal{L}^k(p)$ at each p . Thus,
 $\dim \mathcal{L}^k = \binom{n}{k}$, and $T \in \mathcal{L}^k \Rightarrow$

$$T = T_{|i_1 \dots i_n|} dx^{i_1} \wedge \dots \wedge dx^{i_n}$$

where $|i_1 \dots i_n|$ means sum over all increasing
sequences of indices

- Conclude: every k -form can be written as:

$$T_{i_1 \dots i_k} dx^{i_1} \otimes \dots \otimes dx^{i_k}, \quad T_{i_{\pi(1)} \dots i_{\pi(k)}} = (-1)^\pi T_{i_1 \dots i_k}$$

$$T_{|i_1 \dots i_k|} dx^{i_1} \wedge \dots \wedge dx^{i_k}$$

$$\frac{1}{k!} T_{i_1 \dots i_k} dx^{i_1} \wedge \dots \wedge dx^{i_k}$$

(★)

I.e. $\Lambda^k \ni T = T_{i_1 \dots i_k} dx^{i_1} \otimes \dots \otimes dx^{i_k}$

$$= \sum_{|i_1 \dots i_k|} \sum_{\pi} (-1)^\pi T_{i_1 \dots i_k} dx^{i_{\pi(1)}} \otimes \dots \otimes dx^{i_{\pi(k)}}$$

$$= \sum_{|i_1 \dots i_k|} T_{i_1 \dots i_k} \sum_{\pi} (-1)^\pi dx^{i_{\pi(1)}} \otimes \dots \otimes dx^{i_{\pi(k)}}$$

$$= T_{|i_1 \dots i_k|} dx^{i_1} \wedge \dots \wedge dx^{i_k}$$

$$= \frac{1}{k!} T_{i_1 \dots i_k} dx^{i_1} \wedge \dots \wedge dx^{i_k} \quad \checkmark$$

• Transformation Law:

$$T = T_{|i_1 \dots i_n|} dx^{i_1} \dots dx^{i_n} \quad \underline{x} \rightarrow \underline{y}$$

$$\Rightarrow T = T_{|\alpha_1 \dots \alpha_n|} dy^{\alpha_1} \dots dy^{\alpha_n}$$

where

$$T_{\alpha_1 \dots \alpha_n} = T_{|i_1 \dots i_n|} \det_{\alpha_1 \dots \alpha_n}^{|i_1 \dots i_n|} \frac{\partial x}{\partial y}$$

P.T.

$$T = T_{|i_1 \dots i_n|} dx^{i_1} \otimes \dots \otimes dx^{i_n}$$

$$= T_{|i_1 \dots i_n|} \frac{\partial x^{i_1}}{\partial y^{\alpha_1}} \dots \frac{\partial x^{i_n}}{\partial y^{\alpha_n}} dy^{\alpha_1} \otimes \dots \otimes dy^{\alpha_n}$$

$$= \sum_{|i_1 \dots i_n|} T_{|i_1 \dots i_n|} \sum_{\pi} (-1)^{\pi} \frac{\partial x^{i_{\pi(1)}}}{\partial y^{\alpha_1}} \dots \frac{\partial x^{i_{\pi(n)}}}{\partial y^{\alpha_n}}$$

$$dy^{\alpha_1} \otimes \dots \otimes dy^{\alpha_n}$$

$$= \underbrace{T_{|i_1 \dots i_n|} \det_{\alpha_1 \dots \alpha_n}^{|i_1 \dots i_n|} \frac{\partial x}{\partial y}}_{\text{"}} dy^{\alpha_1} \otimes \dots \otimes dy^{\alpha_n}$$

$$T_{\alpha_1 \dots \alpha_n}$$

Proof of Theorem:

We first show that

$$\{dx^\lambda \equiv dx^{\lambda_1} \cdots dx^{\lambda_n}, \lambda_1 < \cdots < \lambda_n < n\}$$

form an independent set of vectors in L^n .

I.e., if

$$\alpha_\lambda dx^\lambda = 0, \quad (\text{sum over } \lambda \text{ implied})$$

then

$$\alpha_\lambda dx^\lambda \left[\frac{\partial}{\partial x^{\lambda_1}} \cdots \frac{\partial}{\partial x^{\lambda_n}} \right] = 0$$

$\forall \{x_1, \dots, x_n\}$. Choose $X_1, \dots, X_n = \frac{\partial}{\partial x^{\lambda_1}}, \dots, \frac{\partial}{\partial x^{\lambda_n}}$.

Then

$$dx^\mu \left(\frac{\partial}{\partial x^{\lambda_1}} \cdots \frac{\partial}{\partial x^{\lambda_n}} \right) = 0 \quad \mu \neq \lambda,$$

so

$$dx^\lambda \left(\frac{\partial}{\partial x^{\lambda_1}} \cdots \frac{\partial}{\partial x^{\lambda_n}} \right) = 1.$$

Thus

$$0 = \alpha_\lambda dx^\lambda \left[\frac{\partial}{\partial x^{\lambda_1}} \cdots \frac{\partial}{\partial x^{\lambda_n}} \right] = \alpha_\lambda.$$

This holds $\forall \lambda \Rightarrow \{dx^\lambda\}$ are indept.

We now show $\{dx^\lambda\}$ spans $\mathcal{L}^n$.

Let $\omega \in \mathcal{L}^n$. Then

$$\begin{aligned} \omega(X_1 \cdots X_n) &= \omega\left(X_1^{i_1} \frac{\partial}{\partial x^{i_1}} \cdots X_n^{i_n} \frac{\partial}{\partial x^{i_n}}\right) \\ &= X_1^{i_1} \cdots X_n^{i_n} \omega\left(\frac{\partial}{\partial x^{i_1}} \cdots \frac{\partial}{\partial x^{i_n}}\right) \end{aligned}$$

where each i_k sums from $1 \cdots n$. Now partition $\{1, \dots, n\}$ into $\binom{n}{k}$ increasing sequences $\mu_1, \dots, \mu_{\binom{n}{k}}$, each of k -length.

Thus

$$\omega(X_1 \cdots X_n) = \sum_{i=1}^{\binom{n}{n}} \sum_{\lambda=\pi(\mu_i)} X_1^{\lambda_1} \cdots X_n^{\lambda_n} \omega\left(\frac{\partial}{\partial X^{\lambda_1}} \cdots \frac{\partial}{\partial X^{\lambda_n}}\right)$$

where $\mu_i = (i_1 \cdots i_n)$, $\pi(\mu_i) = (i_{\pi(1)} \cdots i_{\pi(n)})$. Thus

$$\omega(X_1 \cdots X_n) = \sum_{i=1}^{\binom{n}{n}} \sum_{\pi} (-1)^{\pi} X_1^{i_{\pi(1)}} \cdots X_n^{i_{\pi(n)}} \omega\left(\frac{\partial}{\partial x^{i_1}} \cdots \frac{\partial}{\partial x^{i_n}}\right)$$

$$= \sum_{i=1}^{\binom{n}{n}} dx^{\mu_i}(X_1 \cdots X_n) \underbrace{\omega\left(\frac{\partial}{\partial x^{i_1}} \cdots \frac{\partial}{\partial x^{i_n}}\right)}_{\alpha_{\mu_i}}$$

$$\omega = \sum_{i=1}^{\binom{n}{n}} \alpha_{\mu_i} dx^{\mu_i} \quad \checkmark$$

• Wedge Product: Our defn of $dx^{i_1} \wedge \dots \wedge dx^{i_k}$ is as a det operator,

$$dx^{i_{\pi(1)}} \wedge \dots \wedge dx^{i_{\pi(k)}} = (-1)^\pi dx^{i_1} \wedge \dots \wedge dx^{i_k}.$$

We can extend $\wedge$ to operation on $\bigcup_k \mathcal{L}^k$ by

$$\mathcal{L}^k \ni \alpha = \alpha_\lambda dx^\lambda$$

linearity &
antisymmetry

$$\mathcal{L}^l \ni \beta = \beta_\mu dx^\mu$$

Defn: $\alpha \wedge \beta \in \mathcal{L}^{k+l}$

$$(1) \quad \alpha \wedge \beta = \alpha_{[\lambda} \beta_{\mu]} dx^\lambda \wedge dx^\mu \quad \left(\begin{array}{l} \text{sum over} \\ \text{increasing} \\ \lambda, \mu \end{array} \right)$$

$$(2) \quad dx^\lambda \wedge dx^\mu \equiv dx^{\lambda_1} \wedge \dots \wedge dx^{\lambda_k} \wedge dx^{\mu_1} \wedge \dots \wedge dx^{\mu_l}.$$

Thm: (1) & (2) define $\alpha \wedge \beta$ as a unique ele of $\mathcal{L}^{k+l}$, indep of which coord system (1), (2) are performed in.

• Said Differently: assuming (2) together with $\alpha \wedge B = (-1)^{k+l} B \wedge \alpha$ and multilinearity in each slot,

$$\alpha \wedge (C_1 B^1 + C_2 B^2) = C_1 \alpha \wedge B^1 + C_2 \alpha \wedge B^2$$

$$(C_1 \alpha^1 + C_2 \alpha^2) \wedge B = C_1 \alpha^1 \wedge B + C_2 \alpha^2 \wedge B,$$

defn (1) is implied.

Problem: It is not so easy to prove directly (and it is a bit remarkable) that (1) defines a coord. indep't operation i.e., given y -words,

$$\alpha = \bar{\alpha}_{\bar{\lambda}} dy^{\bar{\lambda}}$$

$$B = \bar{B}_{\bar{\mu}} dy^{\bar{\mu}}$$

we need

$$\alpha \wedge B = \alpha_{\lambda} B_{\mu} dx^{\lambda} \wedge dx^{\mu} = \bar{\alpha}_{\bar{\lambda}} \bar{B}_{\bar{\mu}} dy^{\bar{\lambda}} \wedge dy^{\bar{\mu}}.$$

We'd like to ~~show~~ argue that

$$\alpha_{\bar{\lambda}} = \alpha_{|\lambda|} \frac{\partial x^{\lambda}}{\partial y^{\bar{\lambda}}}$$

$$\beta_{\bar{\mu}} = \beta_{|\mu|} \frac{\partial x^{\mu}}{\partial y^{\bar{\mu}}}$$

$$\Rightarrow \underbrace{\alpha_{|\lambda|} \beta_{|\mu|}}_{\text{LHS}} dx^{\lambda} \wedge dx^{\mu} = \underbrace{\alpha_{|\bar{\lambda}|} \beta_{|\bar{\mu}|} \frac{\partial x^{\lambda}}{\partial y^{\bar{\lambda}}} \frac{\partial x^{\mu}}{\partial y^{\bar{\mu}}}}_{\text{RHS}} dy^{\bar{\lambda}} \wedge dy^{\bar{\mu}}$$

Problem: although $dx^{\lambda} \wedge dx^{\mu} = 0$ unless

$\gamma = (\lambda, \mu) = (\lambda_1, \dots, \lambda_n, \mu_1, \dots, \mu_r)$ ~~is a~~ is a reordering of $\xi_1, \dots, \xi_{n+r}$, $\exists$ complicated coeff's when we use antisymmetry to ~~reorder~~ collect all the coeff's of the permutation of the increasing sequence. I.e. how to relate

$$\underbrace{\alpha_{|\lambda|} \beta_{|\mu|}}_{\text{LHS}} dx^{\lambda} \wedge dx^{\mu} = \underbrace{\gamma_{|\gamma|}}_{\text{RHS}} dx^{\gamma}$$

- (28)
- To prove theorem we give an invariant formula for λ ; then prove that this agrees on k -forms $dx^{i_1} \wedge \dots \wedge dx^{i_k}$ and is multilinear & anti-symmetric $\Rightarrow$ (1) holds.

Defn (1) Let $\omega^1, \dots, \omega^n$ be elements of T^*M . Define $\omega^1 \wedge \dots \wedge \omega^n$ to be the n -form that acts on n -vectors $X_1, \dots, X_n$ by

$$\omega^1 \wedge \dots \wedge \omega^n (X_1, \dots, X_n) = \det [\omega^i(X_j)]. \quad (3)$$

- Clearly (3) is multilinear, antisymmetric, and agrees with defn of $dx^{i_1} \wedge \dots \wedge dx^{i_k}$
- Any k -form that can be expressed as $\omega^1 \wedge \dots \wedge \omega^k$ some $\omega^i \in T^*M$, is called decomposable.

(29)

Defn (2) Let $\alpha \in \Lambda^k$, $\beta \in \Lambda^l$. Define

$$\begin{aligned} \alpha \wedge \beta (X_1, \dots, X_m, X_{m+1}, \dots, X_{m+l}) \\ = \sum_{|\lambda||\mu|} (-1)^{\pi(\lambda, \mu)} \alpha(X_{\lambda_1}, \dots, X_{\lambda_m}) \beta(X_{\mu_1}, \dots, X_{\mu_l}) \end{aligned} \quad (4)$$

where the sum is over $\lambda = (\lambda_1, \dots, \lambda_m)$, $\mu = (\mu_1, \dots, \mu_l)$

satisfying $\lambda_1 < \dots < \lambda_m$, $\mu_1 < \dots < \mu_l$, where

$(\lambda, \mu) = (\lambda_1, \dots, \lambda_m, \mu_1, \dots, \mu_l)$ is a reordering of $(1, \dots, m+l)$ of parity π .

Note: (4) is an invariant defn $\Rightarrow$ coord. indep meaning.

Theorem: Defn (4) has following properties:

① If α, B are decomposable, then (4) agrees with the defn (3) for $\alpha \wedge B$.

② Antisymmetry: $\alpha \wedge B = (-1)^{kl} B \wedge \alpha$

③ Multilinearity: $\alpha \wedge (c_1 B^1 + c_2 B^2) = c_1 \alpha \wedge B^1 + c_2 \alpha \wedge B^2$
 $(c_1 \alpha^1 + c_2 \alpha^2) \wedge B = c_1 \alpha^1 \wedge B + c_2 \alpha^2 \wedge B$.

④ Associativity: $(\alpha \wedge B) \wedge \gamma = \alpha \wedge (B \wedge \gamma)$

~~Cor: ②, ③ imply $\alpha_{[1} dx^1 \wedge B_{1\mu]} dx^\mu = \alpha_{[1} B_{1\mu]} dx^1 \wedge dx^\mu$~~

Corollary:

⑤ $\alpha_{[1} dx^1 \wedge B_{1\mu]} dx^\mu = \alpha_{[1} B_{1\mu]} dx^1 \wedge dx^\mu$

(Follows directly from ①-③)

Pf of thm :

$$\textcircled{1} \quad \alpha = \omega^1 1 \dots 1 \omega^k \quad \beta = \omega^{k+1} 1 \dots 1 \omega^{k+l}$$

$$\alpha \beta = \sum_{|\lambda||\mu|} (-1)^{\pi(\lambda, \mu)} \det [\omega^i(x_{\lambda_j})]_{i=1, \dots, k} \det [\omega^i(x_{\mu_j})]_{i=k+1, \dots, k+l}$$

$$\det [\omega^i(x_{\lambda_j})] = \sum_{\pi} (-1)^{\pi} \omega^1(x_{\lambda_{\pi(1)}}) \dots \omega^k(x_{\lambda_{\pi(k)}})$$

$$\det [\omega^i(x_{\mu_j})] = \sum_{\pi'} (-1)^{\pi'} \omega^{k+1}(x_{\mu_{\pi'(1)}}) \dots \omega^{k+l}(x_{\mu_{\pi'(l)}})$$

$$\alpha \beta = \sum_{|\lambda||\mu|} (-1)^{\pi(\lambda, \mu)} \left(\sum_{\pi} (-1)^{\pi} \omega^1(x_{\lambda_{\pi(1)}}) \dots \omega^k(x_{\lambda_{\pi(k)}}) \right) \cdot \left(\sum_{\pi'} (-1)^{\pi'} \omega^{k+1}(x_{\mu_{\pi'(1)}}) \dots \omega^{k+l}(x_{\mu_{\pi'(l)}}) \right)$$

$$= \sum_{|\lambda||\mu|} \sum_{\pi, \pi'} (-1)^{\pi} (-1)^{\pi'} (-1)^{\pi(\lambda, \mu)} \omega^1(x_{\lambda_{\pi(1)}}) \dots \omega^k(x_{\lambda_{\pi(k)}})$$

names an arb perm. γ of $(1, \dots, k+l)$

$$= \sum_{\gamma} (-1)^{\gamma} \omega^1(x_{\gamma(1)}) \dots \omega^{k+l}(x_{\gamma(k+l)}) = \det [\omega^i(x_j)] \quad \checkmark$$

Proof of Thm

② $(-1)^{kl}$ is the parity associated with moving l numbers from the right to the left of k numbers in a permutation. w.o. otherwise changing the ordering:

$$\begin{aligned}
 (\lambda_1 \cdots \lambda_k \mu_1 \cdots \mu_l) &\rightarrow (\mu_1 \lambda_1 \cdots \lambda_k \mu_2 \cdots \mu_l) \\
 &\quad \quad \quad \underbrace{\hspace{1.5cm}}_k \\
 &\rightarrow (\mu_1 \mu_2 \lambda_1 \cdots \lambda_k \mu_3 \cdots \mu_l) \\
 &\quad \quad \quad \underbrace{\hspace{1.5cm}}_k \\
 &\quad \quad \quad \vdots \\
 &\rightarrow (\mu_1 \cdots \mu_l \lambda_1 \cdots \lambda_k)
 \end{aligned}$$

Thus: $B \wedge \alpha (X_1 \cdots X_{k+l}) = \sum_{|\lambda||\mu|} (-1)^{\pi(\lambda\mu)} \alpha (X_{\lambda_1} \cdots X_{\lambda_k}) \cdot B(X_{\mu_1} \cdots X_{\mu_l})$

HW $\} = \sum_{|\lambda||\mu|} (-1)^{kl} (-1)^{\pi(\mu\lambda)} \alpha (X_{\lambda_1} \cdots X_{\lambda_k}) B(X_{\mu_1} \cdots X_{\mu_l}) = (-1)^{kl} \alpha \wedge B$ ✓

$$\textcircled{3} \quad \propto \lambda (C_1 B^1 + C_2 B^2)$$

$$= \sum_{|\lambda||\mu|} (-1)^{\pi(\lambda, \mu)} \alpha(X_{\lambda_1} \cdots X_{\lambda_n}) \left\{ C_1 B^1(X_{\mu_1} \cdots X_{\mu_r}) + C_2 B^2(X_{\mu_1}, \dots, X_{\mu_r}) \right\}$$

$$= C_1 \sum_{|\lambda||\mu|} (-1)^{\pi(\lambda, \mu)} \alpha(X_{\lambda_1} \cdots X_{\lambda_n}) B^1(X_{\mu_1} \cdots X_{\mu_r}) + C_2 \sum_{|\lambda||\mu|} (-1)^{\pi(\lambda, \mu)} \alpha(X_{\lambda_1} \cdots X_{\lambda_n}) B^2(X_{\mu_1} \cdots X_{\mu_r})$$

$$= C_1 \propto \lambda B_1 + C_2 \propto \lambda B_2.$$

From Here we conclude $\textcircled{5}$ holds, so it suffices to use $\textcircled{5}$ to verify $\textcircled{4}$

H.W

$$\textcircled{4} \quad \alpha = \alpha_{|\lambda|} dx^\lambda \quad \beta = \beta_{|\mu|} dx^\mu \quad \gamma = \gamma_{|\tau|} dx^\tau \quad \textcircled{34}$$

$$(\alpha \wedge \beta) \wedge \gamma = (\alpha_{|\lambda|} dx^\lambda \wedge \beta_{|\mu|} dx^\mu) \wedge \gamma_{|\tau|} dx^\tau$$

$$= (\alpha_{|\lambda|} \beta_{|\mu|} dx^\lambda \wedge dx^\mu) \wedge \gamma_{|\tau|} dx^\tau$$

$$= \alpha_{|\lambda||\mu||\tau|} (dx^\lambda \wedge dx^\mu) \wedge dx^\tau$$

$$= \alpha_{|\lambda||\mu||\tau|} dx^\lambda \wedge (dx^\mu \wedge dx^\tau)$$

$$= \alpha \wedge (\beta \wedge \gamma)$$

↑
By orig defn as
a determinant ✓

Ans.

Another way to wedge product: (Spirvak) (35)

Defn: $T \in \mathcal{Y}_n^0$

$$\text{Alt}(T)(X_1 \cdots X_n) = \frac{1}{n!} \sum_{\pi} (-1)^{\pi} T(X_{\pi(1)} \cdots X_{\pi(n)})$$

Thm

① $\text{Alt}(T) = T$ iff $T \in \mathcal{L}^k$

② $\text{Alt}(T) \in \mathcal{L}^k$ "the anti-symmetrization of T "

③ $S \wedge T = \frac{(k+l)!}{k! l!} \text{Alt}(S \otimes T)$ $S \in \mathcal{Y}_n^0, T \in \mathcal{Y}_l^0$

• Metric g induces inner product $\langle, \rangle$ on $\Lambda^k \Lambda_n$ that determines k -volumes:

• Notation: Given coord system $\underline{x}$:

$$dx^\lambda \equiv dx^{\lambda_1} \wedge \dots \wedge dx^{\lambda_k}$$

$$\frac{\partial}{\partial x^\lambda} \equiv \frac{\partial}{\partial x^{\lambda_1}} \wedge \dots \wedge \frac{\partial}{\partial x^{\lambda_k}}$$

where $\{dx^\lambda\}$, $\{\frac{\partial}{\partial x^\lambda}\}$ are bases for Λ^k , Λ_k resp.

• Eg., $X_1 \wedge \dots \wedge X_k \equiv a_{\lambda_1} dx^{\lambda_1}$ defines a unique ele. of Λ_k by antisymmetry/multi-linearity of $\wedge$.

• We wish to identify $X_1 \wedge \dots \wedge X_k$ with k -parallelotope spanned by $X_1, \dots, X_k$.

Goal: define metric on Λ_k so

$$\langle X_1 \wedge \dots \wedge X_k, X_1 \wedge \dots \wedge X_k \rangle \equiv k\text{-vol}[X_1, \dots, X_k] \text{ as measured by } g.$$

- Not all ele's of Λ_k can be written as $X_1 \wedge \dots \wedge X_k$: There are special:

Defn: $\omega \in \Lambda^k$, $\alpha \in \Lambda_k$ are decomposable if

$\exists \omega^1, \dots, \omega^k \in \Lambda^1$, $\alpha_1, \dots, \alpha_k \in \Lambda_1$ st

$$\omega = \omega^1 \wedge \dots \wedge \omega^k$$

$$\alpha = \alpha_1 \wedge \dots \wedge \alpha_k,$$

respectively.

Lemma ① Let $X_1, \dots, X_k \in T_p M$, $\{X_i\}$ indep't.

Let $\text{Span}\{X_i\} = \text{Span}\{N_i\}$, $i=1, \dots, k$. Then

$$X_1 \wedge \dots \wedge X_k = \det(X_i^j) N_1 \wedge \dots \wedge N_k$$

where $X_i = X_i^\sigma N_\sigma$.

P.f.

$$\begin{aligned}
 X_1 \wedge \dots \wedge X_n &= X_1^\sigma N_\sigma \wedge \dots \wedge X_n^\sigma N_\sigma \\
 &= \sum_{\pi} (-1)^\pi X_1^{\pi(1)} \dots X_n^{\pi(n)} N_1 \wedge \dots \wedge N_n \\
 &= \det(X_j^i) N_1 \wedge \dots \wedge N_n \quad \checkmark
 \end{aligned}$$

Cor: $X_1 \wedge \dots \wedge X_n \neq 0$ iff $\{X_i\}$ are indept.

P.f. Choose $N_1, \dots, N_n$ indept st. $\text{Span}\{N_1, \dots, N_n\} = \text{Span}\{X_1, \dots, X_n\}$.

Then as above

$$\begin{aligned}
 X_1 \wedge \dots \wedge X_n &= X_1^\sigma N_\sigma \wedge \dots \wedge X_n^\sigma N_\sigma \\
 &= \det(X_j^i) N_1 \wedge \dots \wedge N_n \\
 &= 0 \quad \text{if } n < n \quad \checkmark
 \end{aligned}$$

Cor: $\wedge^k = \wedge_n = 0$ for $k > n$.

(38)

Assume a metric g with components in ON-frame given by

$$g_{ij} = d_{ij} = \text{diag} \{ \sigma_1, \dots, \sigma_n \}$$

$$\sigma_1 = \dots = \sigma_r = -1, \sigma_{r+1} = \dots = \sigma_n = 1, \text{sgn}(g) = n - r$$

• This induces volume form

$$\eta = \sqrt{|\det g|} dx^1 \wedge \dots \wedge dx^n$$

• The metric also induces a map

$$T^*M \leftrightarrow TM$$

$$X^i \frac{\partial}{\partial x^i} \leftrightarrow X_i dx^i$$

$$X^i = g^{i\sigma} X_\sigma$$

$$g^{ij} = (g_{ij})^{-1}$$

• This extends to a map (by raising/lowering)

$$y^k_0 \leftrightarrow y^0_k$$

$$T_{i_1 \dots i_n} \frac{\partial}{\partial x^{i_1}} \otimes \dots \otimes \frac{\partial}{\partial x^{i_n}} \leftrightarrow T_{i_1 \dots i_n} dx^{i_1} \otimes \dots \otimes dx^{i_n}$$

$$T_{i_1 \dots i_n} = T^{\sigma_1 \dots \sigma_n} g_{\sigma_1 i_1} \dots g_{\sigma_n i_n}$$

Lemma 2 This induces a ^{linear} mapping

$$L: \mathcal{L}^k \leftrightarrow \mathcal{L}_k$$

$$g_{\lambda_1} dx^{\lambda_1} \leftrightarrow a^{|\lambda|} \frac{\partial}{\partial x^{\lambda}}$$

$$a_{\lambda} \equiv a_{\lambda_1 \dots \lambda_n} = a^{\sigma_1 \dots \sigma_n} g_{\sigma_1 \lambda_1} \dots g_{\sigma_n \lambda_n} \quad (*)$$

Pf. By (*) page 20,

$$g_{\lambda_1} dx^{\lambda_1} = a_{i_1 \dots i_n} dx^{i_1} \otimes \dots \otimes dx^{i_n}$$

where $a_{\lambda} = a_{\lambda_1 \dots \lambda_n}$ for λ increasing & $a_{i_1 \dots i_n}$ transform by (*) ✓

start

(40)

• Defn: Let $\omega^1, \omega^2 \in \mathcal{L}^k$; $\alpha_1, \alpha_2 \in \mathcal{L}^k$.

$$\langle \omega^1, \omega^2 \rangle_* = \omega^1(L\omega^2)$$

$$\langle \alpha_1, \alpha_2 \rangle_* = (L^{-1}\alpha_1)(\alpha_2)$$

• Lemma 3 In coordinates:

$$\langle a_{|\lambda|} dx^\lambda, b_{|\lambda|} dx^\lambda \rangle_* = a^\lambda b_\lambda = k! a^{|\lambda|} b_{|\lambda|}$$

$$\langle a^{|\lambda|} \frac{\partial}{\partial x^\lambda}, b^{|\lambda|} \frac{\partial}{\partial x^\lambda} \rangle_* = a_\lambda b^\lambda = k! a_{|\lambda|} b^{|\lambda|}$$

Pf. For the 2nd: $\alpha_1 = a^{|\lambda|} \frac{\partial}{\partial x^\lambda} \Rightarrow$

$$L^{-1}\alpha_1 = a_{|\lambda|} dx^\lambda$$

But

$$(a_{|\lambda|} dx^\lambda) (b^{|\lambda|} \frac{\partial}{\partial x^\lambda}) = (a_{i_1 \dots i_k} dx^{i_1} \otimes \dots \otimes dx^{i_k})$$

$$= a_{i_1 \dots i_k} b^{i_1 \dots i_k} = a_\lambda b^\lambda \sum_{|\lambda| \uparrow} \sum_{\pi} (-1)^{\pi(\lambda)} a_{|\lambda|} (1)^{\pi(\lambda)} b^\lambda = k! a_{|\lambda|} b^{|\lambda|}$$

(41)

• Let $\underline{x}$ be an ON basis. In this case

$$\left| \left\langle \frac{\partial}{\partial x^{i_1}} 1 \cdots 1 \frac{\partial}{\partial x^{i_n}}, \frac{\partial}{\partial x^{i_1}} 1 \cdots 1 \frac{\partial}{\partial x^{i_n}} \right\rangle \right| = n!$$

because the metric induced mapping $\Lambda^k \hookrightarrow \Lambda_n$ yields $dx^\lambda \leftrightarrow \pm \frac{\partial}{\partial x^\lambda}$

Lemma 4 if $\underline{x}$ is O.N., then $(d_{ij} = \text{diag}(\sigma_1, \dots, \sigma_n))$

$$dx^\lambda \leftrightarrow \sigma_{\lambda_1} \cdots \sigma_{\lambda_n} \frac{\partial}{\partial x^\lambda}.$$

Proof: $\alpha = \frac{\partial}{\partial x^\lambda} \Rightarrow \alpha = a^{i_1 \cdots i_n} \frac{\partial}{\partial x^{i_1}} \otimes \cdots \otimes \frac{\partial}{\partial x^{i_n}}$

where

$$a^{i_1 \cdots i_n} = \begin{cases} (-1)^\pi & \pi = (i_1 \cdots i_n) \rightarrow (\lambda_1 \cdots \lambda_n) = \lambda \\ 0 & \text{o.w.} \end{cases}$$

Thus

$$\begin{aligned} a_{i_1 \cdots i_n} &= a^{j_1 \cdots j_n} d_{j_1 i_1} \cdots d_{j_n i_n} \\ &= \sigma_{i_1} \cdots \sigma_{i_n} a^{i_1 \cdots i_n} \quad \leftarrow \text{no sum} \end{aligned}$$

$$\Rightarrow \boxed{\alpha = \frac{\partial}{\partial x^\lambda} \leftrightarrow (\sigma_{\lambda_1} \cdots \sigma_{\lambda_n}) dx^\lambda} \quad \checkmark$$

In order that $\{dx^\lambda\}, \{\frac{\partial}{\partial x^\lambda}\}$ be orthon. frame for $\omega^\lambda, \omega_\lambda$ resp., we redefine $\langle, \rangle_*$

Defn: $\langle \omega^i, \omega^j \rangle = \frac{1}{n!} \langle \omega^i, \omega^j \rangle_*$

$$\langle \alpha_i, \alpha_j \rangle = \frac{1}{n!} \langle \alpha_i, \alpha_j \rangle_*$$

so that

$$\langle \alpha_{|\lambda|} dx^\lambda, \beta_{|\lambda|} dx^\lambda \rangle = \alpha_{|\lambda|} \beta^\lambda$$

$$\langle \alpha^{|\lambda|} \frac{\partial}{\partial x^\lambda}, \beta^{|\lambda|} \frac{\partial}{\partial x^\lambda} \rangle = \alpha^{|\lambda|} \beta_\lambda$$

↑
FIP

Let $X_1, \dots, X_n$ be linearly indept vectors in $T_p M$, and let $N_1, \dots, N_m, N_{m+1}, \dots, N_n$ be unit orthogonal vectors satisfying

$$\langle N_i, N_i \rangle = \sigma_{\pi(i)} = \tau_i,$$

$$\text{Span} \{N_1, \dots, N_n\} = \text{Span} \{X_1, \dots, X_n\}.$$

(I.e., we cannot ensure that $N_1, \dots, N_n$ is an o.n. basis since τ_i depends on sign $\langle X_i, X_i \rangle$, and we are not free to make $\tau_i = \sigma_i$.)

• Theorem (Main)

$$\textcircled{1} \det \langle X_i, X_j \rangle = \tau_1 \dots \tau_m [\det X_j^i]^2 \text{ where}$$

$$X_j = X_j^i N_i$$

$$\textcircled{2} |X_1, \dots, X_m|^2 = \langle X_1, \dots, X_m, X_1, \dots, X_m \rangle = \det \langle X_i, X_j \rangle$$

$$\textcircled{3} \eta(X_1, \dots, X_m, N_{m+1}, \dots, N_n)^2 = |\det \langle X_i, X_j \rangle|$$

where $\eta = \sqrt{g} dx^1 \wedge \dots \wedge dx^n$ is the vol-form for g

(44)

• Note: $k\text{-Vol}[X_1, \dots, X_n] \equiv k\text{-volume of } \Pi\text{-space spanned by } X_1, \dots, X_n \text{ induced by } g \text{ is defined as}$

Defn: $k\text{-Vol}[X_1, \dots, X_n]^2 = |\eta(X_1, \dots, X_n, N_{n+1}, \dots, N_n)|$
 > 0

Thus: ① & ② of Thm $\Rightarrow$

$$k\text{-Vol}[X_1, \dots, X_n]^2 = \pm |X_1 \wedge \dots \wedge X_n|^2$$

↑
when g is not pos def,
the norm of a vector can
be negative

We can thus define the signed $k\text{-vol}[X_1, \dots, X_n]^2$:

Defn: The signed $k\text{-vol}[X_1, \dots, X_n]^2$ is

$$|X_1 \wedge \dots \wedge X_n|^2 = \det \langle X_i, X_j \rangle.$$

• Thus the final picture emerges for the metric on $\mathcal{L}_k$ that defines k -dim volume induced by g :

► Fix k -parallelotope $X_1 \wedge \dots \wedge X_k \in \mathcal{L}_k$

► Choose $\underline{x}$ to be an arbitrary orthon. frame at p :

$$\left\langle \frac{\partial}{\partial x^i}, \frac{\partial}{\partial x^j} \right\rangle = d_{ij} = \text{diag}\{\sigma_1 \dots \sigma_n\}$$

$\sigma_1 = \dots = \sigma_r = -1$, $\sigma_{r+1} = \dots = \sigma_n = 1$. Let

$$X_i = X_i^{\hat{\alpha}} \frac{\partial}{\partial x^{\hat{\alpha}}}$$

► Then:

$$X_1 \wedge \dots \wedge X_k = X_1^{\hat{\alpha}_1} \frac{\partial}{\partial x^{\hat{\alpha}_1}} \wedge \dots \wedge X_k^{\hat{\alpha}_k} \frac{\partial}{\partial x^{\hat{\alpha}_k}}$$

$$= X_1^{\hat{\alpha}_1} \dots X_k^{\hat{\alpha}_k} \frac{\partial}{\partial x^{\hat{\alpha}_1}} \wedge \dots \wedge \frac{\partial}{\partial x^{\hat{\alpha}_k}}$$

$$= \sum_{|\lambda|} \sum_{\pi} (-1)^{\pi} X_1^{\pi(\lambda_1)} \dots X_k^{\pi(\lambda_k)} \frac{\partial}{\partial x^{\lambda_1}} \wedge \dots \wedge \frac{\partial}{\partial x^{\lambda_k}}$$

$$= \sum_{|\lambda|} \det^{\lambda} \begin{bmatrix} \frac{1}{x_1} & \dots & \frac{1}{x_k} \\ 1 & & 1 \end{bmatrix} \frac{\partial}{\partial x^{\lambda_1}} \wedge \dots \wedge \frac{\partial}{\partial x^{\lambda_k}}$$

► Let

$$p^\lambda \equiv \det^\lambda \begin{bmatrix} \dot{x}_1 & \dots & \dot{x}_n \\ 1 & & 1 \end{bmatrix}$$

Then $p^\lambda \equiv k$ -volume of the projection of $[x_1 \dots x_n]$ onto the $x_1 \dots x_n$ -coordinate axes as measured in the coordinate system x

I.e., p^λ is the coordinate k -volume of this projection.

► Thus: $X_1 \wedge \dots \wedge X_n = p^{|\lambda|} \frac{\partial}{\partial x^\lambda}$

► Therefore:

$$|X_1 \wedge \dots \wedge X_n|^2 = \langle X_1 \wedge \dots \wedge X_n, X_1 \wedge \dots \wedge X_n \rangle$$

$$= (p^{|\lambda|})^2 \left\langle \frac{\partial}{\partial x^\lambda}, \frac{\partial}{\partial x^\lambda} \right\rangle$$

$$= \text{signed } k\text{-vol } [x_1 \dots x_n]^2$$

Sum repeated up-down increasing sequence

and

(47)

$$\left\langle \frac{\partial}{\partial X^\lambda}, \frac{\partial}{\partial X^\lambda} \right\rangle = \sigma_{\lambda_1} \cdots \sigma_{\lambda_n} \equiv \sigma_\lambda$$

Here: $\sigma_\lambda = \sigma_{\lambda_1} \cdots \sigma_{\lambda_n} = \pm 1$ depending on the signature of the metric

Conclude:

$$\begin{aligned} |X_1 \wedge \cdots \wedge X_n|^2 &= \sum_{|\lambda|} \sigma_\lambda (P^\lambda)^2 \\ &= (\kappa\text{-vol}[X_1, \dots, X_n] \text{ induced by } g) \end{aligned} \quad (*)$$

► In particular, if $d_{ij} = \delta_{ij}$, $\sigma_1 = \cdots = \sigma_n = 1$ the pos def case, then we obtain the standard projection thru an or. basis:

$$|X_1 \wedge \cdots \wedge X_n|^2 \equiv \sum_{|\lambda|} (P^\lambda)^2$$

(48)

► For the non-pos def. case, if we choose our orthogonal unit frame so that

$$\frac{\partial}{\partial x^i} = N_i, \quad i = 1, \dots, k \text{ (at } p),$$

and $\text{Span}\{X_i\} = \text{Span}\{N_i\}$, then $(*)$ $\langle N_i, N_j \rangle = \hat{c}_i = \sigma_{\pi(i)} = \pm 1$
 reduces to

$$|X_1 \wedge \dots \wedge X_k|^2 = \overbrace{\tau_1 \dots \tau_k}^{\sigma_\lambda} (P^\lambda)^2, \quad \lambda = (1, \dots, k)$$

$$P^\lambda = \det(X_i^j), \quad X_i = X_i^j N_j$$

$i, j = 1, \dots, k.$

Conclude: "In on frames aligned with the Π -oriped, metric k -volume = $\pm$ coordinate k -volume of the Π -oriped. In on. frames not-aligned with the Π -oriped, the $\pm$ sign in $|X_1 \wedge \dots \wedge X_k|^2 = \sigma_{|\lambda|} (P^\lambda)^2$ account for "Lorentz Contraction" in non-pos definite metrics."

□ Proof of Theorem (Main):

(49)

- For ① note

$$\begin{aligned}\det \langle X_i, X_j \rangle &= \det \langle X_i^p N_e, X_j^m N_j \rangle \\ &= \det (X_i^p X_j^m \langle N_e, N_m \rangle) \\ &= \det \left(H^t \cdot \begin{bmatrix} \tilde{\tau}_1 & 0 \\ 0 & \tilde{\tau}_m \end{bmatrix} H \right), H = (X_j^i) \\ &= [\det (X_j^i)]^2 \tilde{\tau}_1 \cdots \tilde{\tau}_m \quad \checkmark\end{aligned}$$

- pf of ②: $(|X, 1 \cdots 1 X_m|^2 = \det \langle X_i, X_j \rangle)$

We have

$$X, 1 \cdots 1 X_m = \det (X_j^i) N, 1 \cdots 1 N_m,$$

so that

$$|X, 1 \cdots 1 X_m|^2 = \langle X, 1 \cdots 1 X_m, X, 1 \cdots 1 X_m \rangle$$

$$= (\det X_j^i)^2 \langle N, 1 \cdots 1 N_m, N, 1 \cdots 1 N_m \rangle$$

$$= (\det X_j^i)^2 \tilde{\tau}_1 \cdots \tilde{\tau}_m = \det \langle X_i, X_j \rangle (\tilde{\tau}_1 \cdots \tilde{\tau}_m)^2$$

$$= \det \langle X_i, X_j \rangle \quad \checkmark$$

• Proof of ① : $\eta(X_1, \dots, X_k, N_{k+1}, \dots, N_n)^2 =$

$$= \det \langle X_i, X_j \rangle$$

$$\eta(X_1, \dots, X_k, N_{k+1}, \dots, N_n)$$

$$= \eta(X_1^{i_1} N_{i_1}, \dots, X_k^{i_k} N_{i_k}, N_{k+1}, \dots, N_n)$$

$$= \det(X_j^{i_i}) \eta(N_1, \dots, N_n)$$

~~$$= (-1)^{\pi} \det X_j^{i_i} = (-1)^{\pi_1 + \dots + \pi_k} \det \langle X_{i_1}, X_{i_1} \rangle \dots \det \langle X_{i_k}, X_{i_k} \rangle$$~~

Thus:

$$\eta(X_1, \dots, X_k, N_{k+1}, \dots, N_n)^2 = (\det X_j^{i_i})^2$$

$$= \tau_1 \dots \tau_k \det \langle X_{i_i}, X_{i_i} \rangle$$

Summary:

- Given M^n , no metric, coord syst $\underline{x}: U \rightarrow \mathbb{R}^n$

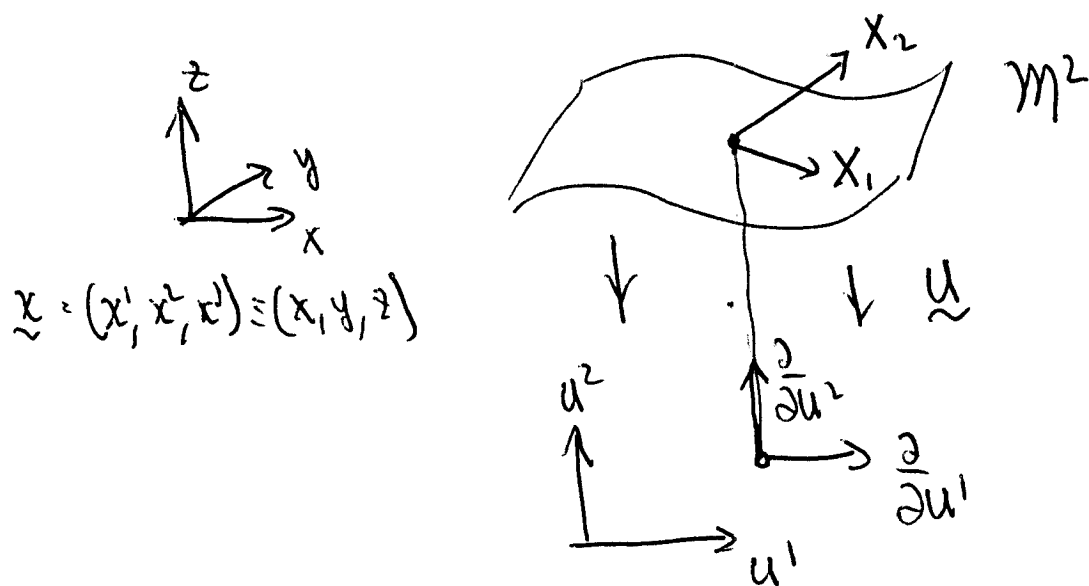
η an n -form $\eta \in \mathcal{L}^n$

$$\dim \mathcal{L}^n = 1 \Rightarrow \eta = \eta(x) dx^1 \wedge \dots \wedge dx^n$$

Defn: $\int_U \eta = \int_{\underline{x}(U)} \eta(x) dx^1 \dots dx^n$

Thm: the integral transforms properly

- What about generalizing the $\mathbb{R}^3$ picture?



$$\tilde{u}^{-1} \equiv \tilde{x}(u)$$

$$= (x(u^1, u^2), y(u^1, u^2), z(u^1, u^2))$$

$$X_i = \tilde{u}^{-1*} \frac{\partial}{\partial u^i} = \frac{\partial x^j}{\partial u^i} \frac{\partial}{\partial x^j} \in \mathbb{R}^3$$

(50) (2)

$|X_1 \times X_2| = \text{area of parallelogram } [X_1, X_2] \text{ in } \mathbb{R}^3$

$$\int_{M^2} ds = \int_{\mathbb{R}^2} |X_1 \times X_2| du^1 du^2$$

• To generalize, let $M^n \supset M^k$, g a metric on M^n

and $\underline{u} : M^k \rightarrow \mathbb{R}^k$ a coord syst on M^k .

$$\underline{u}^{-1}_* \left(\frac{\partial}{\partial u^i} \right) = X_i \in TM^k \subseteq TM^n$$

$|X_1 \wedge \dots \wedge X_k| \equiv \text{"k-vol of parallelepiped inherited from } g \text{ on } M^n \text{"}$

Defn: $\int_{M^k} ds \equiv \int_{\mathbb{R}^k} |X_1 \wedge \dots \wedge X_k| du^1 \dots du^k$

- Defn: M^n orientable if $\exists$ non-vanishing k -form on M^n . If $\underline{u}$ global coord system, then $X_1 \wedge \dots \wedge X_n \neq 0 \Rightarrow X_1 \wedge \dots \wedge X_n$ is a ~~non-vanishing~~ non-vanishing k -form on M^n .

Let: $\sigma = \frac{X_1 \wedge \dots \wedge X_n}{|X_1 \wedge \dots \wedge X_n|}$ denote the "unit"

k -form on M^n . We call σ a (normalized) orientation on M^n .

Let $\omega \in \mathcal{L}^k(M^n)$.

- Defn: $\int_{M^n} \omega = \int_{M^n} \langle \omega, \sigma \rangle \sigma$
 $\uparrow$
 component of ω in direction of σ

- We can extend to more than one coord syst. by "partition of unity".

Application: Spacetime.

- In Einstein's theory of gravity, gravitational effects are determined by a Lorentzian metric g defined on the ~~the~~ 4-d manifold Spacetime:
I.e., in a coordinate system x that is orthonormal at P ,

$$\frac{\partial}{\partial x^0} \equiv \frac{\partial}{\partial ct} \quad (\text{at } P)$$

$$\frac{\partial}{\partial x^1} = \frac{\partial}{\partial x}, \quad \frac{\partial}{\partial x^2} = \frac{\partial}{\partial y}, \quad \frac{\partial}{\partial x^3} = \frac{\partial}{\partial z}, \quad (\text{at } P)$$

and

$$g_{ij} = \text{diag}(-1, 1, 1, 1).$$

- The metric g determines proper time:
along time-like curves and length in meters
along spacelike curves. We assume $c = \text{speed}$
of light in $\frac{\text{meters}}{\text{sec}}$, so that x^0 measures time in meters,
 $x^0 = ct$, $t = \text{seconds}$, $c = \text{speed of light in } \frac{\text{meters}}{\text{sec}}$.

I.e. a curve $c(s)$ thru spacetime is timelike if

$$\left\langle \frac{dc}{ds}, \frac{dc}{ds} \right\rangle < 0$$

all along the curve, and is spacelike if

$$\left\langle \frac{dc}{ds}, \frac{dc}{ds} \right\rangle > 0$$



all along the curve.

- Timelike curves correspond to the world line of physical bodies which move at less than the speed of light

- If c is timelike,

$$\Delta\tau = \int_{\xi_1}^{\xi_2} \sqrt{-\left\langle \frac{dc}{ds}, \frac{dc}{ds} \right\rangle} ds$$

is the proper time change as measured by observer traversing the path $c(s)$, $\xi_1 \leq s \leq \xi_2$.

• If c is spacelike,

(5)

$$\Delta S = \int_{\xi_1}^{\xi_2} \sqrt{\left\langle \frac{dc}{d\xi}, \frac{dc}{d\xi} \right\rangle} d\xi$$

is the "length of the curve in meters"

• Light-like curves correspond to trajectories of light rays.

▴ Assumption of special relativity: $\exists$ a coordinate system x on spacetime in which

$$g_{ij} = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix}$$

at every point.

We show next quarter that in general, the best we can do is to make

$$g_{ij}(p) = \text{diag}(-1, 1, 1, 1)$$

$$\frac{\partial}{\partial x^k} g_{ij}(p) = 0,$$

and the 2nd derivatives of g_{ij} measure the curvature of spacetime.

Einstein's equations determine the metric g_{ij} from the energy and momentum densities and their fluxes:

$$G = \kappa T$$

► 2-d Special Relativity: We construct the ON Frames and c.f. the projection thm.

Assume: Lorentzian coords $\underline{x}$: $g_{ij} = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix}$

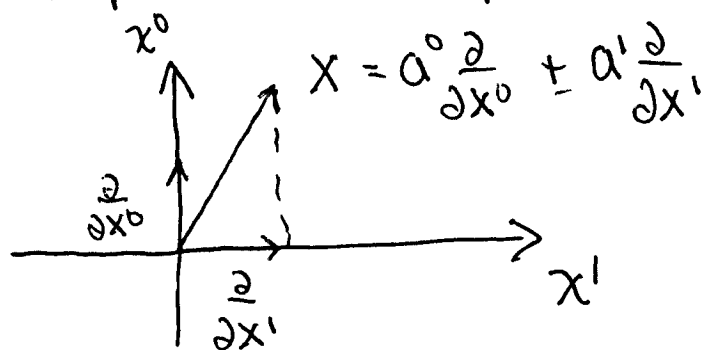
$$x^0 \equiv t, x^1 \equiv x.$$

• Another coordinate system $\underline{y}$ is on. if $g_{ab} = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix}$.

Q: what are the other ON frames x_0, x_1 ?

Ans: Let X be vector $X = a^0 \frac{\partial}{\partial x^0} + a^1 \frac{\partial}{\partial x^1}$

Since metric everywhere the same, we can identify components with pts in the space:

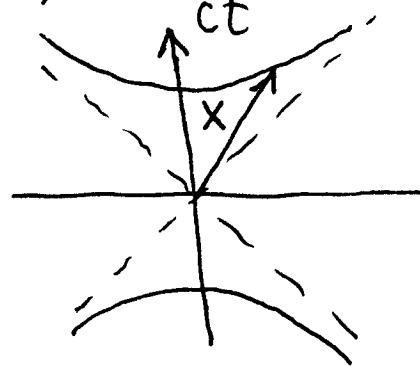


I.e. set $a^0 \equiv ct$, $a^1 \equiv x$, $X = ct \frac{\partial}{\partial x^0} + x \frac{\partial}{\partial x^1}$

Let X be timelike, unit length

$$\langle X, X \rangle = (ct, x) \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} ct \\ x \end{bmatrix} = -(ct)^2 + x^2 = -1$$

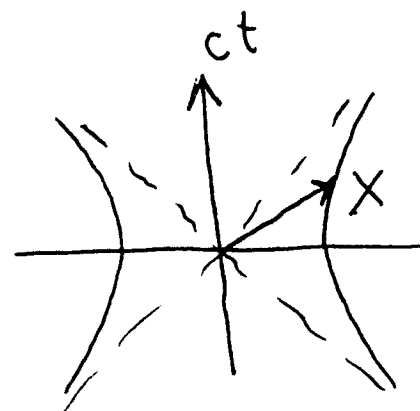
$\Rightarrow X$ lies on unit hyperbola,



Let X be spacelike, unit length

$$\langle X, X \rangle = -(ct)^2 + x^2 = 1$$

$\Rightarrow X$ lies on unit hyperbola,

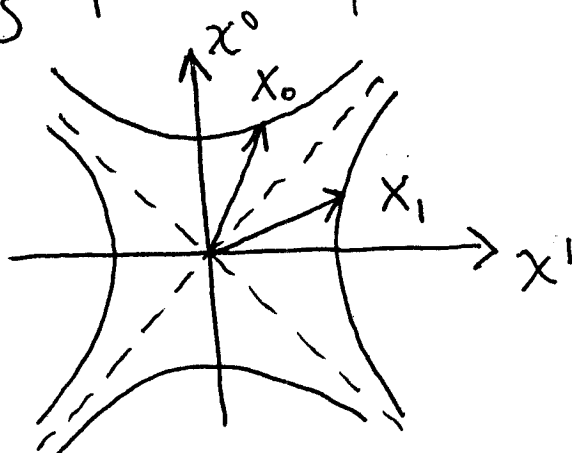


Thus: Assume $\langle X_0, X_1 \rangle = 0$, $\langle X_0, X_0 \rangle = -1$, $\langle X_1, X_1 \rangle = 1$ (56)

$$0 = \langle X_0, X_1 \rangle = (ct_0, x_0) \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} ct_1 \\ x_1 \end{bmatrix} = (-ct_0, x_0) \cdot (ct_1, x_1) = 0$$

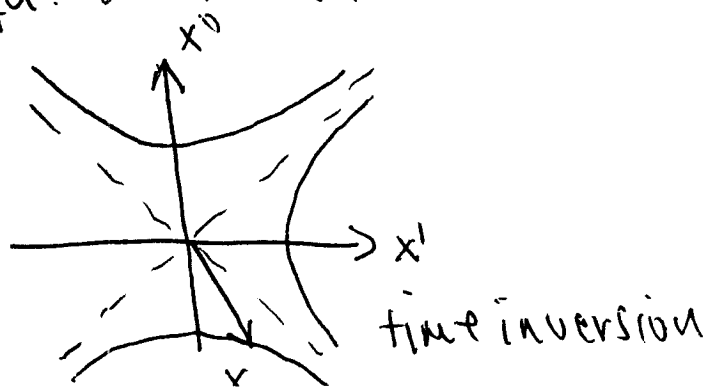
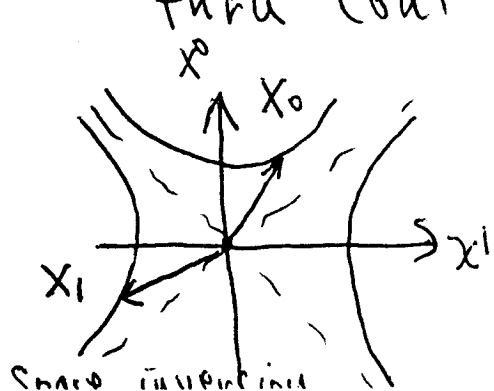
$\Rightarrow "(ct_1, x_1) \parallel \pm (x_0, ct_0) \Rightarrow X_1$ is the reflection of X_0 in line $x = \pm ct$

If we assume X_0 timelike, positive directed and $\{X_0, X_1\}$ positively oriented, then

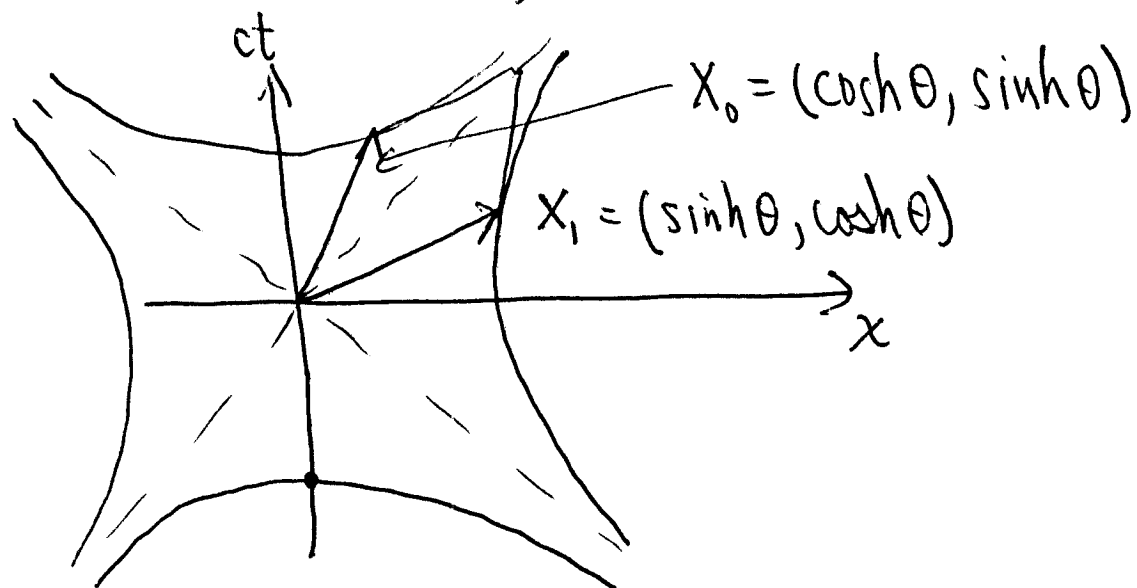


Note ① cannot get to neg oriented frame thru cont transformations

② Cannot get to time inverted or space inverted thru cont sequ. of trans.



Note: $\cosh^2 \theta - \sinh^2 \theta = 1$, thus



Coordinate area of 11-ogram is

$$X_0 \times X_1 = \begin{vmatrix} \cosh \theta & \sinh \theta & 0 \\ \sinh \theta & \cosh \theta & 0 \\ e_1 & e_2 & e_3 \end{vmatrix}$$

$$\begin{vmatrix} \cosh \theta & \sinh \theta & 0 \\ \sinh \theta & \cosh \theta & 0 \\ 0 & 0 & 1 \end{vmatrix} = \cosh^2 \theta - \sinh^2 \theta = 1$$

$\Rightarrow$ x-Coordinate volume of orthonormal bases
is 1 $\Rightarrow$ Volume form $\eta = dx^1 dx^0 = c dx^1 dt$
agrees with coord volume in ON-frames ✓

- Consider 4-d Minkowski spacetime:

Global coordinates on spacetime $\underline{x} = (x^0, \dots, x^3)$

$x^0 = ct$ (x^i in meters $i=1,2,3$, t in seconds,
 c = speed of light meters/sec)

Metric: $\eta_{ij} = \text{diag}(-1, 1, 1, 1)$

Given 3 spacetime vectors

$$X_i = \left\{ X_i^\sigma \frac{\partial}{\partial x^\sigma} \right\} \equiv (X_i^0, \dots, X_i^3)$$

$$\text{3-Vol}[X_1, \dots, X_3]^2 = |X_1 \wedge \dots \wedge X_3|^2$$

$$= \det \langle X_i, X_j \rangle$$

$$= \det \left(\underbrace{-X_i^0 X_j^0}_{\uparrow} + X_i^1 X_j^1 + X_i^2 X_j^2 + X_i^3 X_j^3 \right)$$

minus sign accounts for
 Lorentz contraction